Instrumentation for Flavor Physics - Lesson II



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Outline

Lesson II

Tracking detector

Particle identification (PID) detectors

Muon detectors

Tracking detectors

Stable charged particles: e,μ,π,K,p

Trajectory = path of a charged particle (directional 1-D object in 3-space)

Tracking Detector = highly-segmented detector capable of sensing a charged particle without greatly changing its **P**

Hit = Position measurement in a single active element of a tracking detector

Track = Collection of hits presumably from one particle, reconstructed to estimate the particle **P** and trajectory

LHCb spectrometer

Silicon Tracker: 500 μ m thick, single sided Si strip detector, pitch~200 μ m, vertical and stereo angle strips arrangement (x-u-v-x)=(0°,-5°,+5°,0°)



Magnetic spectrometer resolution



Magnetic spectrometer resolution

Multiple scattering contribution does not depend upon momentum



Drift Chamber - BaBar DCH



$$\vec{F} = q\vec{v} \times \vec{B}$$

momentum \Rightarrow geometry (curvature)
 $\omega = 1/R = qB/p_t$

Magnet+Tracking: magnetic spectrometer

DCH is inside a solenoidal magnetic field B=1.5T

Momentum and angles measurements for charged tracks

40 layers of hexagonal cells providing spatial and ionization loss measurements

Helium-based gas mixture, lowmass wires

DCH components



Aluminum endplates Inner (Be+Al) cylinder Outer (Carbon Fiber) cylinder 5 m³ gas volume (Helium:Isobutane 80:20) wires ~32 KNewton tension

DCH wires





Four stages:

Ionization

Gas amplification

Electronic amplification + digitization

Calibration and reconstruction

Ionization

Charged particles can ionize gas atoms Helium provides most ionization 10x less absolute ionization than Argon Wi=41eV for He (energy for electron/ion pair) Isobutane absorbs photons (quencher) prevents photoelectric electrons from cathode keeps ionization local to particle trajectories



Drifting charges due to electric field Anode = wire

Gas amplification

Electrons and ions drift due to ${\cal E}$

 $v_{\sim} E/m$, $v_{e} \sim 10^{3} v_{ion} (v_{e} \sim 3 v_{h} in Si)$

E~V/r, electrons accelerate near the wire

When $E_e > E_{ionization}$ they ionize the gas, make more electrons (very close to the anode wire)

 \Rightarrow chain reaction

DCH Gain ~ 5×10^4

proportional mode: proportional to initial ionization produced in the detector



Signal amplification and timing

Electrons collected on wire in ~1ns

Electronic signal (voltage) appears on wires as ions drifts to field wires

 $\leq 1 \mu s$ to see total charge

Signal shaping, amplification + digitization

FEE chip: 8 bit ADC (dE/dx), 1ns resolution Time-to-digital converter (TDC) (tracking)

Signal time depend on ionization location

Interpret TDC as "distance to the wire" with time to distance calibration

Drift cell isochrones (100 ns steps)



DCH Calibration

Time to distance calibration using control samples

Position resolution vs drift distance



Tracking and momentum resolution



Distance to wire provide a 2D (x-y) constraint on the track position

Stereo views provide some z information



Exercises

- 1. Calculate the sagitta, s, for a muon (p=100 GeV/c) transversing a region (1m) with \perp magnetic field B=1.5T. Show that $\sigma_p/p = \sigma_s/s$.
- 2. What limits the time resolution of a proportional counter? Can you give an estimate of the maximum uncertainty in terms of wire spacing, d=1cm, and drift velocity, v_d=5cm/µs?
- 3. In a cylindrical proportional counter, the dominant signal is induced by the motions of electron or ions? Can you estimate the ratio of the two signal for a simple case with wire radius=10µm and cell radius 1cm?
- 4. If you wish to measure the momentum of a 10 GeV/c singly-charged particle to 1% accuracy, in a 2 T field, using a 1 m long magnet, how well do you have to know the exit angle (see figure in pag. 7)? Calculate the adequate spacing between detectors, to achieve this precision, if you use a gas detector with 200µm hit resolution and a silicon detector with 20µm hit resolution.

PID detectors

For charged particles, the momentum measurement is provided by the tracking system ρ $\sqrt{1}$

$$p = \beta \gamma m_0 c$$
 $\gamma = \sqrt{\frac{1}{1 - \beta^2}}$

By measuring the velocity of the particle, β , we can determine its mass $m_0 = \frac{p}{\beta \gamma c}$

Measurement of velocity

Cherenkov angle
$$\cos \theta_c = \frac{1}{\beta n}$$
Time of flight $\tau \propto \frac{1}{\beta}$ Energy loss $\frac{dE}{dx} \propto \frac{z^2}{\beta^2} \ln (a\beta\gamma)$

I will not cover transition radiation detectors and calorimeters

A charged particle in a medium that moves faster than the speed of light in the same medium generates Cherenkov radiation.



Energy loss by Cherenkov radiation very small compared to ionization (<1%) Number of photons per unit length and per unit energy interval of the photon:

$$\frac{d^2 N}{dEdx} = \frac{\alpha z^2}{\hbar c} \sin^2 \theta_c = \frac{\alpha^2 z^2}{r_e m_e c^2} \left(1 - \frac{1}{\beta^2 n^2(E)} \right)$$
$$\approx 370 \, \sin^2 \theta_c(E) \, \text{eV}^{-1} \text{cm}^{-1} \qquad (z=1)$$

Equivalently:

$$\frac{d^2 N}{d\lambda dx} = \frac{2\pi\alpha z^2}{\lambda^2} \left(1 - \frac{1}{\beta^2 n^2(\lambda)} \right) = \frac{2\pi\alpha z^2}{\lambda^2} \sin^2 \theta_C$$
Integrate over sensitivity range:
$$\frac{d^2 N}{dx} = \int_{350 \text{ nm}}^{550 \text{ nm}} d\lambda \frac{d^2 N}{d\lambda dx}$$

$$= 475 \ z^2 \sin^2 \theta_C \text{ photons/cm}$$

Cherenkov radiators

Minimal velocity for Cherenkov emission $\beta_{thr} = 1/n$

 $\begin{array}{ll} \mbox{charged particles with} & \beta < 1/n & \mbox{give no light} \\ & \mbox{with} & \beta > 1/n & \mbox{give light} \end{array}$

Parameters of Typical Radiator

Medium	n	$m{eta}_{thr}$	θ _{max} [β=1]	N _{ph} [eV ⁻¹ cm ⁻¹]
Air	1.000283	0.9997	1.36	0.208
Isobutan	1.00127	0.9987	2.89	0.941
Water	1.33	0.752	41.2	160.8
Quartz	1.46	0.685	46.7	196.4

Gas radiators for high momentum particles O(10) GeV Solid radiators for low momentum particles O(1) GeV

Threshold Cherenkov counter

Useful to identify particles in a beam line (with fixed momentum). For example a 50 GeV π ⁺ beam with proton contamination

By choosing a medium with a suitable refractive index, it can be arranged that the π will produce light, but the protons will not

Ex: calculate the refractive index n for the above experiment.



Ring Imaging Cherenkov detector (RICH)

Proposed by Seguinot and Ypsilantis [NIM 142, 377 (1977)

The Cherenkov cone can be imaged into a ring using a spherical mirror

Requires large area photosensitive detectors (e.g. PMT arrays)



The Cherenkov angle θ_c is determined by measuring the radius of the ring r

First RICH detectors at collider experiments

DELPHI and SLD: RICH with 2 radiators to cover a large momentum range for π/K/p separation 0.7-45 GeV/c



RICH in LHCb

RICH-1 and RICH-2 for different momentum range: [2-40] GeV/c, [15-100] GeV/c, respectively

Use a "thin" tilted focusing mirror to bring the images outside detector acceptance

Use different radiators: Aerogel (n=1.03) for low momentum p<10 GeV/c, C_4F_{10} gas (n=1.0014) intermediate momentum, CF₄ gas (n=1.0005) in RICH-2 for high momentum p>20 GeV/c

RICH-1cross section



Cherenkov angle reconstruction

Simulated event in RICH-1 Large (small) rings: Aerogel (C₄F₁₀) Multiple tracks generate overlapping rings→ pattern recognition





Reconstructed Cherenkov angle vs track momentum in C_4F_{10} radiator

RICH performance at LHCb

Cherenkov angle versus particle momentum for the RICH radiators

Kaon efficiency and pion misidentification rate vs momentum



Detector of Internally Reflected Cherenkov light (DIRC)



Ring imaging Cherenkov detector based on total internal reflection.

Preserving angle information due to precision surfaces

Long rectangular bars made from synthetic fused silica ("quartz") are used as both radiator and light guide.

Rings projected in water-filled standoff box (best match to quartz index), where photons are detected with an array of 10K PMTs

DIRC



Light internally reflected in DIRC bar



Typically, a charged track well above the Cherenkov threshold produces 20-60 photons

These are detected by an array of closely packed photomultipliers (PM's) in the Standoff Box

Very good precision in the measurement of the time (a few ns), required to reduce background



 n_2

 n_1

Reconstruction

Display of an $e^+e^- \rightarrow \mu^+\mu^-$ event reconstructed in BABAR with two different time cuts



Reduce background photons outside a fiducial time window around the trigger time Signal patterns are conic sections rather than rings due to different geometry

DIRC performance



Excellent π/K separation over the kinematically allowed region at BaBar

Time Of Flight (TOF) detector

Measure signal time difference between two detectors with good time resolution [can use time of beam crossing as start signal]

Relativistic particles, $E \simeq pc \gg m_i c^2$:

$$\Delta t \approx \frac{L}{pc^2} \left[\left(pc + \frac{m_1^2 c^4}{2pc} \right) - \left(pc + \frac{m_2^2 c^4}{2pc} \right) \right]$$
$$\Delta t = \frac{Lc}{2p^2} \left(m_1^2 - m_2^2 \right)$$



Number of standard deviation separation for a TOF and RICH detector

$$N_{\sigma} = \frac{|m_1^2 - m_2^2| d}{2 p^2 \sigma_t c} \quad \text{TOF} \quad N_{\sigma} = \frac{|m_1^2 - m_2^2|}{2 p^2 \sigma_{\theta} \sqrt{n^2 - 1}} \quad \text{RICH}$$

Factor $1/\sqrt{n^2 - 1}$ allows RICH to cover high momentum spectrum using appropriate radiators (n value close to 1)

TOF performance - example L = 2 m



dE/dx measurement - an example

Measured energy loss in BABAR drift chamber (DCH)



Bethe-Bloch distribution

Large Landau fluctuations: use many dE/dx measurements and truncate large energy-loss values "truncated mean"

dE/dx resolution~7.5%

Excellent kaon-pion separation below 700 MeV/c, where the DIRC is not effective

Separation at to the 2 σ level at higher momentum (3 GeV/c)

Muon detectors

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Muon system - LHCb

Provide information for L0 trigger and offline muon identification



Minimum ionizing particles; penetrates thick iron absorbers (80cm); measure signal behind complete detector

5 stations M1-M5 (for a total active surface of 435 m²) separated by iron filters (iron is only permeable to energetic muons)

1368 Multi-Wire Proportional Chambers and 24 triple GEM (Gas electron Multiplier) detectors (in the M1 region close to the beam) with spatial X-Y readout

Performance - LHCb

$B_s \rightarrow \mu^+ \mu^-$ rare decay a candidate event



A track is a candidate muon when has hits in a minimum number of stations (2-4 depending on momentum value)

Acceptance ~20% for muons from inclusive b semileptonic decays

Muon p>6GeV/c to cross 5 stations (~20 X₀)

Muon ID efficiency ~97% with 1-3% $\pi \rightarrow \mu$ misidentification

Instrumented Flux Return (IFR)- BaBar

Original design: for muon and K_{L} detection

Instrumented Flux Return



Performance - BaBar



Muons with p>1 GeV/c are likely to cross the entire IFR while hadrons are stopped within the first layers

Decays in flight ~2% $\pi \rightarrow \mu$ misidentification

Serious degrading of RPC efficiencies during time: quality of chambers and operations under high temperatures

Barrel replaced with Limited Streamer Tubes (LST) and Forward Endcap with new RPC of better quality



- 1. A Cherenkov counter based on nitrogen gas is located on a charged particle beam with momentum 20 GeV/c. The dependence of the refractive index on the pressure P is $n=1.3 \cdot 10^{-9}P$ (Pa). Calculate the range of P for which the Cherenkov emits light for π but not for K.
- 2. What are the Cherenkov angles for electrons and pions of 1 GeV/c for a radiator with n = 1.4? What will be the ratio of the number of radiated photons for incident electrons and pions?
- 3. What is the fraction of charged pions decaying in flight within L=4m if p=3GeV/c? If L=20m and p=40GeV/c?