

Neutrinos in Cosmology



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8th IDPASC school
VLC, 21-31 May 2018



Suggested References Books

Modern Cosmology, S. Dodelson (Academic Press, 2003)

Kinetic theory in the expanding Universe, Bernstein (Cambridge U., 1988)

Neutrino Cosmology, Mangano, Miele, Lesgourgues & SP (Cambridge U., 2013)

Reviews

Massive neutrinos and cosmology, J. Lesgourgues & SP,
Phys. Rep. 429 (2006) 307-379 [[astro-ph/0603494](#)]

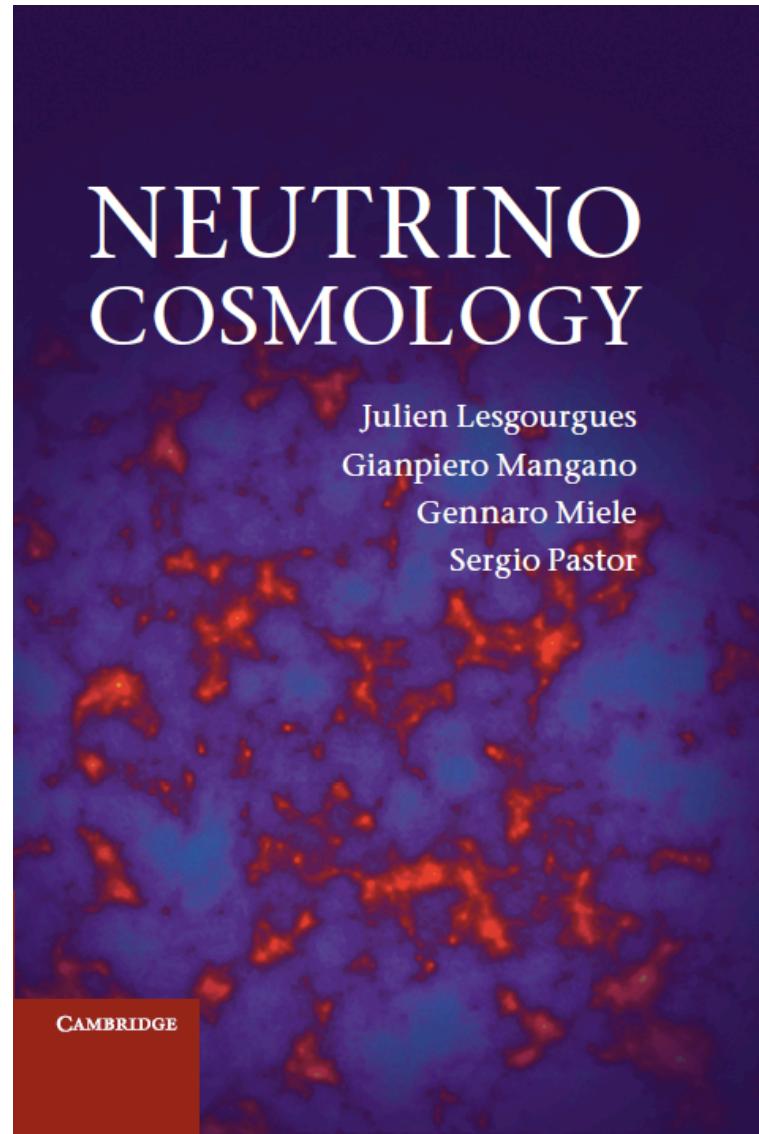
Primordial Nucleosynthesis: from precision cosmology to fundamental physics,
F. Iocco, G. Mangano, G. Miele, O. Pisanti & P.D. Serpico
Phys. Rep. 472 (2009) 1-76 [[arXiv:0809.0631](#)]

Neutrino mass in cosmology: status and prospects, Y.Y.Y. Wong
Ann. Rev. Nucl. Part. Sci. 61 (2011) 69-98 [[arXiv:1111.1436](#)]

Neutrino Physics from the CMB and LSS, K.N. Abazajian, M. Kaplinghat
Ann. Rev. Nucl. Part. Sci. 66 (2016) 401-420

Cosmological and astrophysical constraints on neutrino properties, M. Lattanzi &
M. Gerbino, *Frontiers in Physics* 5 (2018) 70 [[arXiv:1712.07109](#)]

For more details...

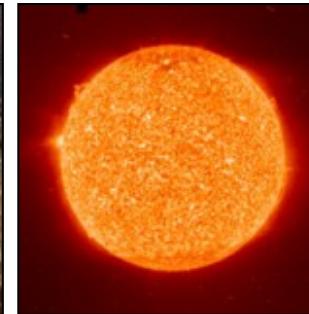
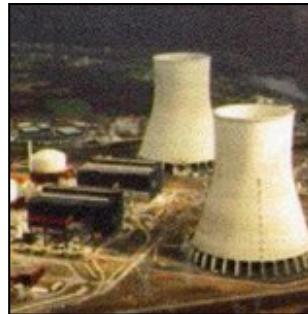


Ed. Cambridge Univ. Press, 2013

Where do neutrinos come from?



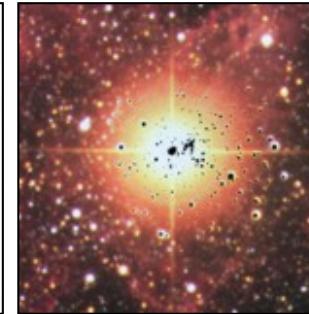
Nuclear reactors



Sun



Particle accelerators

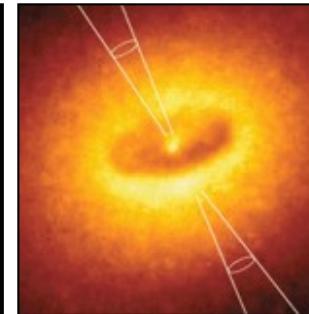
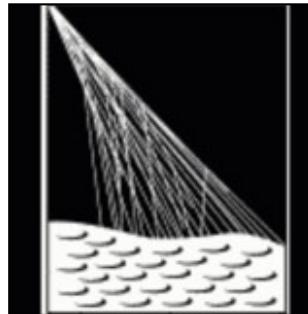


Supernovae

SN 1987A ✓



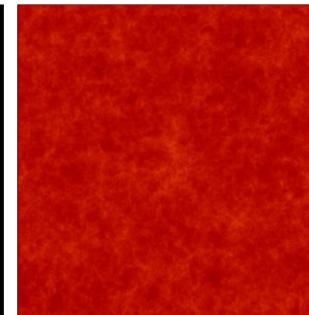
Earth Atmosphere
(Cosmic rays)



Accelerators in
astrophysical sources ?✓

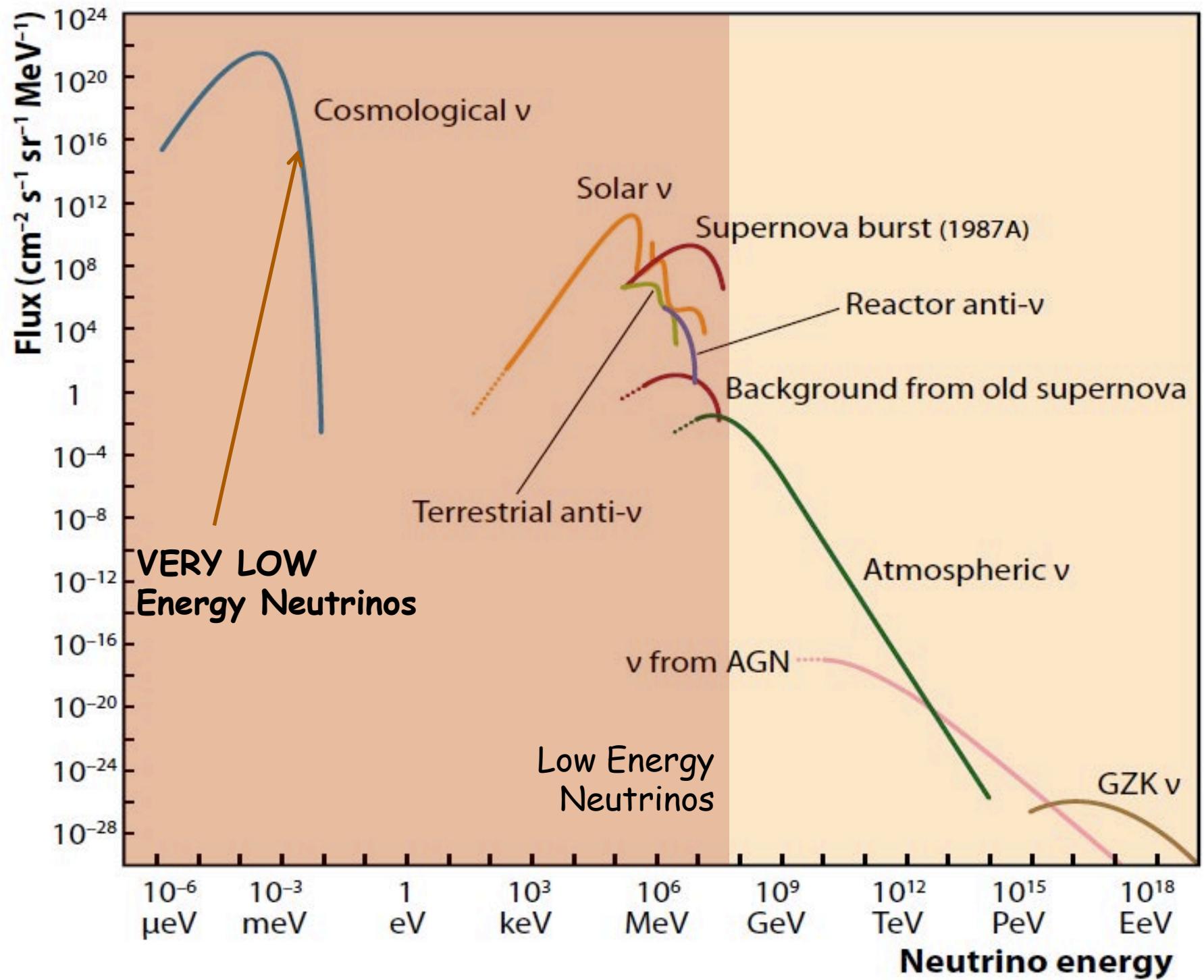


Earth interior
(Natural Radioactivity)

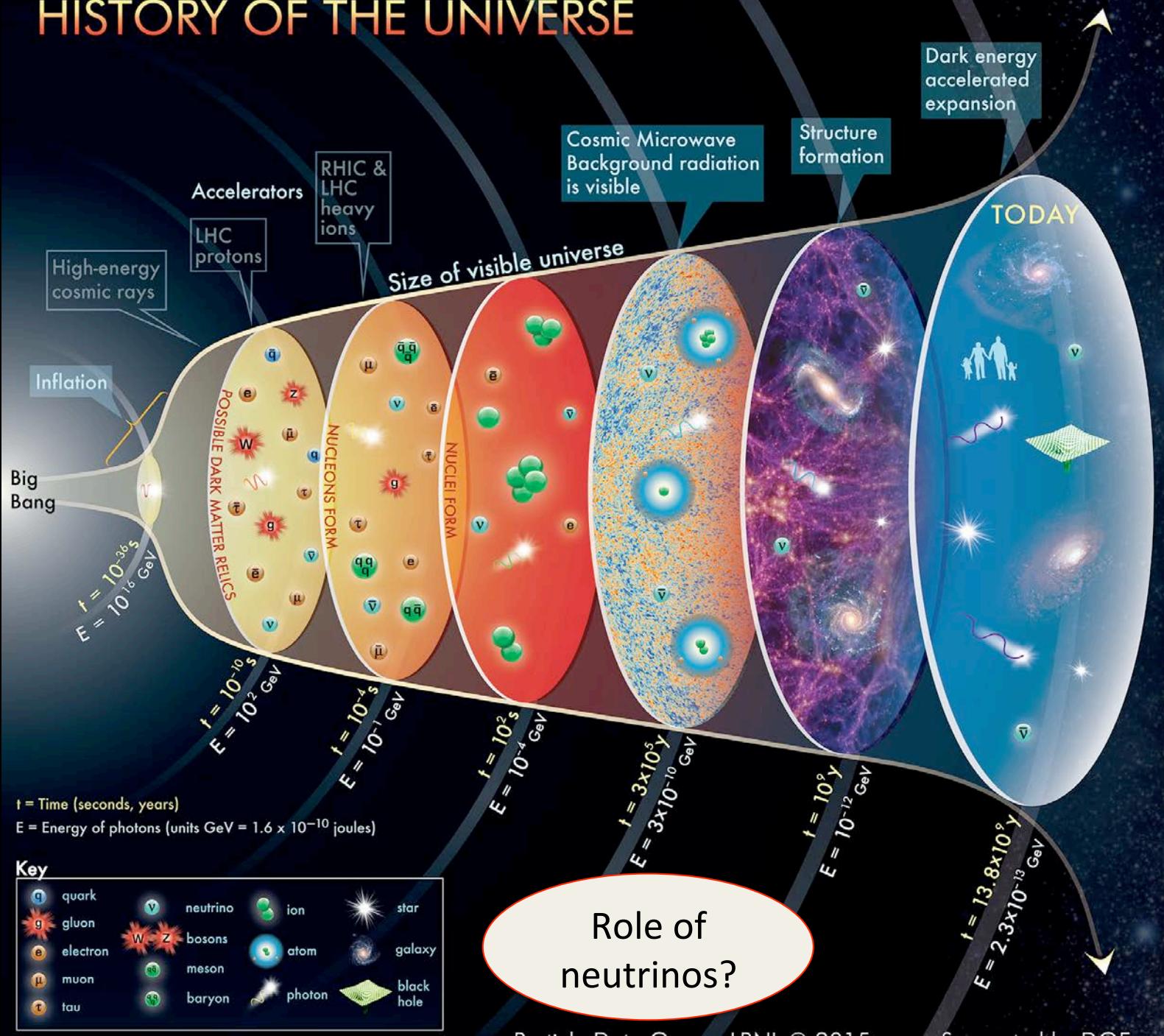


Early Universe
(today 336 v/cm^3)

Indirect evidence



HISTORY OF THE UNIVERSE



The concept for the above figure originated in a 1986 paper by Michael Turner.

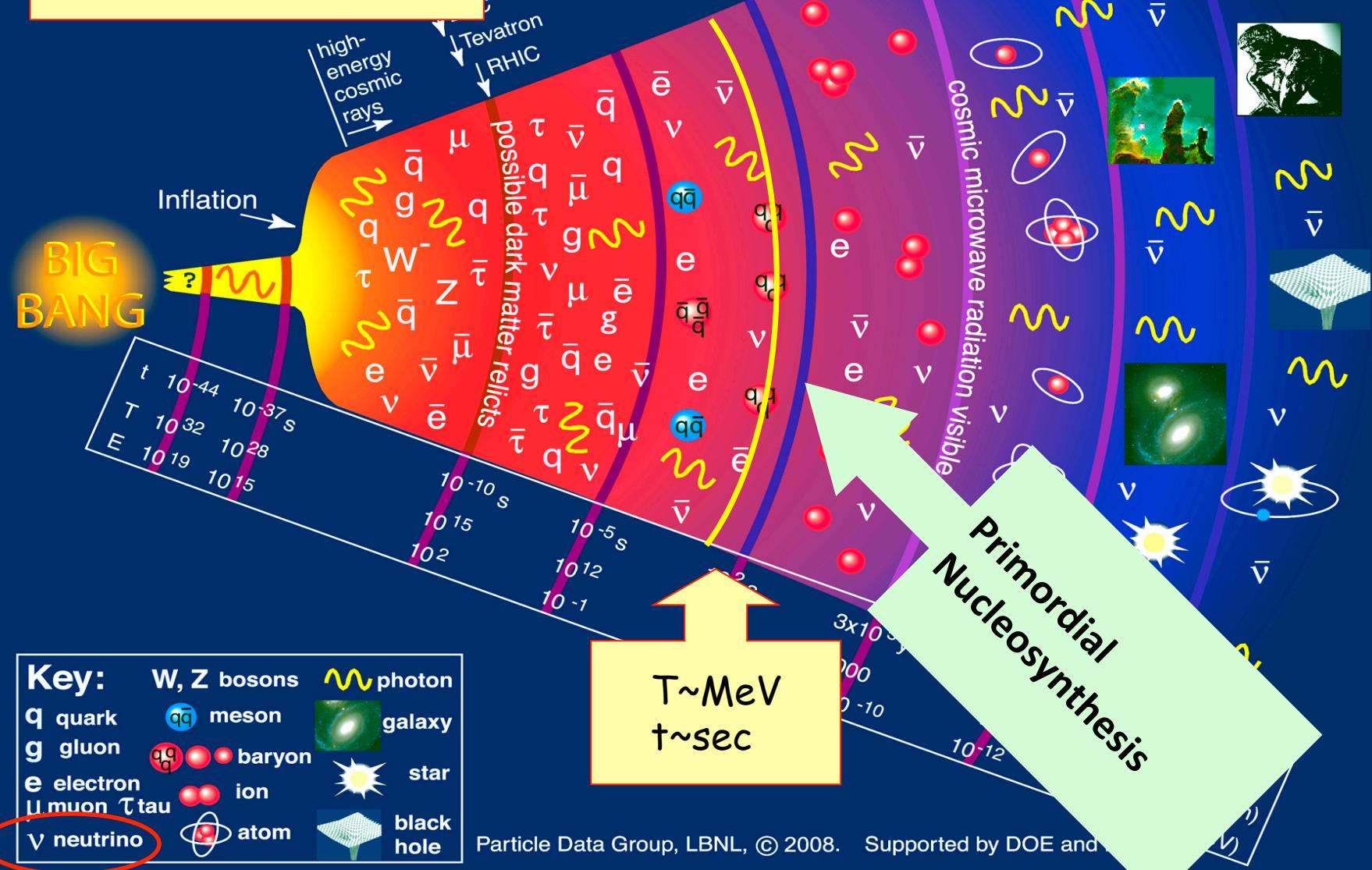
Particle Data Group, LBNL © 2015

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Introduction: neutrinos and the history of the Universe

History of the Universe

Neutrinos coupled by weak interactions

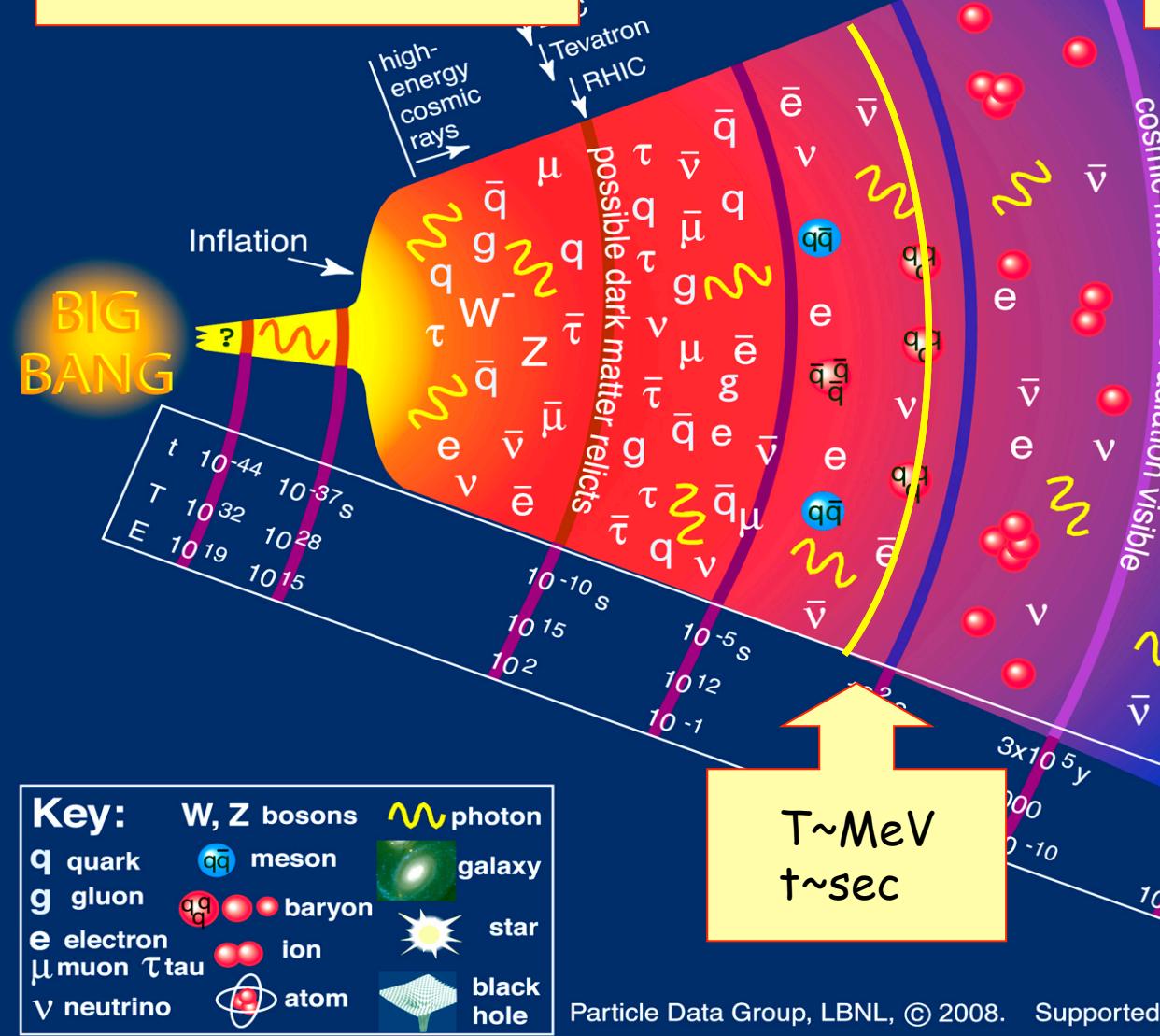


Particle Data Group, LBNL, © 2008. Supported by DOE and

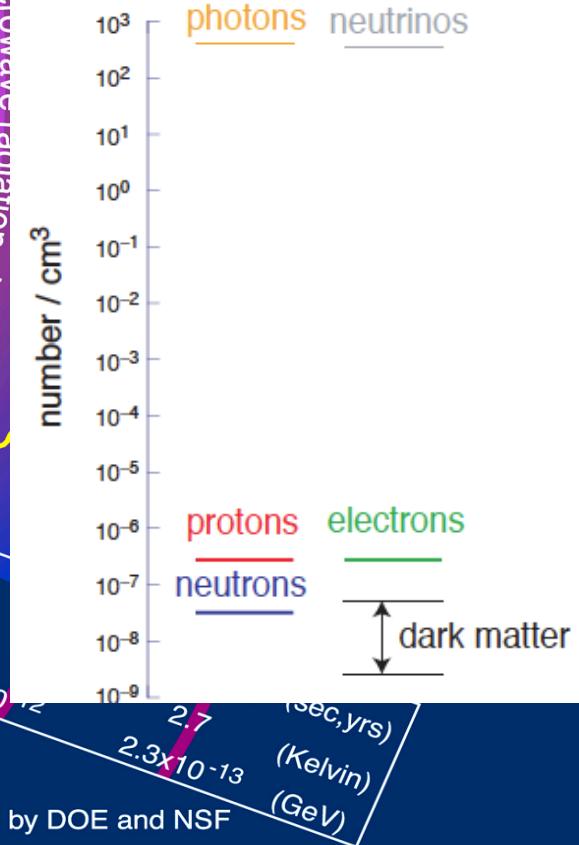
History of the Universe

Neutrinos coupled by weak interactions

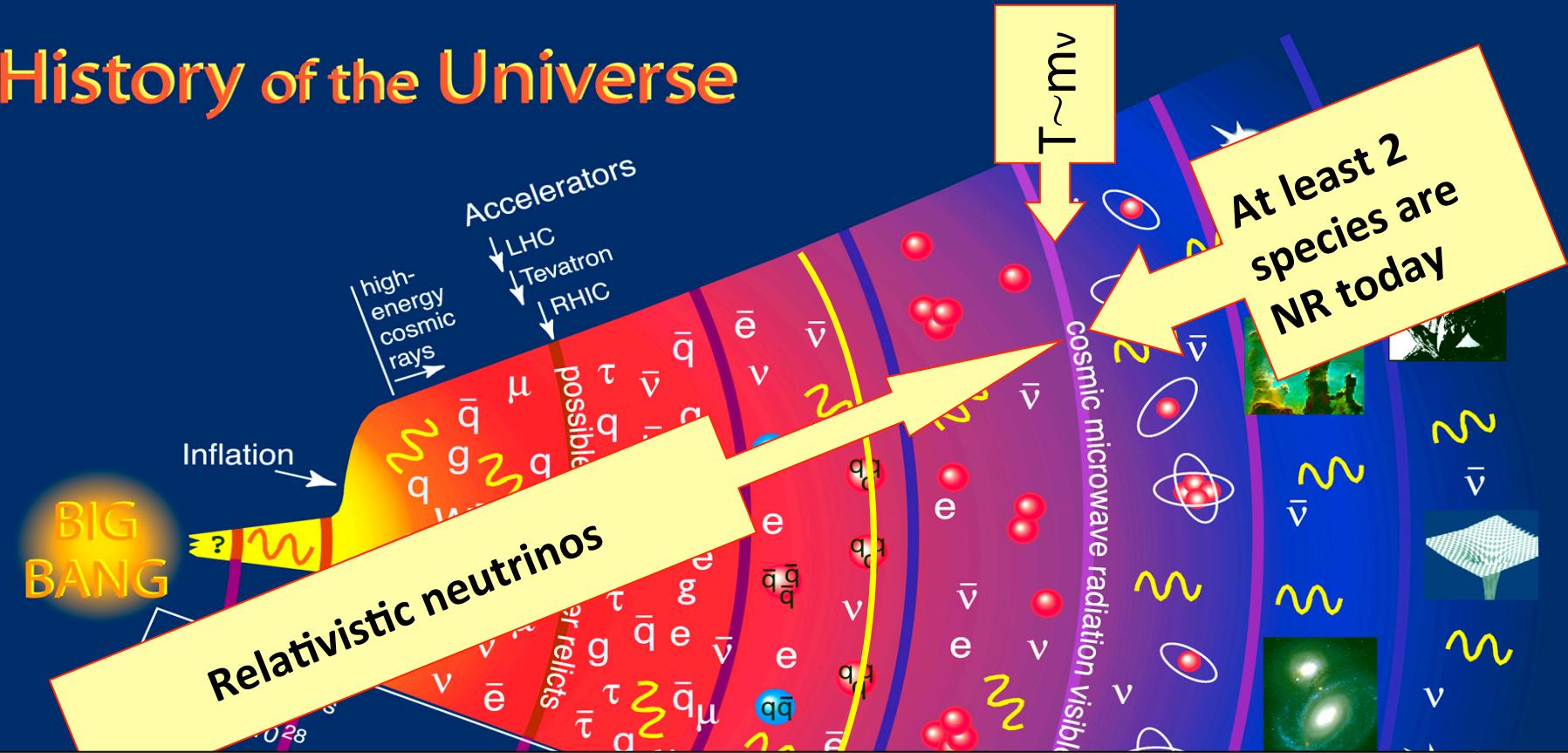
Decoupled neutrinos
(Cosmic Neutrino
Background or CNB)



The Particle Universe



History of the Universe

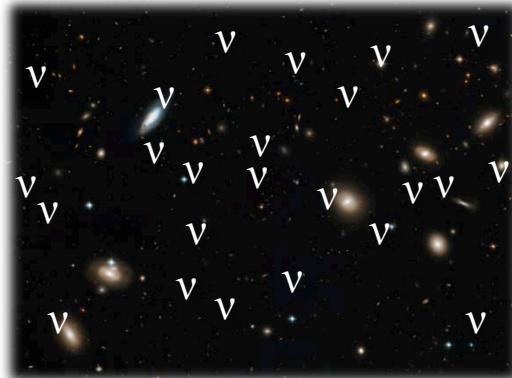


Neutrino cosmology is interesting because Relic neutrinos are very abundant:

- The CNB contributes to **radiation at early times** and to matter at late times (info on the number of neutrinos and their masses)
- Cosmological observables can be used to **test standard or non-standard neutrino properties**



Outline

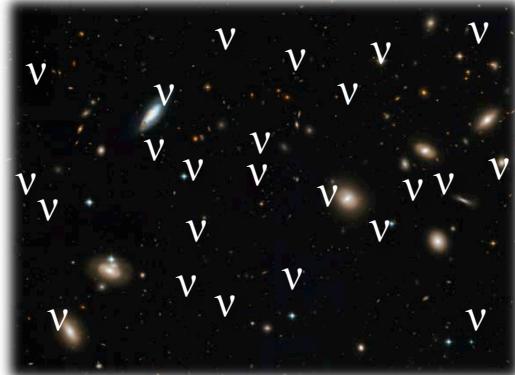


Introduction: neutrinos and the history of the Universe

Basics of Cosmology

Production and decoupling of relic neutrinos

Outline

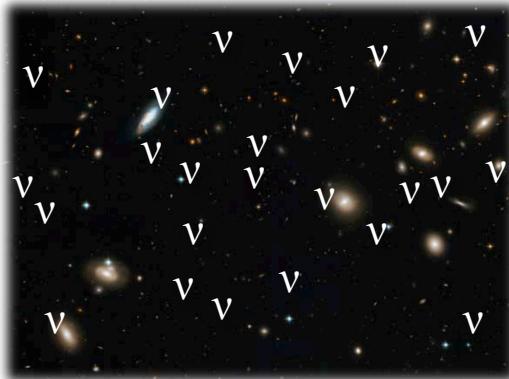


**The radiation content
of the Universe (N_{eff})**

**Neutrinos and Primordial
Nucleosynthesis**

**Neutrino oscillations
in the Early Universe**

Outline



Massive neutrinos as Dark Matter

**Effects of neutrino masses
on cosmological observables**

**Present bounds on neutrino
properties from cosmology**

**Future sensitivities on neutrino
physics from cosmology**

Basics of Cosmology

Eqs in the SM of Cosmology

The FLRW Model describes the evolution of the isotropic and homogeneous expanding Universe

$$ds^2 = g_{\mu\nu}dx^\mu dx^\nu = dt^2 - a(t)^2 \left(\frac{dr^2}{1-kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right)$$

$a(t)$ is the scale factor and $k=-1,0,+1$ the curvature

Einstein eqs



$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi G T_{\mu\nu} - \Lambda g_{\mu\nu}$$

Energy-momentum
tensor of a perfect
fluid



$$T_{\mu\nu} = (p + \rho)u_\mu u_\nu - p g_{\mu\nu}$$

Eqs in the SM of Cosmology

00 component
(Friedmann eq)



$$H(t)^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2}$$

$$\rho = \rho_M + \rho_R + \rho_\Lambda$$

$H(t)$ is the Hubble parameter

$$\frac{k}{H(t)^2 a^2} = \Omega - 1 \quad \Omega = \rho / \rho_{\text{crit}}$$

$$\dot{\rho} = \frac{d\rho}{dt} = -3H(\rho + p)$$

$\rho_{\text{crit}} = 3H^2/8\pi G$ is the critical density

Eq of state $p = \alpha \rho$



$$\rho = \text{const } a^{-3(1+\alpha)}$$

Radiation $\alpha = 1/3$

$$\rho_R \sim 1/a^4$$

Matter $\alpha = 0$

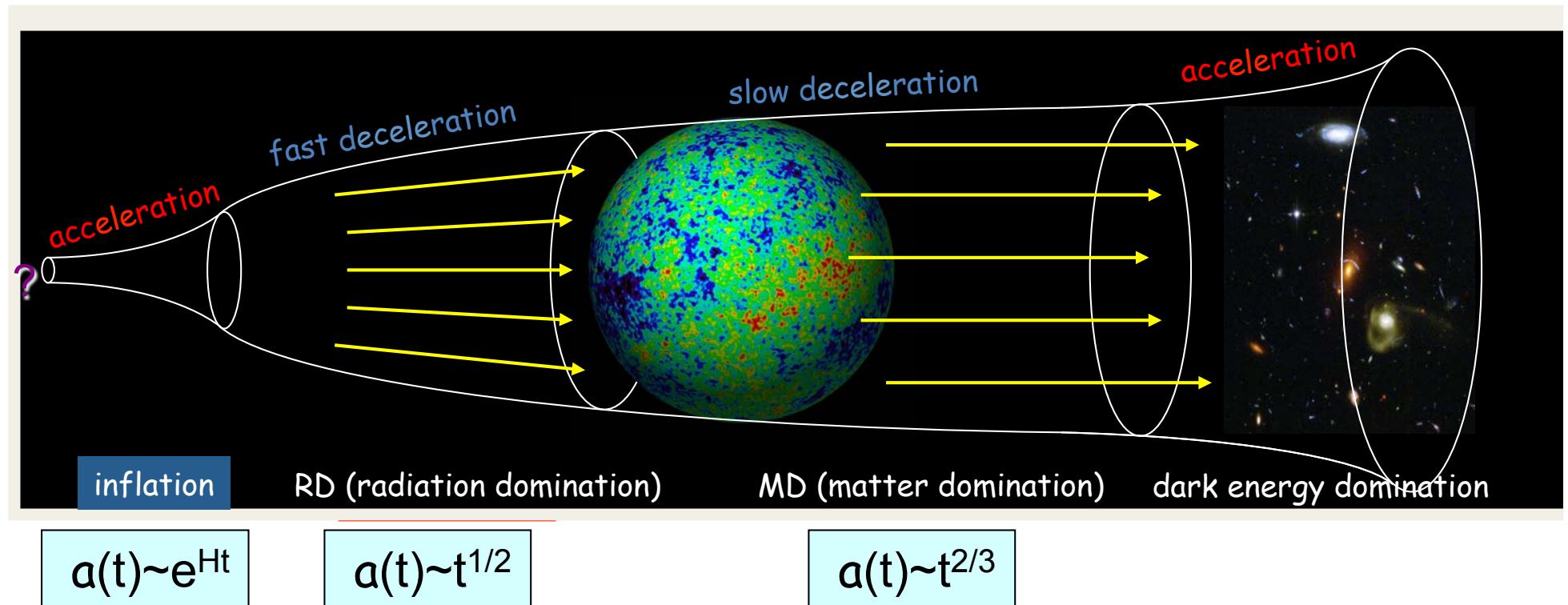
$$\rho_M \sim 1/a^3$$

Cosmological constant $\alpha = -1$

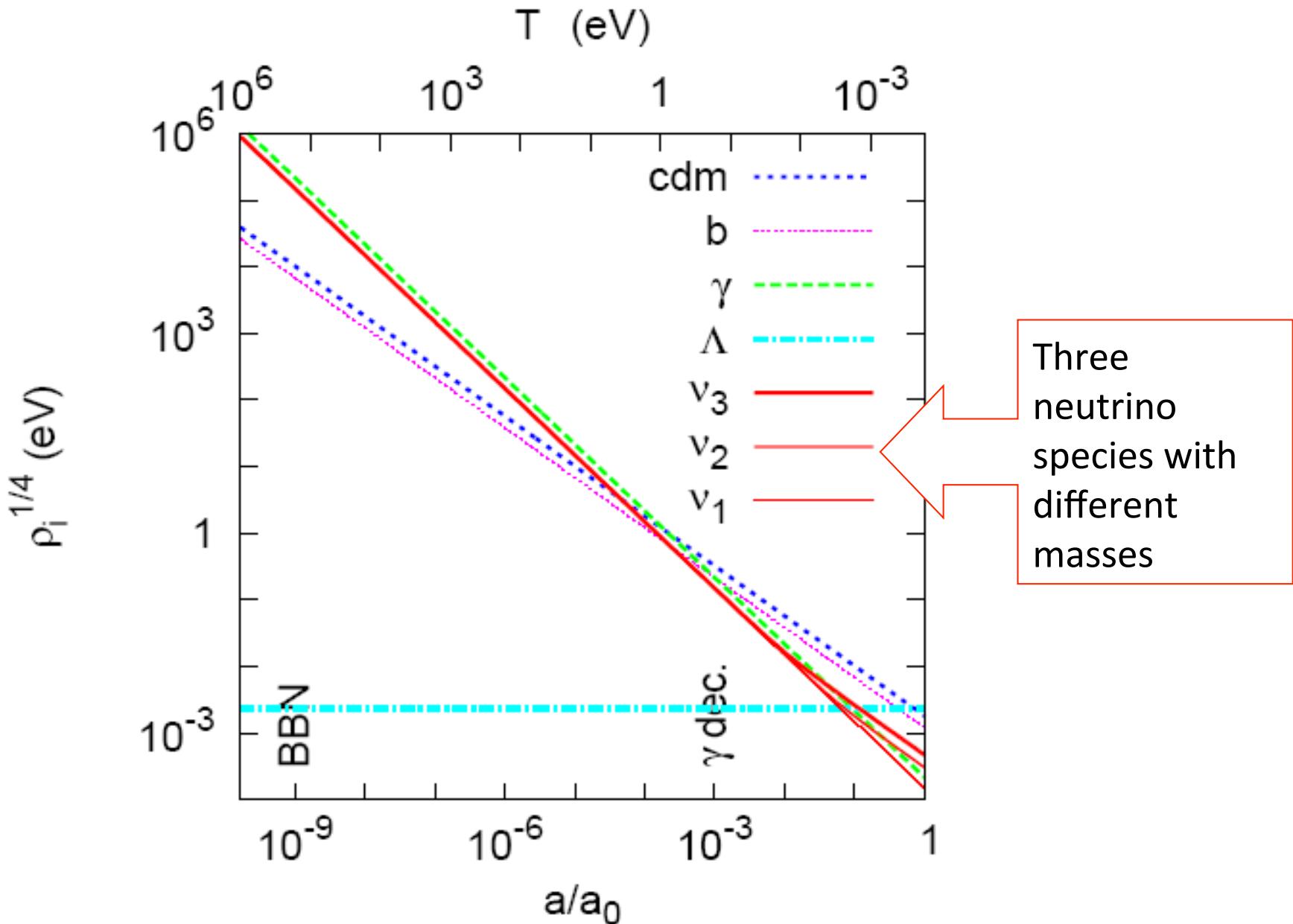
$$\rho_\Lambda \sim \text{const}$$

Evolution of the Universe

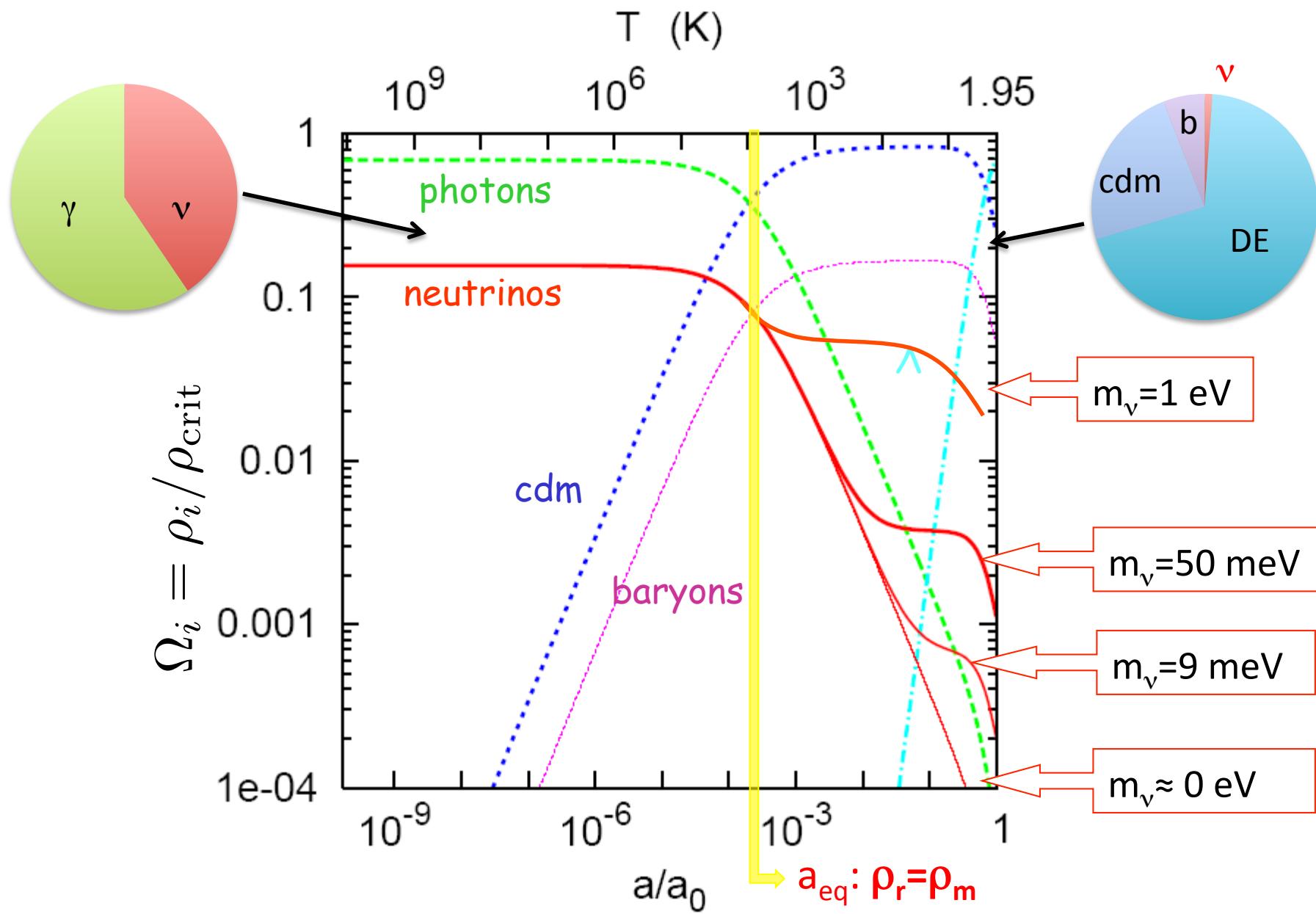
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3}$$



Evolution of the background densities: 1 MeV → now



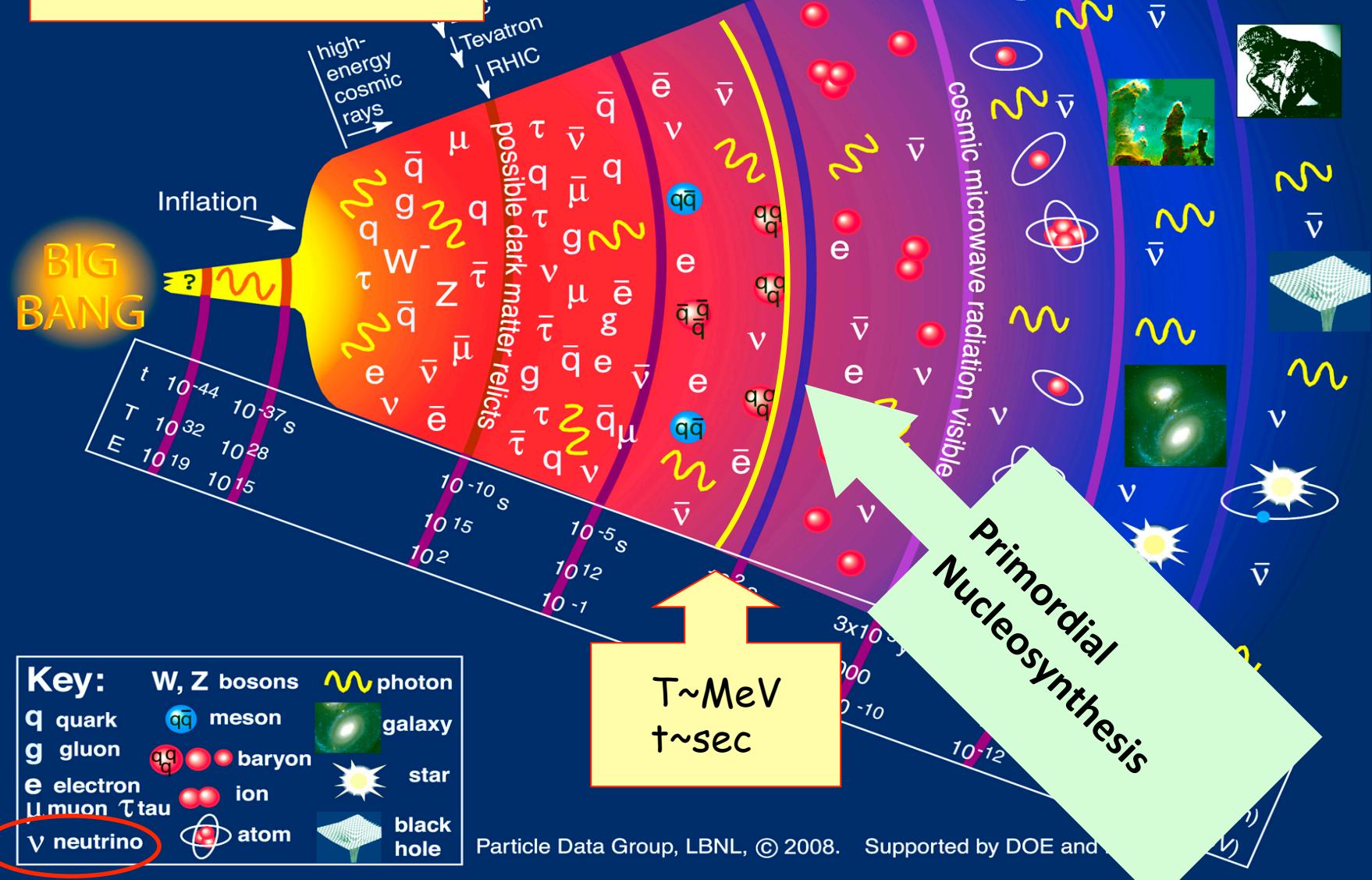
Background densities: 1 MeV → now



Production and decoupling of relic neutrinos

History of the Universe

Neutrinos coupled by weak interactions



Equilibrium thermodynamics

Particles in equilibrium when T are high and interactions effective

$$T \sim 1/a(t)$$

Distribution function of particle momenta in equilibrium

$$f_i^{eq}(p, T) = \left[\exp \left(\frac{E_i - \mu_i}{T} \right) \mp 1 \right]^{-1}$$

Thermodynamical variables

VARIABLE	RELATIVISTIC		NON REL.
	BOSE	FERMI	
n	$\frac{\zeta(3)}{\pi^2} g T^3$	$\frac{3\zeta(3)}{4\pi^2} g T^3$	$g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T}$
ρ	$\frac{\pi^2}{30} g T^4$	$\frac{7\pi^2}{8 \cdot 30} g T^4$	mn
p		$\frac{\rho}{3}$	$nT \ll \rho$
$\langle E \rangle$	$2,701T$	$3,151T$	$m + \frac{3}{2}T$

$$n = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} f_i(p, T) \quad \rho = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} E_i f_i(p, T)$$

$$p = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} \frac{p^2}{3E_i} f_i(p, T) \quad \langle E \rangle = \rho/n$$

Neutrinos in Equilibrium

$$1 \text{ MeV} \lesssim T \lesssim m_\mu$$

$$T_\nu = T_{e^\pm} = T_\gamma$$

$$\begin{aligned}\nu_\alpha \nu_\beta &\leftrightarrow \nu_\alpha \nu_\beta \\ \nu_\alpha \bar{\nu}_\beta &\leftrightarrow \nu_\alpha \bar{\nu}_\beta \\ \nu_\alpha e^\pm &\leftrightarrow \nu_\alpha e^\pm \\ \nu_\alpha \bar{\nu}_\alpha &\leftrightarrow e^+ e^-\end{aligned}$$

$$\mathcal{L}_{\text{SM}} = -2\sqrt{2}G_F \left\{ (\bar{\nu}_e \gamma^\mu L \nu_e)(\bar{e} \gamma_\mu L e) + \sum_{P,\alpha} g_P (\bar{\nu}_\alpha \gamma^\mu L \nu_\alpha)(\bar{e} \gamma_\mu P e) \right\}$$

$$P = L, R = (1 \mp \gamma_5)/2 \quad g_L = -\tfrac{1}{2} + \sin^2 \theta_W \text{ and } g_R = \sin^2 \theta_W$$

Neutrino decoupling

As the Universe expands, particle densities are diluted and temperatures fall. Weak interactions become **ineffective** to keep neutrinos in good thermal contact with the e.m. plasma

Rough, but quite accurate estimate of the decoupling temperature

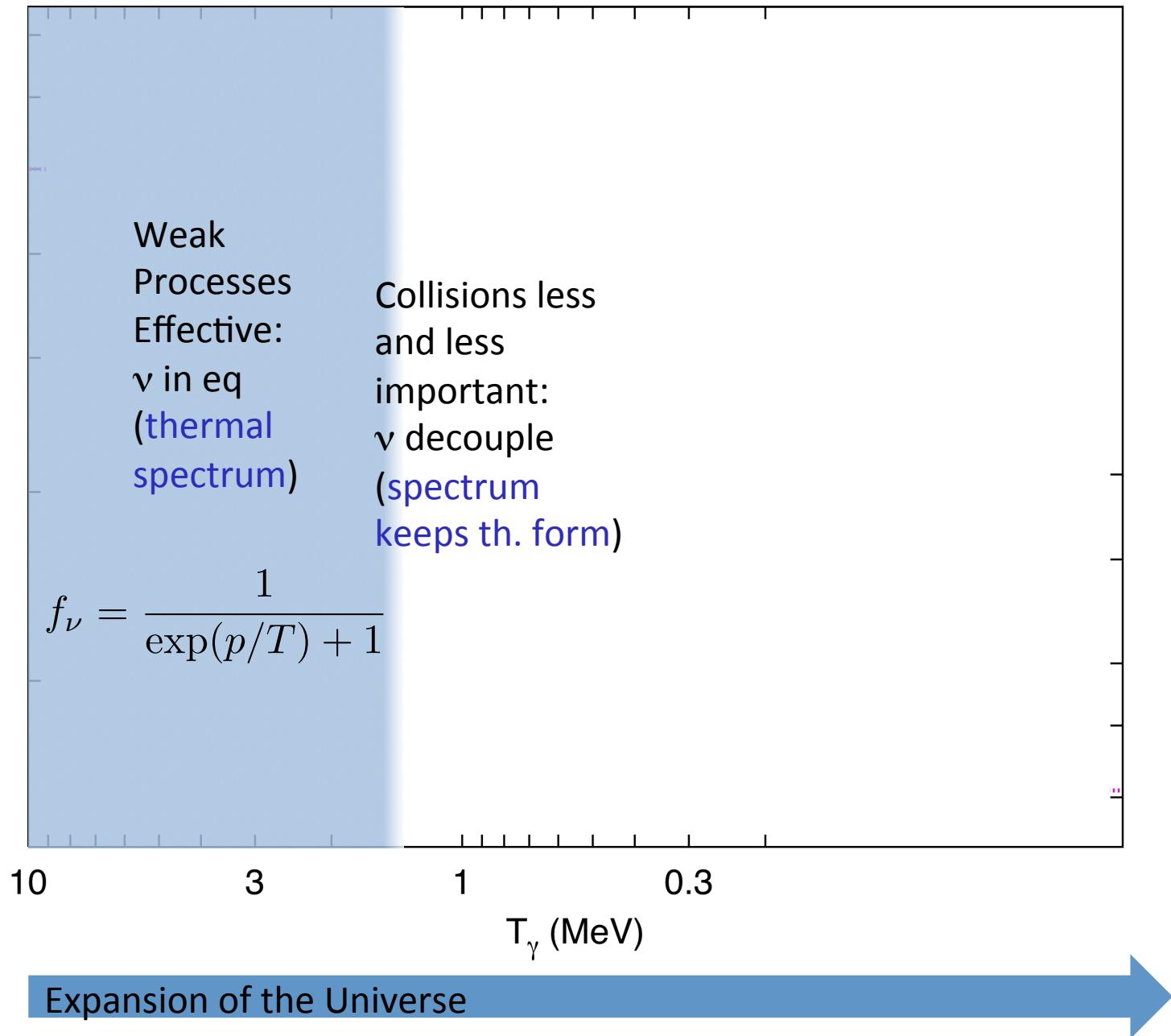
Rate of weak processes \sim Hubble expansion rate

$$\Gamma_W \approx \sigma_W |v| n, \quad H^2 = \frac{8\pi\rho_{\text{rad}}}{3M_P^2} \rightarrow G_F^2 T^5 \approx \sqrt{\frac{8\pi\rho_{\text{rad}}}{3M_P^2}} \rightarrow T_{\text{dec}}(\nu) \approx 1 \text{ MeV}$$

Since ν_e have both CC and NC interactions with e^\pm

$$T_{\text{dec}}(\nu_e) \simeq 2 \text{ MeV} \quad T_{\text{dec}}(\nu_{\mu,\tau}) \simeq 3 \text{ MeV}$$

Neutrino decoupling



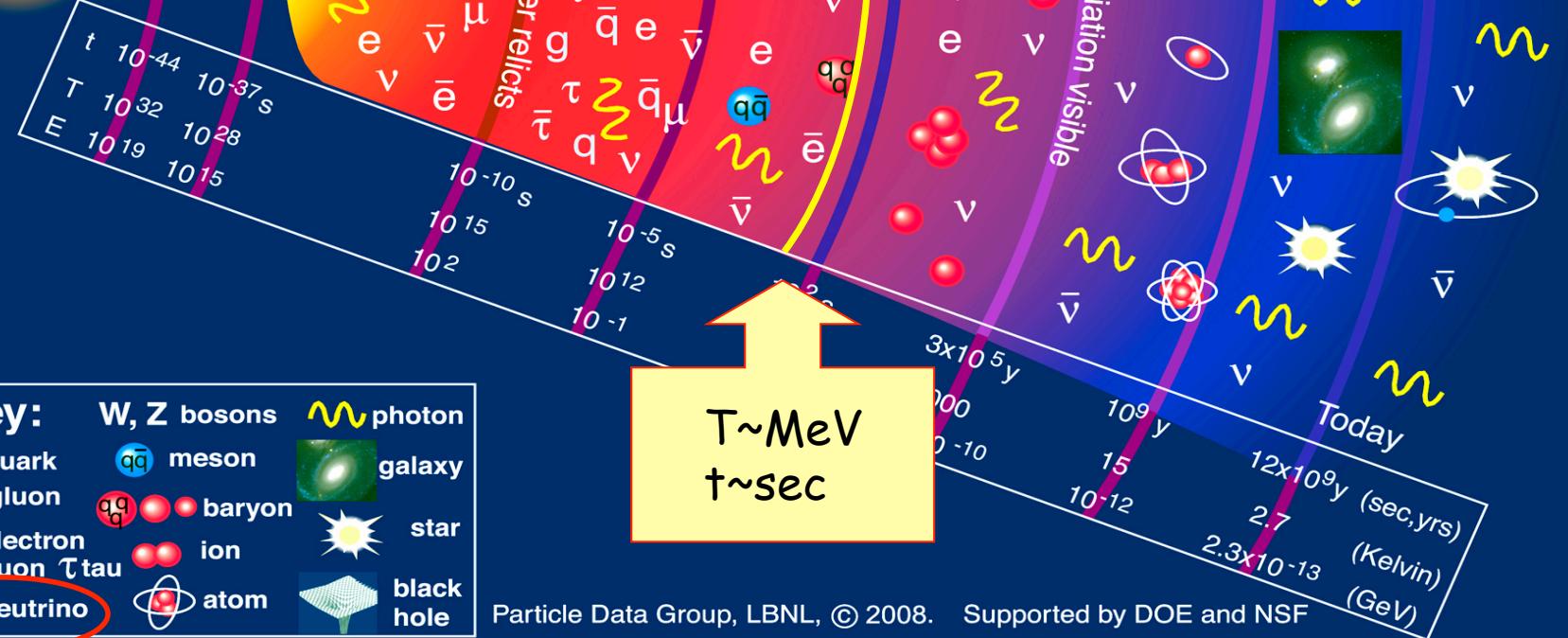
History of the Universe

Neutrinos coupled by weak interactions

$$f_\nu(p, T) = \frac{1}{\exp(p/T) + 1}$$

BIG BANG

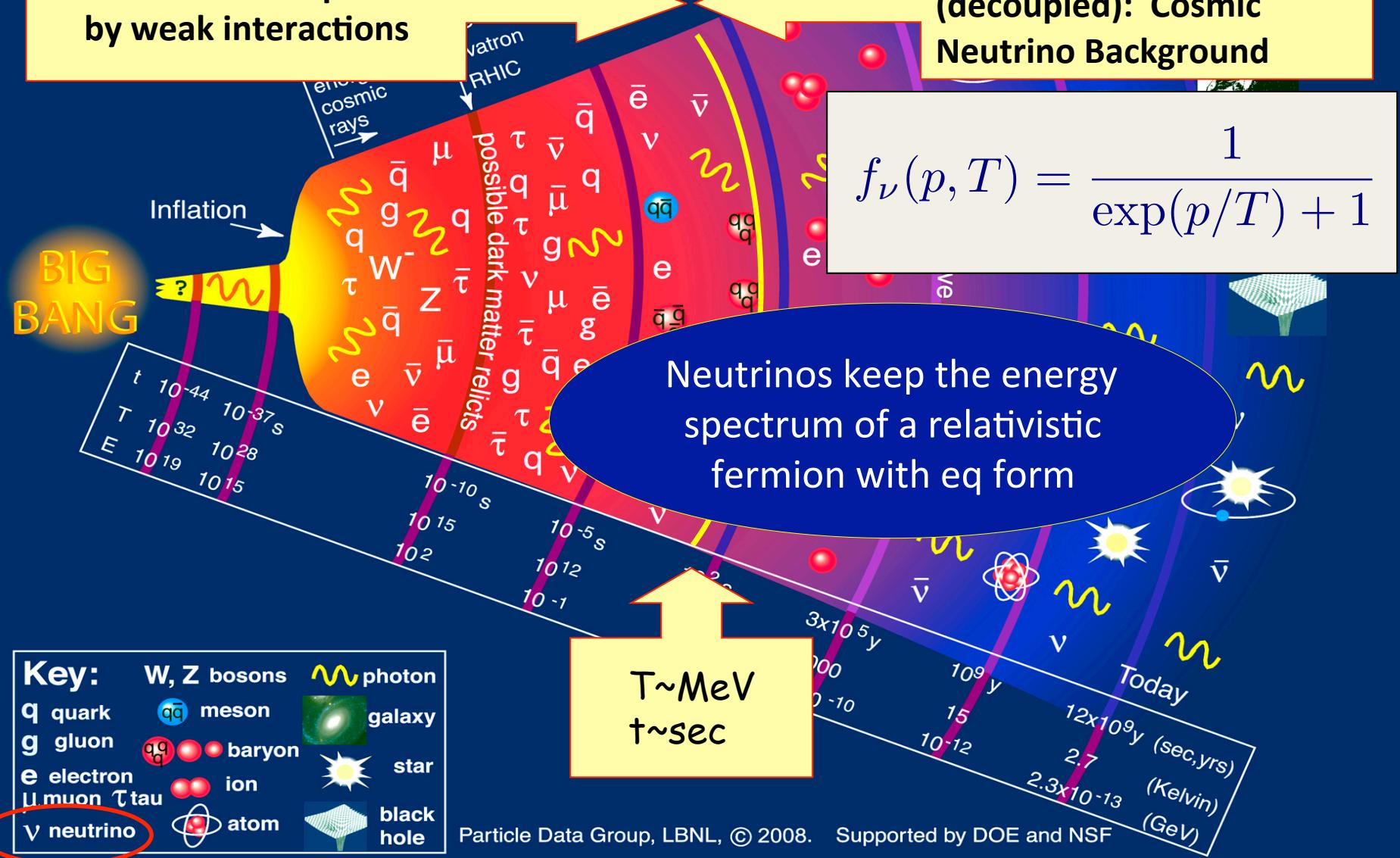
Key:	W, Z bosons	Photon
q quark	qq meson	W
g gluon	qq baryon	Z
e electron	ion	gamma
μ muon	atom	galaxy
τ tau	black hole	star
ν neutrino		



History of the Universe

Neutrinos coupled by weak interactions

Free-streaming neutrinos (decoupled): Cosmic Neutrino Background



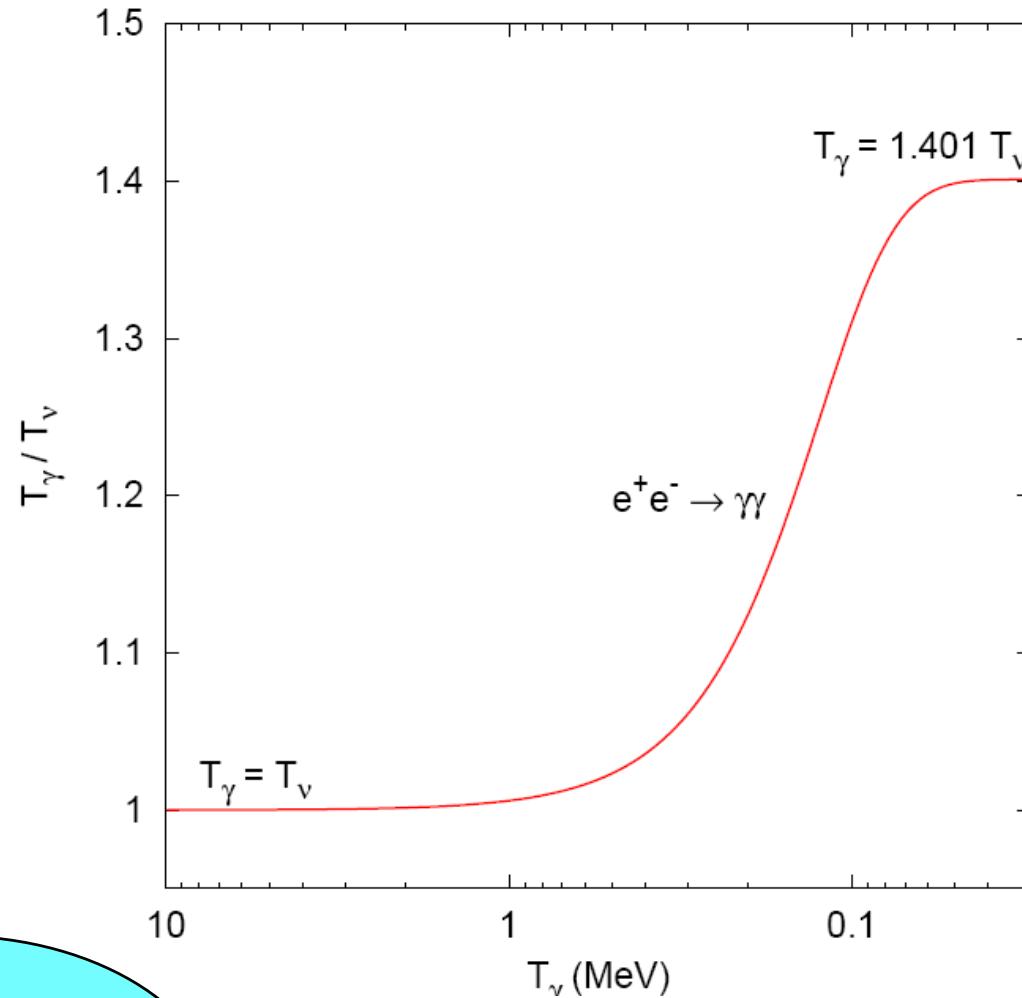
Neutrino and photon (CMB) temperatures

At $T \sim m_e$,
electron-
positron pairs
annihilate



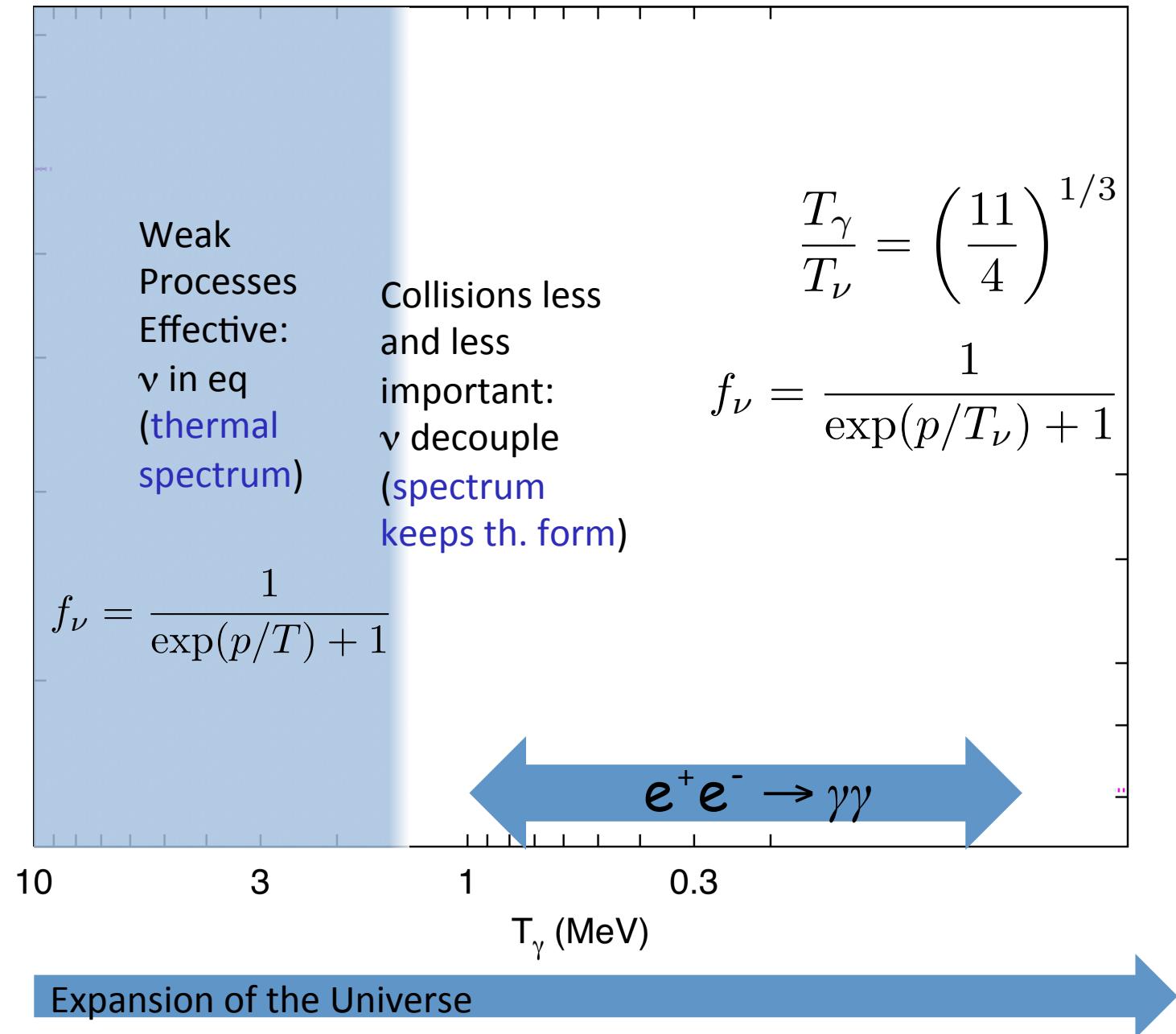
heating photons
but not the
decoupled
neutrinos

$$\frac{T_\gamma}{T_\nu} = \left(\frac{11}{4}\right)^{1/3}$$



$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

Neutrino decoupling and e^\pm annihilations



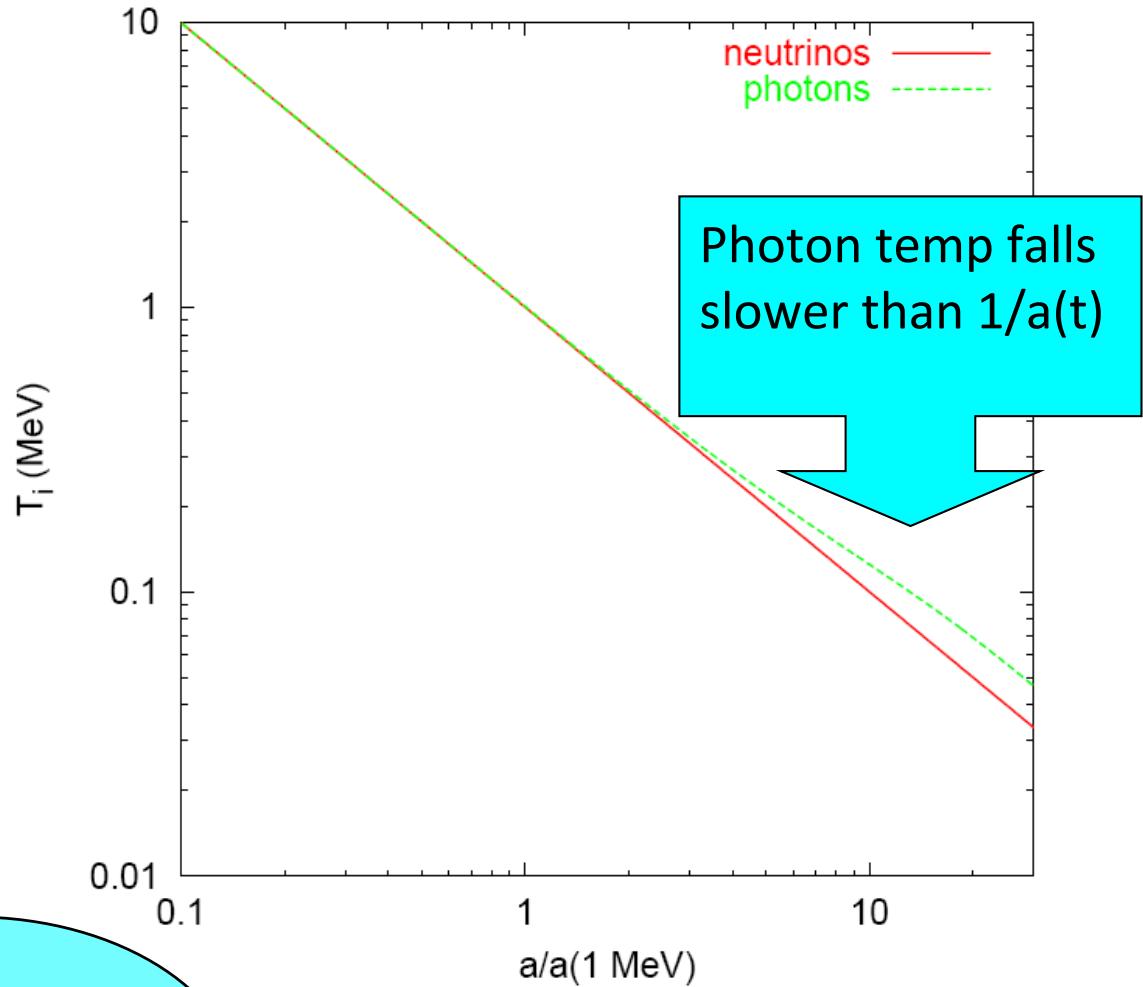
Neutrino and Photon (CMB) temperatures

At $T \sim m_e$,
electron-
positron pairs
annihilate



heating photons
but not the
decoupled
neutrinos

$$\frac{T_\gamma}{T_\nu} = \left(\frac{11}{4}\right)^{1/3}$$



$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

The Cosmic Neutrino Background

Neutrinos decoupled at $T \sim \text{MeV}$, keeping a spectrum as that of a relativistic species

$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

- Number density

$$n_\nu = \int \frac{d^3 p}{(2\pi)^3} f_\nu(p, T_\nu) = \frac{3}{11} n_\gamma = \frac{6\zeta(3)}{11\pi^2} T_{\text{CMB}}^3$$

- Energy density

$$\rho_{\nu_i} = \int \sqrt{p^2 + m_{\nu_i}^2} \frac{d^3 p}{(2\pi)^3} f_\nu(p, T_\nu) \rightarrow \begin{cases} \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_{\text{CMB}}^4 & \text{Massless} \\ m_{\nu_i} n_\nu & \text{Massive } m_\nu \gg T \end{cases}$$

The Cosmic Neutrino Background

Neutrinos decoupled at $T \sim \text{MeV}$, keeping a spectrum as that of a relativistic species

$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

- Number density

At present $112 (\nu + \bar{\nu}) \text{ cm}^{-3}$ per flavour

- Energy density

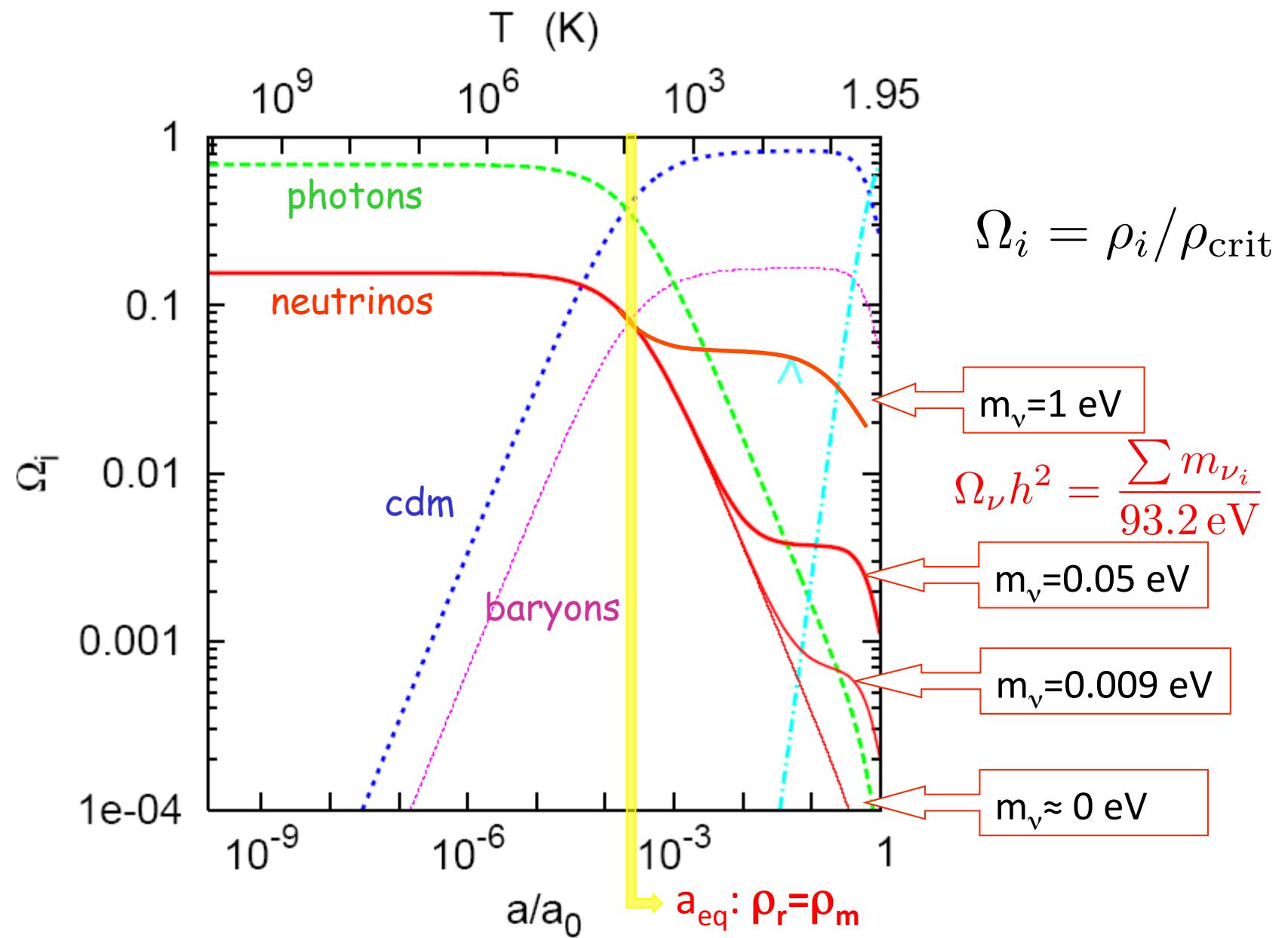
$$\Omega_\nu h^2 \simeq 1.7 \times 10^{-5} \quad \text{Massless}$$

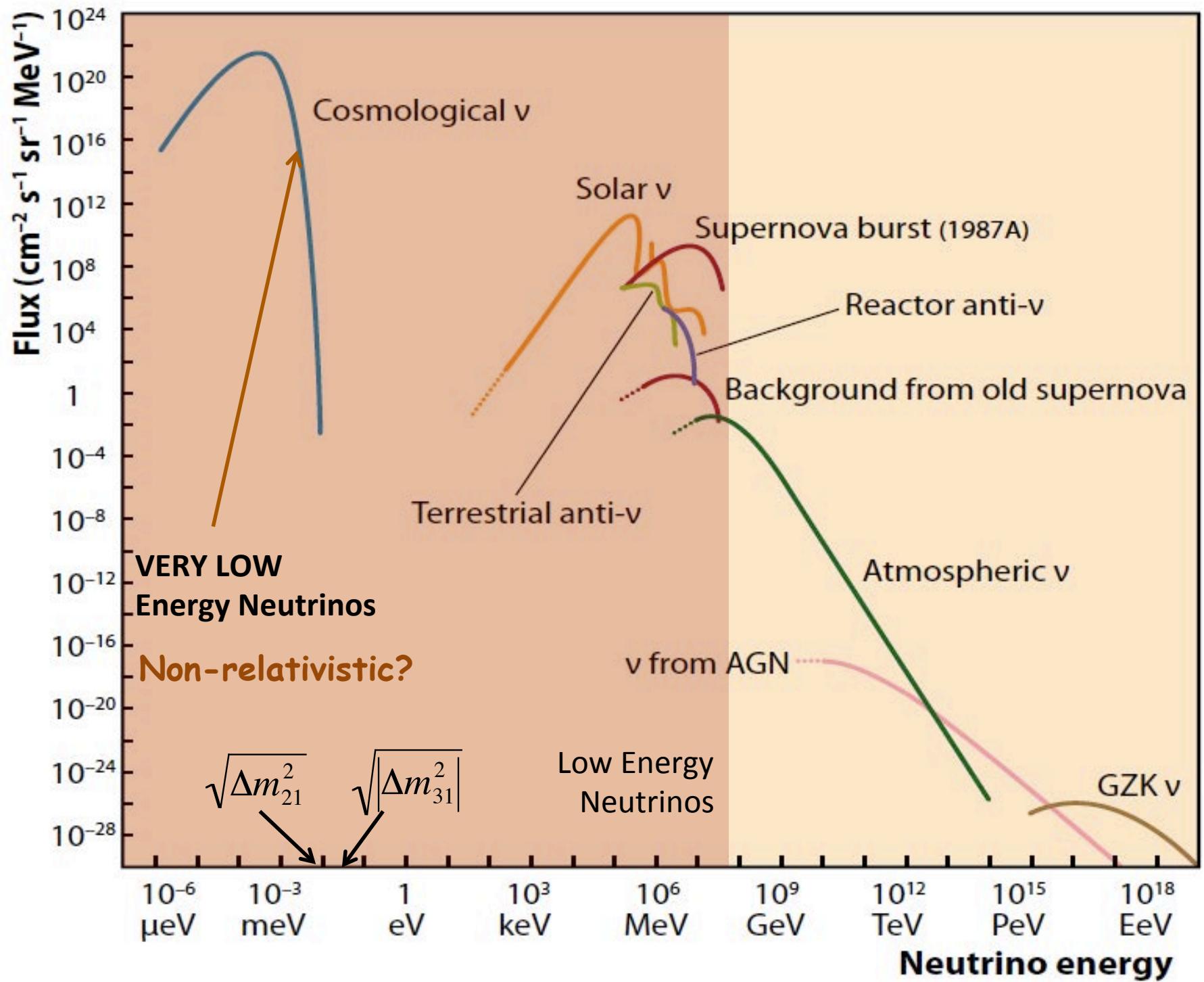
Contribution to the energy density of the Universe

$$\Omega_\nu h^2 = \frac{\sum_i m_{\nu_i}}{94.1 \text{ eV}}$$

Massive
 $m_\nu \gg T$

Evolution of the background densities: 1 MeV → now





The radiation content of the Universe (N_{eff})

Relativistic particles in the Universe

At $T >> m_e$, the radiation content of the Universe is

$$\rho_r = \rho_\gamma + \rho_\nu = \frac{\pi^2}{15} T^4 + 3 \times \frac{7}{8} \times \frac{\pi^2}{15} T^4 = \left[1 + \frac{7}{8} \times 3 \right] \rho_\gamma$$

At $T < m_e$, the radiation content of the Universe is

$$\rho_r = \rho_\gamma + \rho_\nu = \frac{\pi^2}{15} T_\gamma^4 + 3 \times \frac{7}{8} \times \frac{\pi^2}{15} T_\nu^4 = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} 3 \right] \rho_\gamma$$

$$\rho_r = \rho_\gamma + \rho_\nu = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma \quad \frac{T_\nu^4}{T_\gamma^4}$$

of flavour neutrinos: $N_\nu = 2.984 \pm 0.008$ (LEP data)

Relativistic particles in the Universe

At $T < m_e$, the radiation content of the Universe is

$$\rho_r = \rho_\gamma + \rho_\nu + \rho_x = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma$$

Effective number of relativistic neutrino species

Traditional parametrization of ρ stored in relativistic particles

N_{eff} is a way to measure the ratio $\frac{\rho_\nu + \rho_x}{\rho_\gamma}$

- standard neutrinos only: $N_{\text{eff}} \simeq 3$ (3.045)
- $N_{\text{eff}} > 3$ (delays equality time) from **additional relativistic particles** (scalars, pseudoscalars, decay products of heavy particles,...) or **non-standard neutrino physics** (primordial neutrino asymmetries, totally or partially thermalized light sterile neutrinos, non-standard interactions with electrons,...)

Bounds on N_{eff} from
Primordial Nucleosynthesis
and other cosmological
observables (CMB+LSS)

Exercises: try to calculate...

- The present number density of massive/massless neutrinos n_ν^0 in cm^{-3}
 - The present energy density of massive/massless neutrinos Ω_ν^0 and find the limits on the total neutrino mass from $\Omega_\nu^0 < 1$ and $\Omega_\nu^0 < \Omega_m^0$
 - The final ratio T_γ/T_ν using the conservation of entropy density before/after e^\pm annihilations
 - The decoupling temperature of relic neutrinos using $\Gamma_w \approx H$
-
- The photon temperature / redshift of the matter radiation equality for $m_\nu = 1 \text{ eV}$

End of 1st lecture