

Flavour and CP Phenomenology

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1 CP-Generalities

2 FCNC $b \rightarrow s\gamma$

3 FCNC, CP $\Delta F = 2$

4 More stuff

Disclaimer

- *Before I came here I was confused about this subject.
Having listened to your lecture I am still confused.
But on a higher level.*

E. Fermi

CP – Generalities

CP Violation in Neutral Meson Mixing

■ Decays $\Delta F = \Delta Q$

$$\begin{array}{ll} M^0 \rightarrow X\ell^+ & \bar{M}^0 \nrightarrow X\ell^+ \\ M^0 \nrightarrow X\ell^- & \bar{M}^0 \rightarrow X\ell^- \end{array}$$

$$\lambda_f = \frac{q}{p} \frac{\bar{A}_f}{A_f} \quad \text{give} \quad C_{X\ell^\pm} = \pm 1, \quad S_{X\ell^\pm} = R_{X\ell^\pm} = 0$$

■ Decays $\Leftrightarrow$ Transition Probabilities

$$\begin{aligned} |\langle X\ell^- | \mathcal{T} | M^0(t) \rangle|^2 &\sim |\langle \bar{M}^0 | \mathcal{T} | M^0(t) \rangle|^2 \\ |\langle X\ell^+ | \mathcal{T} | \bar{M}^0(t) \rangle|^2 &\sim |\langle M^0 | \mathcal{T} | \bar{M}^0(t) \rangle|^2 \end{aligned}$$

$$\begin{aligned} |\langle f | \mathcal{T} | M^0(t) \rangle|^2 &= e^{-\Gamma t} \left\{ \mathcal{C}_h[M^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[M^0, f] \cos(\Delta Mt) \right. \\ &\quad \left. + \mathcal{S}_h[M^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[M^0, f] \sin(\Delta Mt) \right\} \\ |\langle f | \mathcal{T} | \bar{M}^0(t) \rangle|^2 &= e^{-\Gamma t} \left\{ \mathcal{C}_h[\bar{M}^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[\bar{M}^0, f] \cos(\Delta Mt) \right. \\ &\quad \left. + \mathcal{S}_h[\bar{M}^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[\bar{M}^0, f] \sin(\Delta Mt) \right\} \end{aligned}$$

CP Violation in Neutral Meson Mixing

- We have

$$|\langle \bar{M}^0 | \mathcal{T} | M^0(t) \rangle|^2 \sim e^{-\Gamma t} \left\{ \mathcal{C}_h[M^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[M^0, f] \cos(\Delta Mt) \right. \\ \left. + \mathcal{S}_h[M^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[M^0, f] \sin(\Delta Mt) \right\}$$

$$|\langle M^0 | \mathcal{T} | \bar{M}^0(t) \rangle|^2 \sim e^{-\Gamma t} \left\{ \mathcal{C}_h[\bar{M}^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[\bar{M}^0, f] \cos(\Delta Mt) \right. \\ \left. + \mathcal{S}_h[\bar{M}^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[\bar{M}^0, f] \sin(\Delta Mt) \right\}$$

with $(\theta = 0)$

$$\mathcal{C}_h[M^0, f] = \frac{\Gamma_f(1-\delta)}{2(1-C_f\delta)} \{2\}$$

$$\mathcal{C}_c[M^0, f] = \frac{\Gamma_f(1-\delta)}{2(1-C_f\delta)} \{2C_f\}$$

$$\mathcal{S}_h[M^0, f] = \frac{\Gamma_f(1-\delta)}{2(1-C_f\delta)} \{2R_f\}$$

$$\mathcal{S}_c[M^0, f] = \frac{\Gamma_f(1-\delta)}{2(1-C_f\delta)} \{-2S_f\}$$

$$\mathcal{C}_h[\bar{M}^0, f] = \frac{\Gamma_f(1+\delta)}{2(1-C_f\delta)} \{2\}$$

$$\mathcal{C}_c[\bar{M}^0, f] = \frac{\Gamma_f(1+\delta)}{2(1-C_f\delta)} \{-2C_f\}$$

$$\mathcal{S}_h[\bar{M}^0, f] = \frac{\Gamma_f(1+\delta)}{2(1-C_f\delta)} \{2R_f\}$$

$$\mathcal{S}_c[\bar{M}^0, f] = \frac{\Gamma_f(1+\delta)}{2(1-C_f\delta)} \{2S_f\}$$

CP Violation in Neutral Meson Mixing

- We have

$$|\langle \bar{M}^0 | \mathcal{T} | M^0(t) \rangle|^2 \sim e^{-\Gamma t} \left\{ \mathcal{C}_h[M^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[M^0, f] \cos(\Delta Mt) \right. \\ \left. + \mathcal{S}_h[M^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[M^0, f] \sin(\Delta Mt) \right\}$$

$$|\langle M^0 | \mathcal{T} | \bar{M}^0(t) \rangle|^2 \sim e^{-\Gamma t} \left\{ \mathcal{C}_h[\bar{M}^0, f] \cosh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{C}_c[\bar{M}^0, f] \cos(\Delta Mt) \right. \\ \left. + \mathcal{S}_h[\bar{M}^0, f] \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \mathcal{S}_c[\bar{M}^0, f] \sin(\Delta Mt) \right\}$$

with $(\theta = 0)$

$$\mathcal{C}_h[M^0, \bar{M}^0] = \frac{\Gamma(1-\delta)}{(1+\delta)} \quad \mathcal{C}_c[M^0, \bar{M}^0] = -\frac{\Gamma(1-\delta)}{(1+\delta)}$$

$$\mathcal{S}_h[M^0, \bar{M}^0] = 0 \quad \mathcal{S}_c[M^0, \bar{M}^0] = 0$$

$$\mathcal{C}_h[\bar{M}^0, M^0] = \frac{\Gamma(1+\delta)}{(1-\delta)} \quad \mathcal{C}_c[\bar{M}^0, M^0] = -\frac{\Gamma(1+\delta)}{(1-\delta)}$$

$$\mathcal{S}_h[\bar{M}^0, M^0] = 0 \quad \mathcal{S}_c[\bar{M}^0, M^0] = 0$$

CP Violation in Neutral Meson Mixing

- We have with ($\theta = 0$)

$$|\langle \bar{M}^0 | \mathcal{T} | M^0(t) \rangle|^2 \sim e^{-\Gamma t} \Gamma \left\{ \cosh\left(\frac{\Delta\Gamma}{2}t\right) - \cos(\Delta Mt) \right\} \frac{(1-\delta)}{(1+\delta)}$$

$$|\langle M^0 | \mathcal{T} | \bar{M}^0(t) \rangle|^2 \sim e^{-\Gamma t} \Gamma \left\{ \cosh\left(\frac{\Delta\Gamma}{2}t\right) - \cos(\Delta Mt) \right\} \frac{(1+\delta)}{(1-\delta)}$$

CP Violation in Neutral Meson Mixing

- We have with $(\theta = 0)$
- CP Violating asymmetry

$$\left| \langle \bar{M}^0 | \mathcal{T} | M^0(t) \rangle \right|^2 - \left| \langle M^0 | \mathcal{T} | \bar{M}^0(t) \rangle \right|^2 \sim \frac{-4\delta}{1 - \delta^2} e^{-\Gamma t} \Gamma \left[\cosh \left(\frac{\Delta\Gamma}{2} t \right) - \cos(\Delta M t) \right]$$

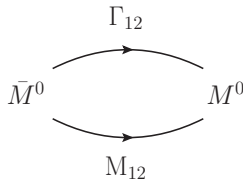
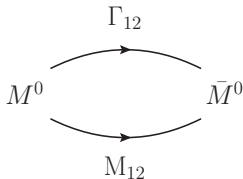
CP Violation in Neutral Meson Mixing

- H with $\theta = 0$

$$H = \begin{pmatrix} \mu & \frac{p}{q} \frac{\Delta\mu}{2} \\ \frac{q}{p} \frac{\Delta\mu}{2} & \mu \end{pmatrix}$$

- $\frac{H_{12}}{H_{21}} = \frac{q^2}{p^2}$

$$\left| \frac{q}{p} \right|^2 = \frac{|M_{12}^* - \frac{i}{2}\Gamma_{12}^*|^2}{|M_{12} - \frac{i}{2}\Gamma_{12}|^2} \Leftrightarrow \delta = -\frac{\text{Im}(M_{12}\Gamma_{12}^*)}{|M_{12}|^2 + \frac{1}{4}|\Gamma_{12}|^2}$$



Direct CP Violation

- CP transformation

$$\mathbf{CP}|f\rangle = e^{i\varphi_f}|\bar{f}\rangle \quad \mathbf{CP}|\bar{f}\rangle = e^{-i\varphi_f}|f\rangle$$

- Invariance under CP of transition operator $\mathcal{T}$

$$\mathbf{CP} \mathcal{T} \mathbf{CP} = \mathcal{T}$$

- CP invariance of $f \rightarrow g$ transition

$$\langle g|\mathcal{T}|f\rangle = e^{i(\varphi_f - \varphi_g)}\langle \bar{g}|\mathcal{T}|\bar{f}\rangle$$

- Simplest CP asymmetry

$$\mathcal{A}_{CP}(f \rightarrow g) = |\langle g|\mathcal{T}|f\rangle|^2 - |\langle \bar{g}|\mathcal{T}|\bar{f}\rangle|^2$$

$\mathcal{A}_{CP}(f \rightarrow g) \neq 0$ in decays of CP conjugate channels:

direct CP Violation

Direct CP Violation

- In general

$$\langle g|\mathcal{T}|f\rangle = A e^{i\delta} e^{i\phi} \quad \langle \bar{g}|\mathcal{T}|\bar{f}\rangle = \bar{A} e^{i\delta} e^{-i\phi}$$

δ : *strong* phase, ϕ : *weak* phase

$$\mathcal{A}_{CP}(f \rightarrow g) = A^2 - \bar{A}^2 \quad A \neq \bar{A} \Rightarrow \text{CP Violation}$$

but in SM, from $\mathcal{L}_{CC}$, expect $A = \bar{A}$ at tree level

- Consider two amplitudes

$$\langle g|\mathcal{T}|f\rangle = A_1 e^{i(\delta_1+\phi_1)} + A_2 e^{i(\delta_2+\phi_2)}$$

$$\langle \bar{g}|\mathcal{T}|\bar{f}\rangle = A_1 e^{i(\delta_1-\phi_1)} + A_2 e^{i(\delta_2-\phi_2)}$$

$$\mathcal{A}_{CP}(f \rightarrow g) = -4A_1A_2 \sin(\delta_2 - \delta_1) \sin(\phi_2 - \phi_1)$$

Indirect CP Violation

- Indirect CP Violation
 - Interference CP Violation
 - Mixing $\times$ Decay
- CP Violating asymmetry

$$|\langle f|\mathcal{T}|M^0(t)\rangle|^2 - |\langle \bar{f}|\mathcal{T}|\bar{M}^0(t)\rangle|^2$$

- If No CP Violation

$$\langle f|\mathcal{T}|M^0\rangle = e^{i(\varphi_B - \varphi_f)} \langle \bar{f}|\mathcal{T}|\bar{M}^0\rangle \quad \langle f|\mathcal{T}|\bar{M}^0\rangle = e^{-i(\varphi_B + \varphi_f)} \langle \bar{f}|\mathcal{T}|M^0\rangle$$

$$\langle M^0|\mathcal{H}|\bar{M}^0\rangle = e^{-i2\varphi_B} \langle \bar{M}^0|\mathcal{H}|M^0\rangle$$

Indirect CP Violation

- That is

$$A_f = e^{i(\varphi_B - \varphi_f)} \bar{A}_{\bar{f}} \quad \bar{A}_f = e^{-i(\varphi_B + \varphi_f)} A_{\bar{f}}$$

$$\frac{p}{q} = e^{-i2\varphi_B} \frac{q}{p}$$

- Combine

$$\frac{q}{p} \frac{\bar{A}_f}{A_f} = \frac{p}{q} \frac{A_{\bar{f}}}{\bar{A}_{\bar{f}}} \Leftrightarrow \lambda_f \lambda_{\bar{f}} = 1$$

$$C_f = -C_{\bar{f}} \quad S_f = -S_{\bar{f}} \quad R_f = R_{\bar{f}}$$

- For a CP eigenstate f_{CP}

$$\lambda_{f_{CP}} = \pm 1$$

$$C_{f_{CP}} = 0 \quad S_{f_{CP}} = 0 \quad R_{f_{CP}} = \pm 1$$

Rephasing

■ Arbitrary rephasings

- States: $|M^0\rangle \mapsto e^{i\omega}|M^0\rangle, |\bar{M}^0\rangle \mapsto e^{i\bar{\omega}}|\bar{M}^0\rangle$
- States: $|f\rangle \mapsto e^{i\omega_f}|f\rangle$
- CP definition: $\mathbf{CP}|f\rangle = e^{i\varphi_f}|\bar{f}\rangle, \mathbf{CP}|\bar{f}\rangle = e^{-i\varphi_f}|f\rangle$

■ Decay amplitudes

$$\left. \begin{aligned} \langle f|\mathcal{T}|M^0\rangle &= A_f \mapsto e^{i\omega}e^{-i\omega_f}A_f \\ \langle f|\mathcal{T}|\bar{M}^0\rangle &= \bar{A}_f \mapsto e^{i\bar{\omega}}e^{-i\omega_f}\bar{A}_f \end{aligned} \right\} |A_f|, |\bar{A}_f| \text{ are invariant}$$

■ Meson Mixing $H_{21} \mapsto e^{i(\omega-\bar{\omega})}H_{21}$

$$\frac{q}{p} \mapsto e^{i(\omega-\bar{\omega})}\frac{q}{p} \Rightarrow |q/p| \text{ is invariant } (\Leftrightarrow \delta)$$

■ Meson Mixing $\times$ Decay amplitudes

$$\frac{q}{p} \frac{\langle f|\mathcal{T}|\bar{M}^0\rangle}{\langle f|\mathcal{T}|M^0\rangle} = \frac{q}{p} \frac{\bar{A}_f}{A_f} = \lambda_f \text{ is invariant}$$

Neutral Meson Mixing (again)

- H with $\theta = 0$

$$H = \begin{pmatrix} \mu & \frac{p}{q} \frac{\Delta\mu}{2} \\ \frac{q}{p} \frac{\Delta\mu}{2} & \mu \end{pmatrix}$$

- $H_{12}H_{21} = \frac{1}{4}(\Delta\mu)^2$

$$H_{12}H_{21} = |M_{12}|^2 - \frac{1}{4}|\Gamma_{12}|^2 - i\text{Re}(M_{12}\Gamma_{12}^*)$$

$$\frac{(\Delta\mu)^2}{4} = \frac{1}{4} \left[(\Delta M)^2 - \frac{1}{4}(\Delta\Gamma)^2 - i(\Delta M)(\Delta\Gamma) \right]$$

- $\frac{H_{12}}{H_{21}} = \frac{q^2}{p^2}$

$$\left| \frac{q}{p} \right|^2 = \frac{|M_{12}^* - \frac{i}{2}\Gamma_{12}^*|^2}{|M_{12} - \frac{i}{2}\Gamma_{12}|^2} \Leftrightarrow \delta = -\frac{\text{Im}(M_{12}\Gamma_{12}^*)}{|M_{12}|^2 + \frac{1}{4}|\Gamma_{12}|^2}$$

B_d, B_s systems

$$|M_{12}|^2 - \frac{1}{4}|\Gamma_{12}|^2 - i\text{Re}(M_{12}\Gamma_{12}^*) = \frac{1}{4} \left[(\Delta M)^2 - \frac{1}{4}(\Delta\Gamma)^2 - i(\Delta M)(\Delta\Gamma) \right]$$

$$\delta = -\frac{\text{Im}(M_{12}\Gamma_{12}^*)}{|M_{12}|^2 + \frac{1}{4}|\Gamma_{12}|^2}$$

- If $|\Gamma_{12}| \ll |M_{12}|$

$$4|M_{12}|^2 = (\Delta M)^2 - \frac{1}{4}(\Delta\Gamma)^2, \quad 0 = (\Delta M)(\Delta\Gamma), \quad \delta = 0$$

- No CP Violation in mixing: $\delta = 0 \Leftrightarrow \left| \frac{q}{p} \right| = 1$

- $\Delta\Gamma = 0$ and $2|M_{12}| = \Delta M$

$$\frac{q}{p} = e^{-i \arg(M_{12})} = \frac{M_{12}^*}{|M_{12}|} = \frac{p^*}{q^*}$$

Kaons

$$|M_{12}|^2 - \frac{1}{4}|\Gamma_{12}|^2 - i\text{Re}(M_{12}\Gamma_{12}^*) = \frac{1}{4} \left[(\Delta M)^2 - \frac{1}{4}(\Delta\Gamma)^2 - i(\Delta M)(\Delta\Gamma) \right]$$

$$\delta = -\frac{\text{Im}(M_{12}\Gamma_{12}^*)}{|M_{12}|^2 + \frac{1}{4}|\Gamma_{12}|^2}$$

- Relevant fact 1:

$$|\Delta\Gamma| \simeq 2|\Delta M| \Rightarrow 2|M_{12}| \simeq |\Gamma_{12}|$$

- Relevant fact 2:

$$\begin{aligned} \text{Re}(M_{12}\Gamma_{12}^*) &\simeq |M_{12}\Gamma_{12}| = (\Delta M)(\Delta\Gamma) \\ \Rightarrow \text{Re}(M_{12}\Gamma_{12}^*) &\simeq \frac{1}{2}(\Delta M)^2 \quad 2|M_{12}| \simeq \Delta M \end{aligned}$$

Kaons

$$2|M_{12}| \simeq |\Gamma_{12}| \simeq \Delta M$$

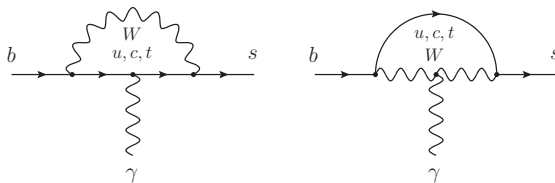
- CP Violation in mixing

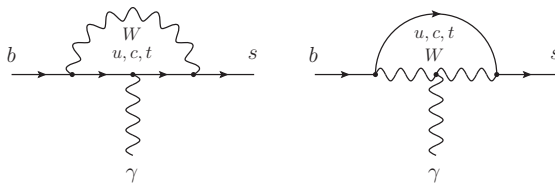
$$\delta = -\frac{\text{Im}(M_{12}\Gamma_{12}^*)}{|M_{12}|^2 + \frac{1}{4}|\Gamma_{12}|^2} = -\frac{\text{Im}(M_{12}e^{-i\arg(\Gamma_{12})})}{\Delta M}$$

- In kaons, since $K_S \sim K_+$ and $K_L \sim K_-$,
the “traditional” notation is not q and p :

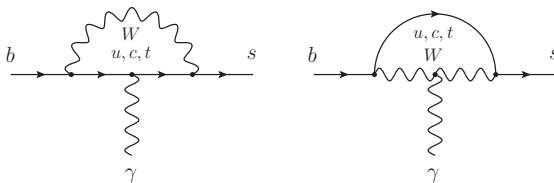
$$\frac{q}{p} = \frac{1 - \epsilon_K}{1 + \epsilon_K} \Rightarrow \delta = \frac{2\text{Re}(\epsilon_K)}{1 + |\epsilon_K|^2}$$

$$\text{FCNC} \quad b \rightarrow s\gamma$$

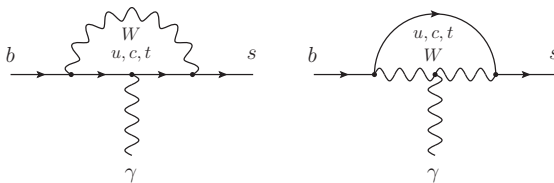
 e



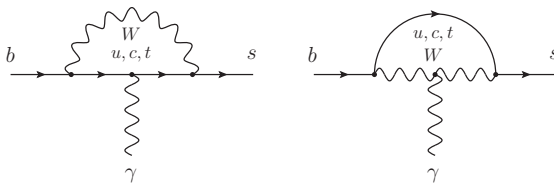
$$e g^2$$



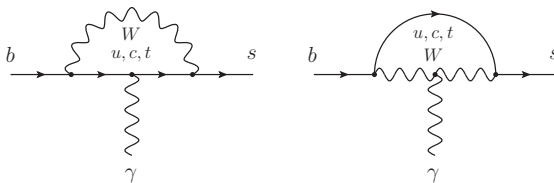
$$e g^2 V_{jb} V_{js}^*$$



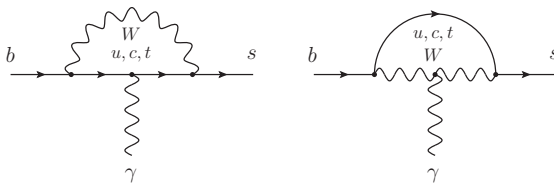
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2}$$



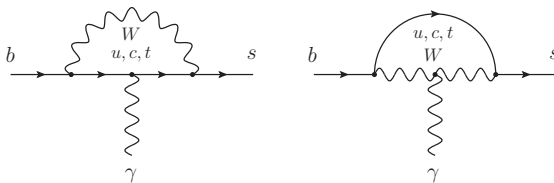
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu}$$



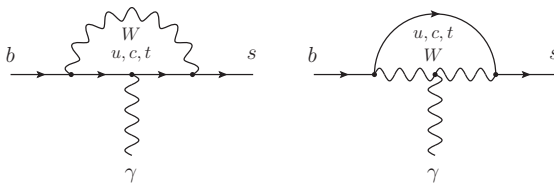
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L b_L$$



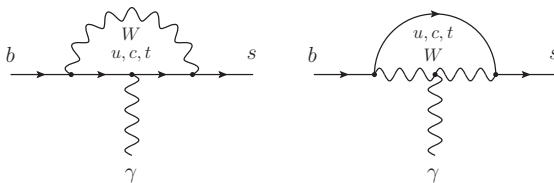
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_L$$



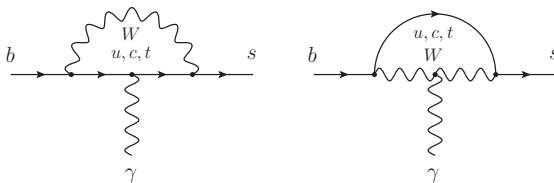
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_L$$



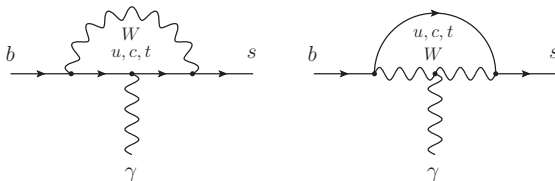
$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_R m_b$$



$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_R m_b \hat{I}(m_{u_j}^2, M_W^2)$$



$$e g^2 V_{jb} V_{js}^* \frac{1}{16\pi^2} F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_R m_b \hat{I}(x_j) M_W^{-2}$$



$$\sum_{j=u,c,t} \frac{eg^2}{16\pi^2 M_W^2} V_{jb} V_{js}^* \hat{I}(x_j) m_b F_{\mu\nu} \bar{s}_L \sigma^{\mu\nu} b_R, \quad x_j = \frac{m_{u_j}^2}{M_W^2}$$

Glashow-Iliopoulos-Maiani (in $\Delta F = 1$)

- Loop function $\hat{I}(x_j)$ in

$$\sum_{j=1}^3 V_{jb} V_{js}^* \hat{I}(x_j)$$

- CKM unitarity

$$\sum_{j=1}^3 V_{jb} V_{js}^* = 0 \quad \sum_{j=1}^3 V_{jb} V_{js}^* \hat{I}(0) = 0$$

- Then

$$\sum_{j=1}^3 V_{jb} V_{js}^* \hat{I}(x_j) = \sum_{j=1}^3 V_{jb} V_{js}^* I(x_j)$$

with

$$I(x_j) \equiv \hat{I}(x_j) - \hat{I}(0), \quad \lim_{x \rightarrow 0} I(x) = 0$$

- In $b \rightarrow s\gamma$, t contribution dominant

FCNC and CP in $\Delta F = 2$

GIM (in $\Delta F = 2$)

- Loop function $\hat{S}(x_j, x_k)$

$$\sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* \hat{S}(x_j, x_k) \quad \text{sym. } j \leftrightarrow k$$

- CKM unitarity *twice*

$$\sum_{k=1}^3 V_{k\alpha} V_{k\beta}^* \left[\sum_{j=1}^3 V_{j\alpha} V_{j\beta}^* \hat{S}(0, x_k) \right] = 0 \quad \left[\sum_{j=1}^3 V_{j\alpha} V_{j\beta}^* \right] \left[\sum_{k=1}^3 V_{k\alpha} V_{k\beta}^* \right] \hat{S}(0, 0) = 0$$

- Then

$$\sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* \hat{S}(x_j, x_k) = \sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* S(x_j, x_k)$$

$$S(x_j, x_k) \equiv \hat{S}(x_j, x_k)$$

- Largest function: $S(x_t, x_t) \gg S(x_j, x_t)$

GIM (in $\Delta F = 2$)

- Loop function $\hat{S}(x_j, x_k)$

$$\sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* \hat{S}(x_j, x_k) \quad \text{sym. } j \leftrightarrow k$$

- CKM unitarity *twice*

$$\sum_{k=1}^3 V_{k\alpha} V_{k\beta}^* \left[\sum_{j=1}^3 V_{j\alpha} V_{j\beta}^* \hat{S}(0, x_k) \right] = 0 \quad \left[\sum_{j=1}^3 V_{j\alpha} V_{j\beta}^* \right] \left[\sum_{k=1}^3 V_{k\alpha} V_{k\beta}^* \right] \hat{S}(0, 0) = 0$$

- Then

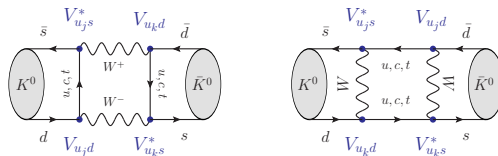
$$\sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* \hat{S}(x_j, x_k) = \sum_{j,k=1}^3 V_{j\alpha} V_{j\beta}^* V_{k\alpha} V_{k\beta}^* S(x_j, x_k)$$

$$S(x_j, x_k) \equiv \hat{S}(x_j, x_k) - \hat{S}(x_j, 0) - \hat{S}(0, x_k) + \hat{S}(0, 0), \quad \lim_{x,y \rightarrow 0} S(x, y) = 0$$

- Largest function: $S(x_t, x_t) \gg S(x_j, x_t)$

ϵ_K, δ in kaons

■ For M_{12}



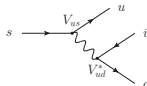
$$M_{12} \propto \sum_{j,k=1}^3 V_{js} V_{jd}^* V_{ks} V_{kd}^* L(x_j, x_k)$$

ϵ_K, δ in kaons

■ For Γ_{12}

$$\Gamma_{12} = 2\pi \sum_n \delta(M_0 - E_n) \langle i | \mathcal{H}_w | n \rangle \langle n | \mathcal{H}_w | j \rangle$$

■ $\pi\pi$ decays



$$\rightarrow \Gamma_{12} \propto (V_{us} V_{ud}^*)^2$$

■ Back to $\delta = -\frac{\text{Im}(M_{12} e^{-i \arg(\Gamma_{12})})}{\Delta M}$

$$\delta \propto \text{Im} \left(\sum_{j,k=1}^3 V_{js} V_{jd}^* V_{ks} V_{kd}^* L(x_j, x_k) \frac{(V_{us}^* V_{ud})^2}{|V_{us}^* V_{ud}|^2} \right)$$

ϵ_K, δ in kaons

■ Loop function

$$L(x_u) \simeq 1.4 \times 10^{-9} \quad L(x_c) \simeq 7.4 \times 10^{-5} \quad L(x_t) \simeq 2.4$$

$$L(x_u, x_c) \simeq 4.2 \times 10^{-8} \quad L(x_u, x_t) \simeq 3.0 \times 10^{-8} \quad L(x_c, x_t) \simeq 8.0 \times 10^{-4}$$

■ CKM's with Wolfenstein parameters

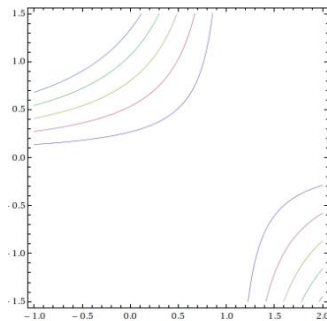
$$V_{cs} V_{cd}^* = -\lambda(1 - \lambda^2/2) \quad V_{ts} V_{td}^* = -A^2 \lambda^5 (1 - \rho + i\eta)$$

■ Imaginary part

$$\begin{aligned} \text{Im}\left(M_{12} e^{-i \arg(\Gamma_{12})}\right) &\propto A^2 \lambda^6 \eta \left[2L(x_c, x_t) + A^2 \lambda^4 L(x_t)(1 - \rho)\right] \\ &\sim \eta [0.15 + 3.6(1 - \rho)] \end{aligned}$$

ϵ_K, δ in kaons

Contours $2L(x_c, x_t) + A^2\lambda^4 L(x_t)(1 - \rho) = [\text{Const.}]$



$\sin 2\beta$

■ CP Asymmetry in $B_d^0 \rightarrow J\Psi K_S$

■ Decays:

$$\langle J/\Psi K_S | \mathcal{T} | B_d^0 \rangle = + \frac{1}{2p_K} \langle J/\Psi K^0 | \mathcal{T} | B_d^0 \rangle$$

$$\langle J/\Psi K_S | \mathcal{T} | \bar{B}_d^0 \rangle = - \frac{1}{2q_K} \langle J/\Psi \bar{K}^0 | \mathcal{T} | \bar{B}_d^0 \rangle$$

■ Tree level: $b \rightarrow c\bar{c}s \propto V_{cb}V_{cs}^*$

■ Kaon mixing: $\frac{q_K}{p_K} = \frac{V_{cs}^* V_{cd}}{V_{cs} V_{cd}^*} e^{-i2\chi'}$

■ $\delta = 0$ in $B_d^0 - \bar{B}_d^0$: $\frac{q}{p} \frac{M_{12}^*}{|M_{12}|} = \frac{(V_{tb}^* V_{td})^2}{|V_{tb}^* V_{td}|^2}$

■ Then:

$$\lambda_{J/\Psi K_S} = \frac{q}{p} \frac{\bar{A}_{J/\Psi K_S}}{A_{J/\Psi K_S}} = - \frac{(V_{tb}^* V_{td})^2}{|V_{tb}^* V_{td}|^2} \frac{(V_{cb} V_{cs}^*)^2}{|V_{cb} V_{cs}^*|^2} \frac{(V_{cs} V_{cd}^*)^2}{|V_{cs} V_{cd}^*|^2} e^{i2\chi'} = - e^{-i2(\beta - \chi')} \simeq - e^{-i2\beta}$$

$$C_{J/\Psi K_S} = 0 \quad S_{J/\Psi K_S} = \sin 2\beta \quad R_{J/\Psi K_S} = -\cos 2\beta$$

$\sin 2\beta$

- Time dependent rates

$$|\langle J/\Psi K_S | \mathcal{T} | B_d^0(t) \rangle|^2 = e^{-\Gamma t} \Gamma_{J/\Psi K_S} \{1 - \sin 2\beta \sin(\Delta M t)\}$$

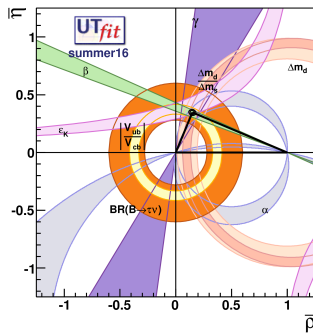
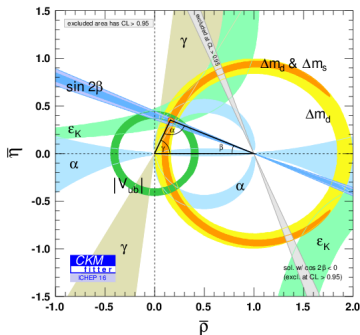
$$|\langle J/\Psi K_S | \mathcal{T} | \bar{B}_d^0(t) \rangle|^2 = e^{-\Gamma t} \Gamma_{J/\Psi K_S} \{1 + \sin 2\beta \sin(\Delta M t)\}$$

- CP Violating Asymmetry

$$\frac{|\langle J/\Psi K_S | \mathcal{T} | B_d^0(t) \rangle|^2 - |\langle J/\Psi K_S | \mathcal{T} | \bar{B}_d^0(t) \rangle|^2}{|\langle J/\Psi K_S | \mathcal{T} | B_d^0(t) \rangle|^2 + |\langle J/\Psi K_S | \mathcal{T} | \bar{B}_d^0(t) \rangle|^2} = -2 \sin 2\beta \sin(\Delta M t)$$

- Experimentally: $\sin 2\beta \simeq 0.68 \pm 0.02$

Remember CKM Fits?



More stuff

- Effective approach
- QCD effects, RGE
- New Physics sensitivity of flavour processes