Messengers from the High Energy Universe What we know



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Messengers from the High Energy Universe



Photons Cosmic rays Neutrinos Gravitational Waves ???

The Physics



The Physics (γ , p/nucleus, v , ? beams)

Astrophysics

Sources: AGNs, Pulsares, GRBs, SNRs...

Aceleration mechanisms, propagation





γ sky > 100 GeV Nov 2011

Cosmology/Fundamental Physics

Lorentz symmetry violation, extra-dimensions, dark matter, dark energy ...

Energy scales

10²¹ ZeV (zeta) **10**¹⁸ Black holes / AGN EeV (exa) **10**¹⁵ Supernovae's PeV (peta) **10**¹² TeV (tera) Man Accelerators 10⁹ GeV (giga) Proton mass 10⁶ MeV (mega) Nuclear rectors **10**³ KeV (kilo) X rays, TV **Battery** 1 eV



The beginnings

Free electrical charges in the atmosphere !!!

Coulomb, 1785



Discharge of charged conductors ...

End of XIX century





Radioactivity?

Make good electroscopes!







The two friends Julius Elster and Hans Geiter, gymnasium teachers in Wolfenbuttel, around 1900





Theodor Wulf, German Jesuit, perfected the electroscope in 1908-09, up to a sensitivity of 1 volt

By sea and by sky







the winning idea !!!: immersing an electroscope 3m deep in the sea Pacini finds a significant (20% at 4.3σ) reduction

The radiation comes from the sky!!!

Viktor Hess, 1912 Several flights up to 6 Km Day/night, eclipses, ...



"Cosmic Rays" ! (Milikan)





The confirmation



Kolhörster





First shower energy estimates !

Particle Physics Birth

- Positron (Anderson1932) Antimatter! (Dirac) $\gamma \rightarrow e^+e^-$ (Einstein)
- μ (Anderson1937)
 - Rossi, 1940:
 - Muon life time.
 - Time dilation!
- π (Latter, Powell1947) Strong interactions (Yukawa)
- K, Λ, ... (Leprince Ringuet 1944, Rochester, Butter 1947, ...)
 - Strangeness



The present data

Charged cosmic rays (p/nucleus)



John Linsley, 1962



Forward Region Extreme Energies

Small fluxes indirect measurements

Anthropomorphic representation



Air shower energy spectrum (E >10¹⁴ eV)



The highest energy region (E >10¹⁸ eV)



Origin (E ~<10¹⁵ eV)?

Energy density (cosmic rays)

$$\rho_{E} \sim \int E \frac{dN}{dE} dE \sim 10^{-12} \text{ erg/cm}^{3}$$

$$P \sim \frac{\rho_E V_{galaxy}}{\tau_{esc}} \sim 5.10^{40} erg/s$$

For example, power dissipated by a Supernova (the remnant of a collapsed star)

 $P \sim 10^{42} \text{ erg/s}$

Supernovae in our Galaxy

1 SN ~30 Years



Where can be these accelerators in the Universe?

Large Hadron Collider











 $\mathbf{E} \propto \mathbf{B} \mathbf{R}$

 $R \sim 10 \text{ km}, B \sim 10 \text{ T}$ $E \sim 10 \text{ TeV}$

A few new terms

- Stellar endproducts. A star heavier than the Sun collapses at the end of its life into a neutron star (R ~ few km, which can be pulsating – a pulsar) or into a BH, and ejects material in an explosion (SuperNova Remnant).
 - Very large B fields are in the pulsar; magnetic fields also in the SNR
- The centres of galaxies host black holes, often supermassive (millions or even billion solar masses). They might accrete at the expense of the surrounfing matter, and accelerate particles in the process. When they are active, they are called Active Galactic Nuclei.





They surpass human-made accelerators





High Luminosity Sophisticated detectors Central region

Energy limited

Energy spectrum $(10^9 < E < 10^{15} eV)$



Cosmic-Ray Composition

| Ζ | Element | F | Z | Element | F |
|---------|------------|------|---------|---------|------|
| 1 | Н | 540 | 13-14 | Al-Si | 0.19 |
| 2 | ${\rm He}$ | 26 | 15 - 16 | P-S | 0.03 |
| 3 - 5 | Li-B | 0.40 | 17 - 18 | Cl-Ar | 0.01 |
| 6-8 | C-O | 2.20 | 19 - 20 | K-Ca | 0.02 |
| 9–10 | F-Ne | 0.30 | 21 - 25 | Sc-Mn | 0.05 |
| 11 - 12 | Na-Mg | 0.22 | 26 - 28 | Fe-Ni | 0.12 |

Table 26.1: Relative abundances F of cosmic-ray nuclei at 10.6 GeV/nucleon normalized to oxygen ($\equiv 1$) [6]. The oxygen flux at kinetic energy of 10.6 GeV/nucleon is 3.26×10^{-2} (m² s sr GeV/nucleon)⁻¹. Abundances of hydrogen and helium are from Refs. [2,3]. Note that one can not use these values to extend the cosmic ray flux to high energy because the power law indicies for each element may differ slightly.

Solar vs Galactic Cosmic Rays



ACE News 83 (2004) - NASA

Electron/positron



Figure 26.2: Differential spectrum of electrons plus positrons multiplied by E^3 (data from [17–23]). The line shows the proton spectrum multiplied by 0.01. The inset shows the positron to electron ratio measured by PAMELA [24] compared to the expected decrease [25].

UHECR Correlation with AGNs



Vernon-Cetty-Vernon AGN catalog

 $E\!>57$ EeV, z<0.018, distance <3.1 deg.

 $P = 0.006, f = 33 \pm 5\%$



Cen A



Closest (4.6 Mpc) powerful radio galaxy with characteristics jets and lobes, candidate for UHECR acceleration. Auger South.





Cosmic magnetic fields & CR propagation (CR for astronomy?)

- Gyroradius for a proton
 - Typical B field in the Galaxy: ~ 1-3 μG
 - Intergalactic B field largely unknown
 1 nG < B < 1 fG
- If you want to look at the GC (d ~ 8 kpc) you need E > 2 10¹⁹ eV
 - But only 1 particle / km2 / year
 - And: no galactic emitters expected at this energy
- But in principle one could look outside the galaxy, were B is smaller and there are SMBHs...

$$\frac{r}{1 \text{ pc}} \cong \frac{\frac{E}{1 \text{ PeV}}}{\frac{B}{1 \mu G}}$$

extremely difficult

=> Locating the sources of charged CR is impossible AT GROUND

- We need to detect a large statistics of protons with E between 10¹⁹ eV and 10²¹ eV
- Areas ~ 30 thousand km² (1/10 of Italy), ~impossible to find
- Needs satellites
- VHE photons
- Neutrinos
 - Detected only from Sun and SN1987A
- Gravitons
 - Undetected up to now; we shall not discuss them

Cosmic Microwave Background:

2.7 K

A Galactic gas cloud called Rho Ophiuchi: 60 K Dim star near the center of the Orion Nebula: 600 K

> the Sun: 6000 K

Cluster of very bright stars, Omega Centauri, as observed from the space: 60,000 K



Accretion Disks can reach temperatures >> 10⁵ K



But this is still ~1 keV, in the X-Ray band!



Cosmic γ rays: two different production mechanisms expected to be at work



Leptonic vs. Hadronic models for γ emission

- SSC: currently explain most emissions
 - easy to accelerate electrons to TeV energies
 - easy to produce synchrotron and IC gamma-rays

But:

- recent results would require more sophisticated leptonic models
- Hadronic Models:
 - protons interacting with ambient hadronic targets -> neutrinos (1)
 - But needs adequate targets: works well for SNR, more difficult for AGN
 - protons interacting with photons (hadronic photoproduction) -> neutrinos (2)
 - proton synchrotron (no neutrinos)
 - very large magnetic fields

High Energy γ rays: non-thermal Universe

- Particles accelerated in extreme environments interact with medium
 - Gas and dust; Radiation fields Radio, IR, Optical, …;
 Intergalactic Magnetic Fields, …
- Gamma rays traveling to us!
 - HE: 30 MeV to 30 GeV
 - VHE: 30 GeV to 30 TeV



No deflection from magnetic fields, gammas point ~ to the sources

Gamma rays can trace cosmic rays at energies ~10x

- Large mean free path
 - Regions otherwise opaque can be transparent to X/γ

Studying Gamma Rays allows us to see these aspects of the Universe
VHE sources have been located using gammas

- Thanks mostly to Cherenkov telescopes, imaging of VHE (> 30 GeV) galactic sources and discovery of many new galactic and extragalactic sources: ~ 150 (and >200 papers) in the last 7 years
 - And also a better knowledge of the diffuse gammas and electrons
- A comparable success in HE (the Fermi realm); a 10x increase in the number of sources
- A new tool for cosmic-ray physics and fundamental physics
- In the Galaxy, mostly associated to stellar endproducts (PWN in particular). Extragalactic emission associated to AGN





Cen A

- Located at a distance of ~13 Mly, the best candidate as a CR emitter from Auger
- Very faint VHE gamma-ray emitter (HESS 2009)







Neutrinos







Solar radiation: 98 % light 2 % neutrinos At Earth 66 billion neutrinos/cm² s

Hans Bethe (1906–2005, Nobel prize 1967)

Thermonuclear reaction chains (1938)

Neutrinos

- Neutrinos play a very special role in particle astrophysics. Their cross section is small, and they can leave production sites without interacting => can carry information about conditions deep in the core of astrophysical objects that produce them
- They are produced in the nuclear reaction chains by which stars generate energy. Each conversion of 4p into He produces two neutrinos - the Sun emits ~2x10³⁸ neutrinos per second
- They are also produced in nature's most violent explosions, including the Big Bang
- They are produced as secondary byproducts of cosmic-ray collisions also in the atmosphere. Detection of these neutrinos can help constraining properties of the primary cosmic ray spectrum, more effectively than photons. Neutrinos produced by reactions of ultra-high-energy cosmic rays can provide information on otherwise inaccessible cosmic accelerators

CR: clear synergy gammas-neutrinos

- Smoking gun would be a correlation of VHE CR (difficult: magnetic fields) or a neutrino signal
- Many model uncertainties present in the gamma/p relation disappear when studying gamma/neutrino

$$\frac{dN_{\nu}}{dE} \approx \frac{1}{2} \frac{dN_{\gamma}}{dE}$$

(reflects the fact that pions decay into gamma rays and neutrinos that carry 1/2 and 1/4 of the energy of the parent. This assumes that the four leptons in the charged pion decay equally share the charged pion's energy)



neutrino sky



v energy spectrum



IceCube 40 strings

northern sky: 14139 neutrinos

operated 375.5 days



search for

southern sky: 23151 muons

- clustering
- high energy (>> 100 TeV)

? beams

- New neutral particles ?
 - Neutralino, axion, glueballino, enhanced neutrino-nucleon cross-section ...
- Decay products of super-heavy X particles ~ GUT scale Gauge bosons, Higgs bosons, super-heavy fermions ... cosmic strings, magnetic monopoles...
 - Several "anomalous" events
 - Almost no photons
 - highly penetrating hadrons

Are they real ??





The sources

Where to find these accelerators in the Universe?

- Stellar endproducts
- AGN

(one has energy from the gravitational potential energy, and magnetic fields from plasma motions)

Stellar endproducts (supernova remnants)

A massive star begins its lifetime burning the H in its core under conditions of equilibrium. When the H is exhausted, the core contracts until $3\alpha \rightarrow 12C$ can take place. He is then burned to exhaustion. This pattern (fuel exhaustion, contraction, heating, and ignition of the ashes of the previous cycle) might repeat several times depending on the mass, leading finally to an explosion. A 25-solar-mass star would go through the set of burning cycles ending up in the burning of Si to Fe in about 7 My, with the final explosion taking ~days





Evolution of Stars depends on mass

| M < 0.08 M _{sun} | Never ignites hydrogen → cools ("hydrogen white dwarf") | | Brown dwarf |
|-------------------------------------|---|-------------------------------|---|
| $0.08 < M \lesssim 0.8 M_{sun}$ | Hydrogen burning not completed in Hubble time | | Low-mass main-squence star |
| $0.8 \lesssim M \lesssim 2 M_{sun}$ | Degenerate helium core after hydrogen exhaustion | | • Carbon-oxygen white dwarf |
| $2 \leq M \leq 5-8 M_{sun}$ | Helium ignition non-degenerate | | • Planetary nebula |
| 8 M _{sun} ≲ M < ??? | All burning cycles → Onion skin structure with degenerate iron core | Core collapse supernova | Neutron star (often pulsar) Sometimes black hole? Supernova remnant (SNR), e.g. crab nebula |

Rate of SN in our galaxy: ~ 1 every 30 years (CR considerations); ~1/century (astrophysical cons.)

Observed: ~1/400 years

SNR (continued)

If mass is large enough to explode/collapse, the supernova remnant (SNR) is the structure left over after a supernova explosion: a high density n star (or a BH) lies at the center of the exploded star, whereas the ejecta appear as an expanding bubble of hot gas that shocks and sweeps up the interstellar medium.



SNR (continued)

- Typical velocities for the expulsion of the material are of the order of c/100
- The shock slows down over time as it sweeps up the ambient medium, but it can expand over tens of thousands of years and over tens of pc before its speed falls below the local sound speed
 - Typical duration depends on the mass, and can go between 10ky and 100 ky
- Based on their emission and morphology (which are actually related, and related to age and mass), SNRs are generally classified under three categories: shell-type, pulsar-wind nebulae (PWN), and composite

SNR (continued)

- Shell supernova remnants. As the shockwave from the supernova explosion plows through space, it heats and stirs up any interstellar material it encounters, which produces a big shell of hot material in space. Magnetic fields strengths are estimated to be from 10 µG to 1 mG.
- Pulsar-wind nebulae (PWN) are SNR where a young pulsar is slowing down in rotation: the rate of decrease of rotational energy is in the range 10³² to 10³⁹ erg/s. A small part of this energy goes into the pulsed electromagnetic radiation observed from the pulsar, most of it goes into an outflowing relativistic magnetized particle wind: due to external pressure, this wind abruptly decelerates at a termination shock, beyond which the pulsar wind radiates synchrotron emission so resulting in a PWN. The best know case is the Crab Nebula.
- Composite are the SNR that cannot be classified in any of the above





AGN

- Supermassive black holes (SMBH) of 1M-10G solar masses reside in the cores of most galaxies (the centre of our Galaxy hosts a BH ~4 million solar masses, its mass having been determined by the orbital motion of nearly stars.
- Their fueling by infalling matter can produce a spectacular activity; when they are active, i.e., they are powered by accretion at the expenses of nearby matter, these BH are called Active Galactic Nuclei (AGN).





- It is believed that a "unified model" accounts for all AGN. The SMBH and its inner accretion disk is surrounded by a dusty torus of matter, and the type of active galaxy that is seen depends on the orientation of the torus and jets relative to the observer's line of sight.
- The jet radiates strongly along its axis, also due to the Lorentz boost. The view of an observer, who is looking very close to this axis, will be dominated by the jet emission, and a possibly variable source with no spectral lines will be seen: this is called a BL Lac object, or blazar
- Looking at a modest angle to the jet, an observer will see an unobscured compact source inside the torus; this is in general called a quasar
- From a viewpoint closer to the plane of the torus, the central engine is hidden, and one observes the jets and lobes (extended radio-emitting clouds) of a radio galaxy
- Typical values of B are of the order of 10 G to 10000 G.
- The jet can be traced down to a scale size of order 0.01 pc which is < 100 000 times its total length and <100 times the radius of the BH

AGN



Slanted story. Seemingly diverse astronomical objects may be different views of galactic cores.

- For short period of times, ~ seconds
 - GRB have a bimodal distribution of duration; accordingly, they are classified as "long" or "short"
 - a source emits more energy in Gamma ray than the rest of the Universe
- Long GRBs are thought to come from a hypernova (an old star of large mass), or a core-collapse supernova. There is a prevailing consensus that the basic mechanism of GRB emission is an expanding relativistic fireball, with the beamed radiation due to shock waves. No direct confirmation that they are CR emitters
- During the abrupt compression the magnetic field could be squeezed to extremely large values, of the order of 10¹²G to 10¹⁴G, in a radius of some tens of kilometers

Gamma-Ray Bursts



END OF 1ST LECTURE

Acceleration mechanism

Fermi 2nd order (1949)

particles accelerated in stochastic collisions with massive interstellar clouds (collisions to a moving diffusive wall!)

In the cloud reference frame

 $E_1^* = \gamma E_1(1 - \beta \cos \theta_1)$ $E_2^* = E_1^*$

Back to the Lab reference frame

$$E_2 = \gamma E_2^* (1 + \beta \cos \theta_2^*)$$

Then:

$$\frac{\Delta E}{E} = \frac{1 - \beta \cos \theta_1 + \beta \cos \theta_2^* - \beta^2 \cos \theta_1 \cos \theta_2^*}{1 - \beta^2} - 1$$

But:

$$\langle \cos \theta_2^* \rangle = 0$$

$$\langle \cos \theta_1 \rangle = \frac{\int_{-1}^1 \cos \theta_1 (1 - \beta \cos \theta_1) d\cos \theta_1}{\int_{-1}^1 (1 - \beta \cos \theta_1) d\cos \theta_1} = -\frac{\beta}{3}$$



 $\beta \sim 10^{-1}$

Acceleration mechanism Fermi 1st order

Shock formation :

- Sudden release of Energy (CMEs, SNRs, GRBs,...)
- Supersonic flow hits an obstacle (AGNs jets, pulsar winds, ...)

Particles gain energy by consecutive crossings of the shock front!

$$\frac{\Delta E}{E} = \frac{1 - \beta \cos \theta_1 + \beta \cos \theta_2^* - \beta^2 \cos \theta_1 \cos \theta_2^*}{1 - \beta^2} - 1$$

Now (plane shock front):

$$\langle \cos \theta_1 \rangle = \frac{\int_{-1}^0 \cos^2 \theta_1 \, d\cos \theta_1}{\int_{-1}^0 \cos \theta_1 \, d\cos \theta_1} = -\frac{2}{3}$$

$$\langle \cos \theta_2^* \rangle = \frac{\int_0^1 \cos^2 \theta_2^* \, \mathrm{d} \cos \theta_2^*}{\int_0^1 \cos \theta_2^* \, \mathrm{d} \cos \theta_2^*} = \frac{2}{3}$$

Solar coronal mass ejection 9 Mar 2000





 $\left< \frac{\Delta E}{E} \right> \simeq \frac{4}{3} \beta$

The power law

In each cycle the particle gains a small fraction of energy ϵ . After n cycles:

 $\mathbf{E}_{\mathbf{n}} = \mathbf{E}_{\mathbf{0}} \ (1 + \varepsilon)^{\mathbf{n}}$

Or the number of cycles to attain an energy E is:

 $n = \ln(E/E_0)/\ln(1+\varepsilon)$



The particle may escape from the shock region with some probability P_i . Then the probability to escape with $E > E_n$ is:

$$P_{E_n} = P_i \sum_{j=n}^{\infty} (1 - P_i)^n = (1 - P_i)^n$$

and

$$\frac{N}{N_0} = P_{E_n} = \left(\frac{E}{E_0}\right)^{-\alpha} \qquad \alpha = -\frac{\ln(1-P_i)}{\ln(1+\varepsilon)} \cong \frac{P_i}{\varepsilon} \quad \left(\frac{dN}{dE}\right)_{Source} \approx E^{-2}$$
$$\frac{dN}{dE} \propto \left(\frac{E}{E_0}\right)^{-\gamma} \qquad \left(\frac{dN}{dE}\right)_{Earth} \propto \left(\frac{dN}{dE}\right)_{Source} \cdot \tau_{esc}(E) \propto E^{-2.7}$$

10³ Proof of the origin of CR up to 100 almost the knee

- Evidence that SNR are sources of CR up to ~1000 TeV (almost the knee) came from morphology studies of RX J1713-3946 (H.E.S.S. 2004) with photons
- Striking evidence from the morphology of SNR IC443 (MAGIC + Fermi/Agile 2010)

0.3

02

0.1

0.0

-0.1 В

-0.3

0.4

3.0 [deg]



1 m⁻² s⁻¹

Molecular clouds close to IC 443, W51, RX J1713.7-3946

- VHE $\gamma\text{-ray}$ excess compatible with cloud

• Differential energy spectrum prefers π^0 production





The propagation



Propagation

Charged cosmic rays diffuse and interact in the Galactic randomly magnetized ISM . Confinement times are quite long $(\tau \sim 10^7$ years) and directions become basically isotropic.

Transport equation :



by: T.Gaisser

PISM~ Cm

halo

disk

Galactic

300pc

-15 kpc -

Leaky-Box

A box where charged cosmic rays freely propagates having however some probability to escape by the walls. The sources are uniformely distributed.



Simplified stationary equation :



Constraining propagation models

Secondary/primary ratios

Unstable/stable isotopos



Box size, Diffusion coef., escape time, ...

Radioactive clocks

Energy dependence (τ_{esc})

Primary/Primary Secondary/Primary B/C Ratio C/O Ratio 1 1 1 1 1 1 1 1 $D(E) \propto E^{\delta}$ 1/3-----10⁻¹ $\delta = 0.7$ 1 1 1 1 11 11 10⁻¹ 10⁻² 10² 10^{3} 10^{3} 10² 10 10 Energy(GeV/n) Energy(GeV/n)

CREAM 2008

Confinement and composition

Magnetic Field:

- Galactic $\sim 1-3 \mu G$ -Intergalatic 1 nG > B > 1 fG

Larmor Radius:

 $R = \frac{E}{ZeB}$

Several knees?



Galactic Magnetic Field deflection (p)



T.Stanev

Above 10¹⁹ : Astronomy !

An unique opportunity to measure the galactic magnetic field ?

The Greisen-Zatsepin-Kuzmin (GZK) cutoff



≈ 6 Mpc

GZK is model dependent



Interaction with the Solar System

and the Earth

We see the Moon!

- Moon diameter 0.5°
- Angular resolution: $<1^\circ$





Argo-YBJ



Ice Cube
Interaction in the atmosphere/Earth





p/nucleus



Fluxes in the atmosphere



Figure 26.3: Vertical fluxes of cosmic rays in the atmosphere with E > 1 GeV estimated from the nucleon flux of Eq. (26.2). The points show measurements of negative muons with $E_{\mu} > 1$ GeV [32–36]. J. Beringer et al (PDG) PR D86 010001 (2012)

Hajo Drescher, Frankfurt U.

time = -900 µs

Hajo Drescher, Frankfurt U.

time = -800 µs

Hajo Drescher, Frankfurt U.

time = -700 µs

Hajo Drescher, Frankfurt U.

time = -600 µs

Hajo Drescher, Frankfurt U.

time = -500 µs

Hajo Drescher, Frankfurt U.

time = -400 µs

The events: first interaction

Hajo Drescher, Frankfurt U.

time = -300 µs

The events: shower development

Hajo Drescher, Frankfurt U.

time = -200 µs

The events: shower development

Hajo Drescher, Frankfurt U.

time = -100 µs

The events: shower hits Earth surface

 $P(Fe) Air \rightarrow Baryons (leading, net-baryon \neq 0)$

Hajo Drescher, Frankfurt U.

Extensive Air Showers (EAS)



EAS longitudinal profiles

Gaisser



EAS transverse profiles

NKG (Nishimura, Kamata, Greisen)

$$\rho(\mathbf{r}) = c(s)N_{e} / r_{0}^{2}(\mathbf{r} / r_{0})^{s-2}(1 + r / r_{0})^{s-4.5}$$

J. Knapp et al. | Astroparticle Physics 19 (2003) 77-99





P_t distributions Multiple Coulomb scattering

Shower development

 $P(Fe) Air \rightarrow Baryons (leading, net-baryon \neq 0)$

$$\rightarrow \pi^{0} \qquad (\pi^{0} \rightarrow \gamma\gamma \rightarrow e^{+}e^{-}e^{+}e^{-} \rightarrow ...)$$

$$\rightarrow \pi^{\pm} \qquad (\pi^{\pm} \rightarrow \nu \mu^{\pm} \text{ if } L_{\text{decay}} < L_{\text{int}})$$

$$\rightarrow K^{\pm}, D. ...$$

- e.m. and weak interations
 - well known !

hadronic interations

- large uncertainties !
- forward region, small p_t , very high \sqrt{s}
- main parameters: σ_{in} , k_{in} , $\langle n \rangle$, (fraction π^0 , Nb of Baryons, ...)

Nuclear fragmentation

- Nuclei are not just a superposition of nucleons !

Missing Energy

- 5% to 10% ...



Muon production



Assumptions:

- cascade stops at $E_{\text{part}} = E_{\text{dec}}$
- each hadron produces one muon

Primary particle proton

 π^0 decay immediately

 π^{\pm} initiate new cascades

$$N_{\mu} = \left(\frac{E_0}{E_{\text{dec}}}\right)^{\alpha}$$
$$\alpha = \frac{\ln n_{\text{ch}}}{\ln n_{\text{tot}}} \approx 0.82\dots 0.95$$

(Matthews, Astropart. Phys. 22, 2005)

Muon spectrum at Earth Surface



Figure 26.4: Spectrum of muons at $\theta = 0^{\circ}$ (\blacklozenge [41], \blacksquare [46], \checkmark [47], \blacktriangle [48], \times , + [43], \circ [44], and \bullet [45] and $\theta = 75^{\circ} \diamond$ [49]). The line plots the result from Eq. (26.4) for vertical showers. J. Beringer et al (PDG) PR D86 010001 (2012)

Under Earth

At the Earth surface : ~70 m⁻² s⁻¹ sr⁻¹ ~1 cm⁻² min⁻¹ (Ω ~ π)



Propagation of γ-rays



For gamma rays, relevant background component is optical/infrared (EBL)
different models for EBL: minimum density given by cosmology/star formation

 e^+



The γ "GZK" ($\gamma \gamma \rightarrow e^+ e^-$)





2011-11-20 - Up-to-date plot available at http://www.mpp.mpg.de/~rwagner/sources/



Y rates sandre De Alogelis

Selection bias?



ν Propagation and the ν "GZK"



Neutrinos propagate unabsorbed and without deflection throughtout the universe up to $E_v \sim 10^{25} \text{ eV}$

QED HE processes: bremsstrahlung for electrons...

- (and pair production for photons).
 Forbidden in vacuo by 4-momentum conservation
 - Require interaction with the medium
- Bremsstrahlung (braking radiation): photons of momentum q<E_e emitted with probability ~proportional to 1/q
 - (and collimated: ~ m_e/E)

ie, energy emission is ~constant for each interval of photon energy; total is propto E

 The dependence on the material appears through the radiation length Xo:

 $dE_e/dx = -1/Xo$

- Xo can be found in tables. It is ~300 m for air at NTP, ~43 cm for water; for density 1 g/cm³ roughly proportional to A/Z²
- Collision energy loss is almost constant (plateau)





QED HE processes: ...pair production for photons

• Pair production: $\lambda = (9/7)$ Xo for $E_{\gamma} >> 2m_e$

Energy spectrum ~ flat







Electromagnetic showers

- When a high-energy e or γ enters an absorber, it initiates an em cascade as pair production and bremsstrahlung generate more e and γ with lower energy
- The ionization loss becomes dominant < a critical energy ${\rm E_c}$
 - E_c ~ 84 MeV in air, ~73 MeV in water; ~ (550/Z)MeV
 - Approximate scaling in y = E/E_c
 - The longitudinal development ~scales as the radiation length in the material: t = x/Xo
 - The transverse development scales approximately with the Moliere radius $R_{\rm M}$ ~ (21 MeV/E_c) Xo
 - In average, only 10% of energy outside a cylinder w/ radius R_M
 - In air, $R_M \sim 80$ m; in water $R_M \sim 9$ cm
- Electrons/positrons lose energy by ionization during the cascade process
- Not a simple sequence: needs Monte Carlo calculations

An analytic model: Rossi's "approximation B"

 $t_{max} + 1.4$

٧

- Rossi in 1941 published an analytical formulation for the shower development as a set of 2 integro-differential equations under the approximation that:
 - Electrons lose energy by ionization & bremsstrahlung; asymptotic formulae hold
 - Photons undergo pair production only; asymptotic formulae hold ($E > 2 m_e$)
- Very good approximation until $E \sim E_c$



y

Peak of shower, tmax Centre of gravity, tmed Number e^+ and e^- at peak Total track length T

A snapshot of Rossi's equations (you can solve them, but...)

$$\frac{\partial \pi(E,t)}{\partial t} = 2 \int_0^1 \gamma \left(\frac{E}{u},t\right) \psi_0(u) \frac{du}{u} - \int_0^1 \left[\pi(E,t) - \frac{1}{1-v} \pi \left(\frac{E}{1-v},t\right)\right] \varphi_0(v) dv + \epsilon \frac{\partial \pi(E,t)}{\partial E}$$

$$\frac{\partial \gamma(W, t)}{\partial t} = \int_0^1 \pi \left(\frac{W}{v}, t\right) \varphi_0(v) \frac{dv}{v} - \sigma_0 \gamma(W, t)$$



A simplified approach (Heitler)

- Qualitative features may be
 obtained from a simple model
 - Each electron with E>E_C travels 1 Xo and then gives up half of its energy to a bremsstrahlung photon
 - Each photon with E>E_c travels
 1 Xo and then creates an e+epair with each particle taking E/2
 - 3. Electrons with E<E_C cease to radiate and lose the rest of their energy by collisions
 - Ionization losses are negligible for E>E_C



Results from the simplified approach

 If the initial electron has energy E₀>>E_C, after t Xo the shower will contain 2^t particles. ~equal numbers of e+, e-, γ, each with an average energy

$$\mathsf{E}(\mathsf{t}) = \mathsf{E}_0/2^{\mathsf{t}}$$

 The multiplication process will cease when E(t)=E₂

$$t_{max} = t(E_C) \equiv \frac{\ln \left(E_0 / E_C \right)}{\ln 2},$$

and the number of particles at this point will be

$$N_{max} = \exp\left(t_{max} \ln 2\right) = E_0 / E_C$$



Energy measurement

 Errors asymptotically dominated by statistical fluctuations:

$$\frac{\sigma_E}{E} \cong \frac{k_E}{\sqrt{E}} \oplus c$$



Hadronic showers

- Although hadronic showers are qualitatively similar to em, shower development is more complex because many different processes contribute
 - Larger fluctuations
- Some of the contributions to the total absorption may not give rise to an observable signal in the detector
 - Examples: nuclear excitation and leakage of secondary muons and neutrinos
- Depending on the proportion of π^0 s produced in the early stages of the cascade, the shower may develop predominantly as an electromagnetic one because of the decay $\pi^0 \rightarrow \gamma \gamma$
- The scale of the shower is determined by the nuclear absorption length λ_{H}
 - Typically $\lambda_H > Xo$

Hard and soft hadronic interactions



J. Knapp

String fragmentation

Think of the gluons being exchanged as a spring... which if stretched too far, will snap! Stored energy in spring \rightarrow mass !

Hadrons!



In this way, you can see that quarks are always confined inside hadrons (that's **CONFINEMENT**) !



"Standard" Hadronic models (low pt)

Minimum string configuration

Ralph Engel


Hadronic model parameters

p air



Cross sections extrapolations ...





Cosmic Rays and LHC



Cosmic Rays and LHC



Small-x region (LHC as a pathfinder for CR, and vice-versa)



LHC Kinematics (14 TeV)



Most of the energy flows into very forward region

LHCf: forward photon production



Re-tuning of models needed, size of effect still unclear, distributions for neutrons needed

Cosmic Rays and LHC: total cross section



- Test Glauber model
- Tune EAS simulations

