Particle identification

G.Unal (CERN)

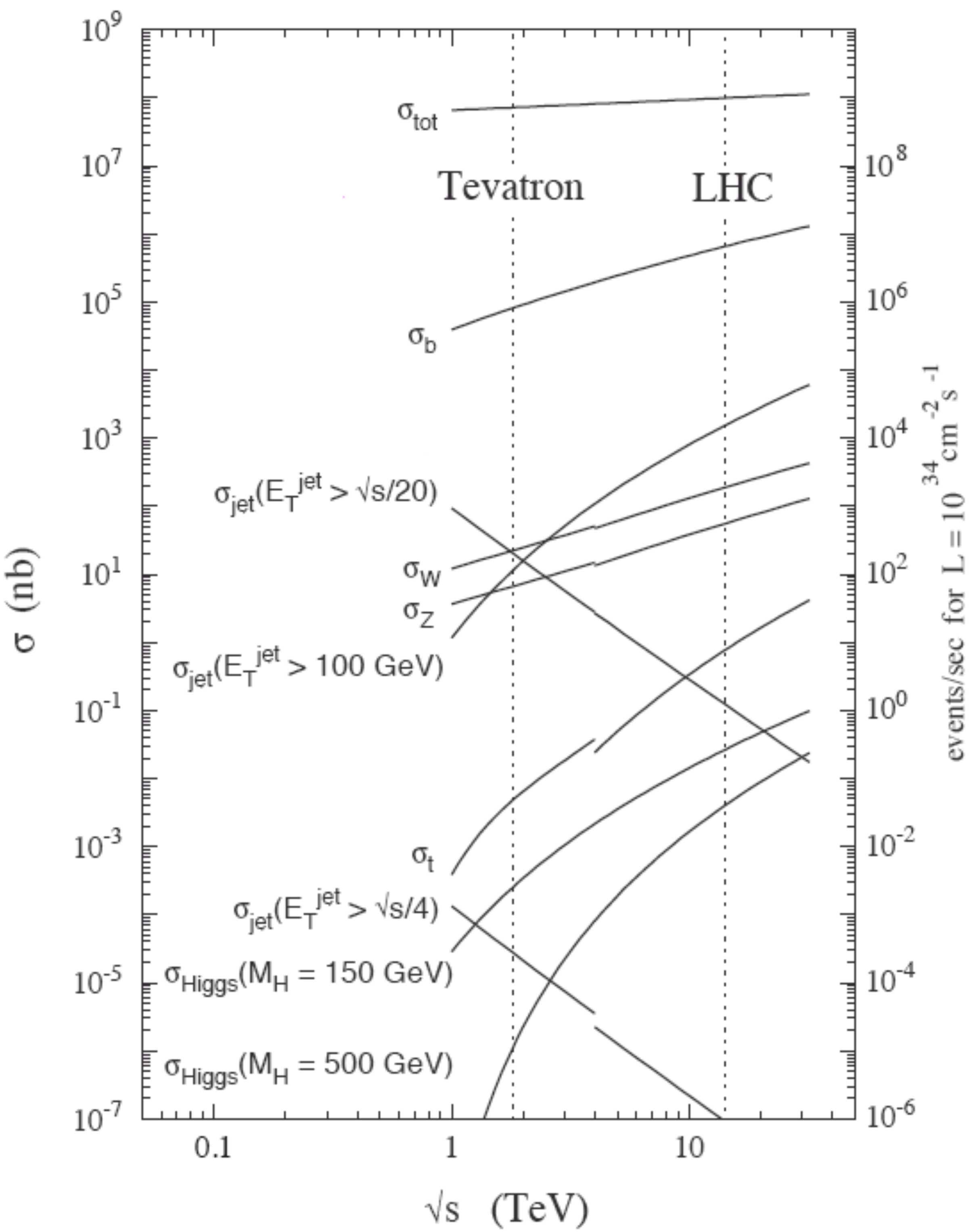
European Scientific Institute

European School of Instrumentation in Particle & Astroparticle Physics

Why particle identification ?

- *Is particle X decaying to electrons or muons ? Which are the corresponding branching ratio ?*
    - Understand properties (couplings) of this particle
- *Use particle Identification to separate signal and backgrounds*
    - To search for H->gamma gamma at LHC identify photons in the final state
- *Use particle Identification to optimize measurement of complicated final state*
    - «particle flow» event reconstruction in collider experiments

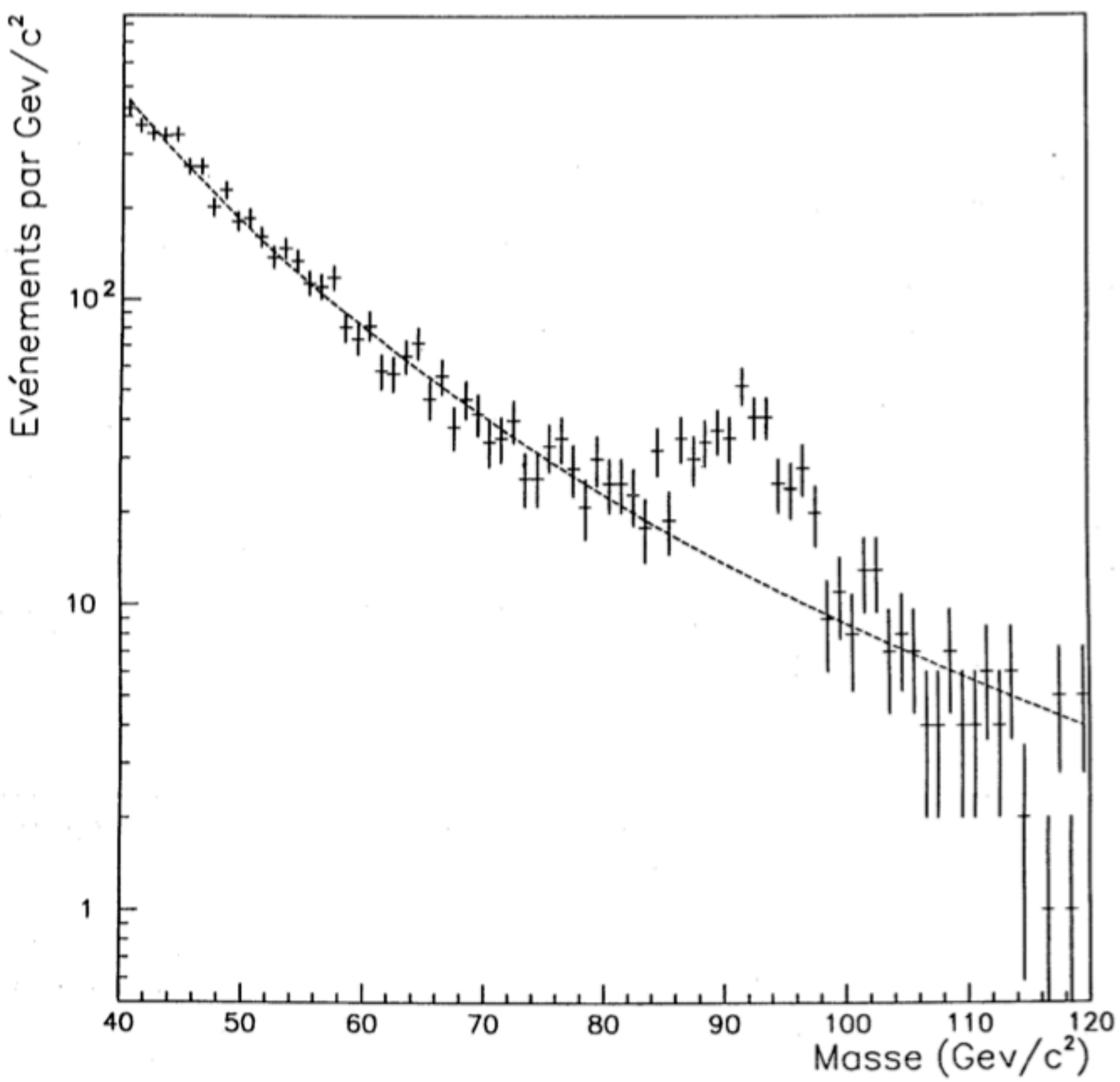
Cross-sections in hadron collider



High energy leptons give access to interesting physics processes

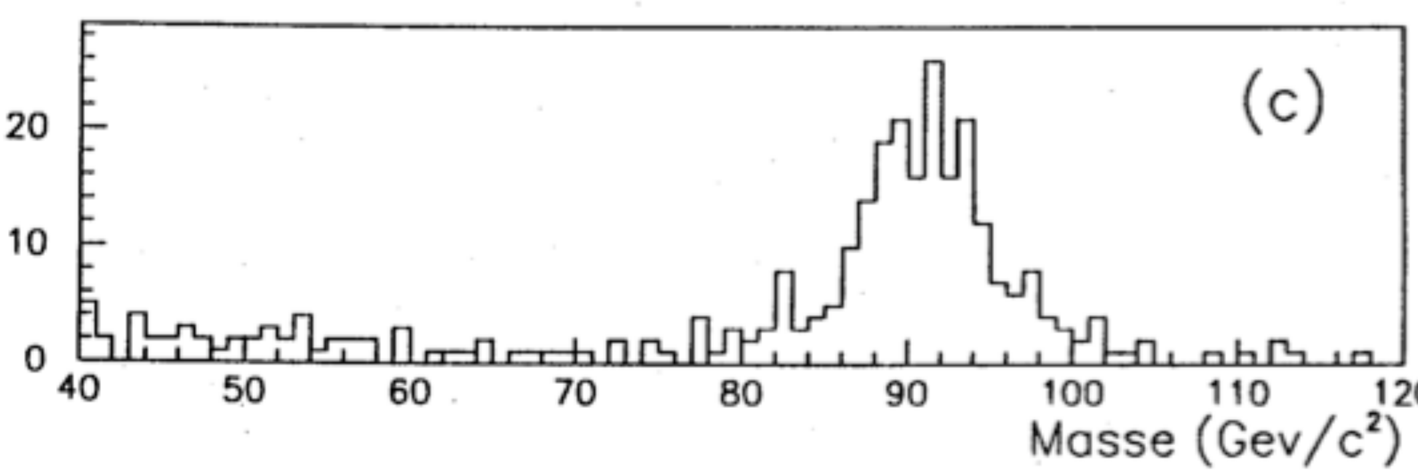
Some of these selections have to be done in real time (trigger) to reduce data rate to an acceptable level

Example of Z->ee sample in UA2 experiment (1988-1990 data)

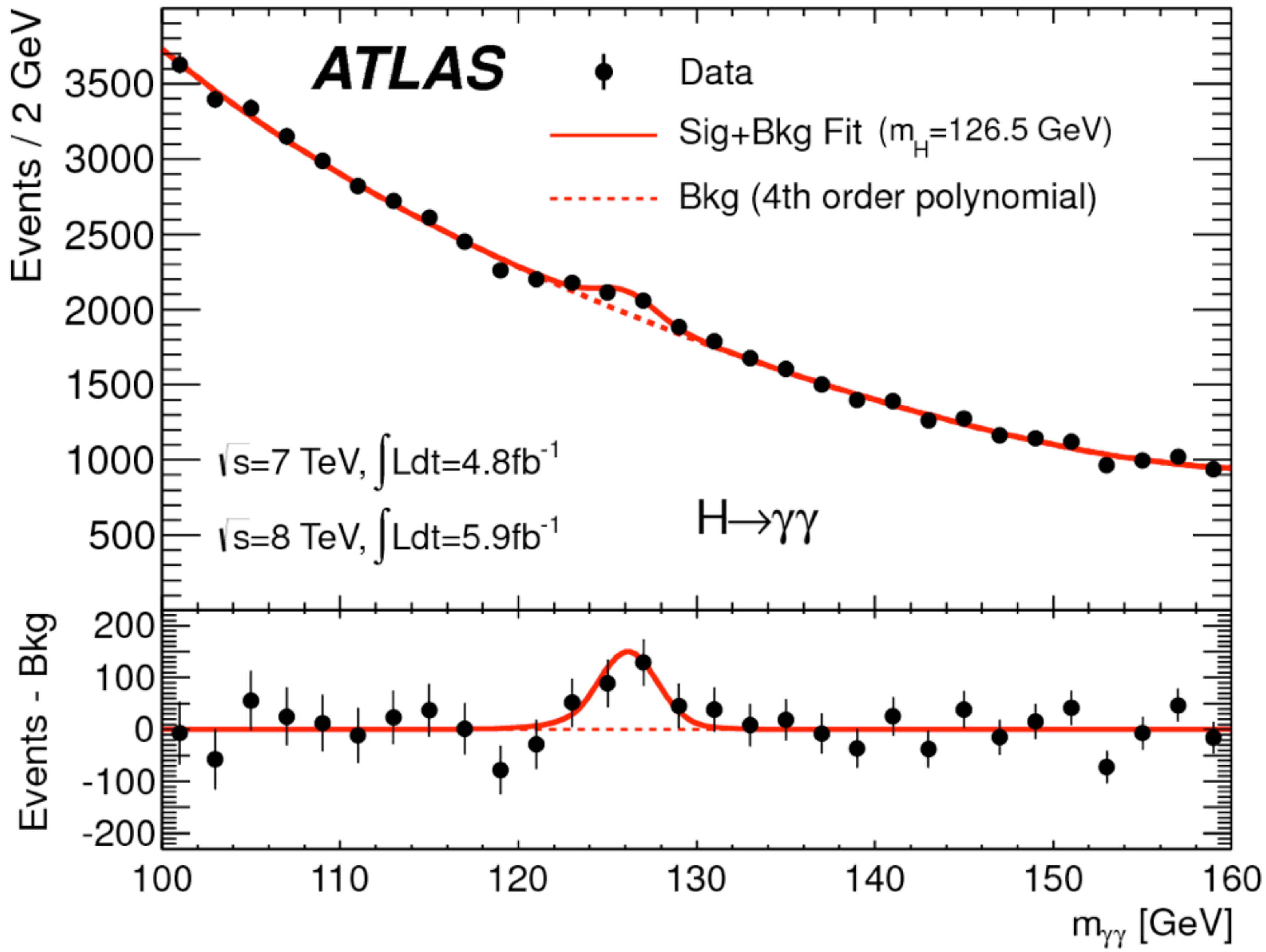


only calorimeter information
S/B ~1/1

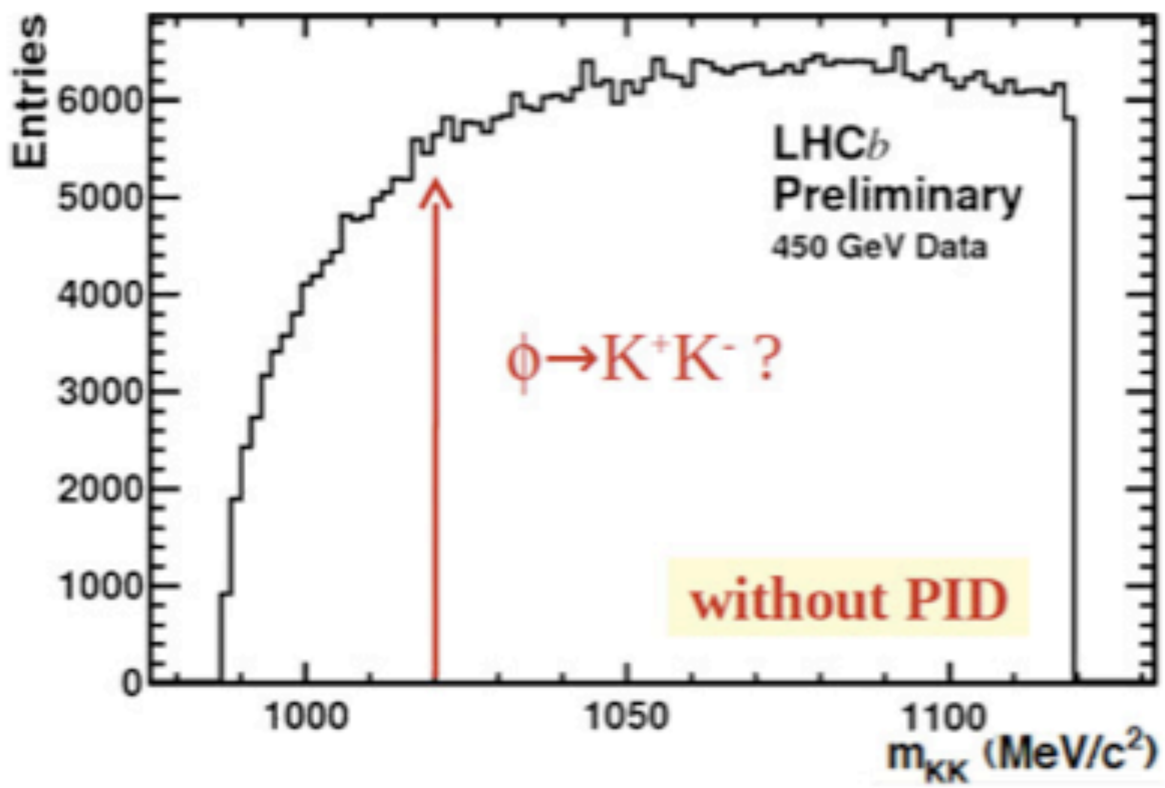
Adding matching to reconstructed track for electron identification

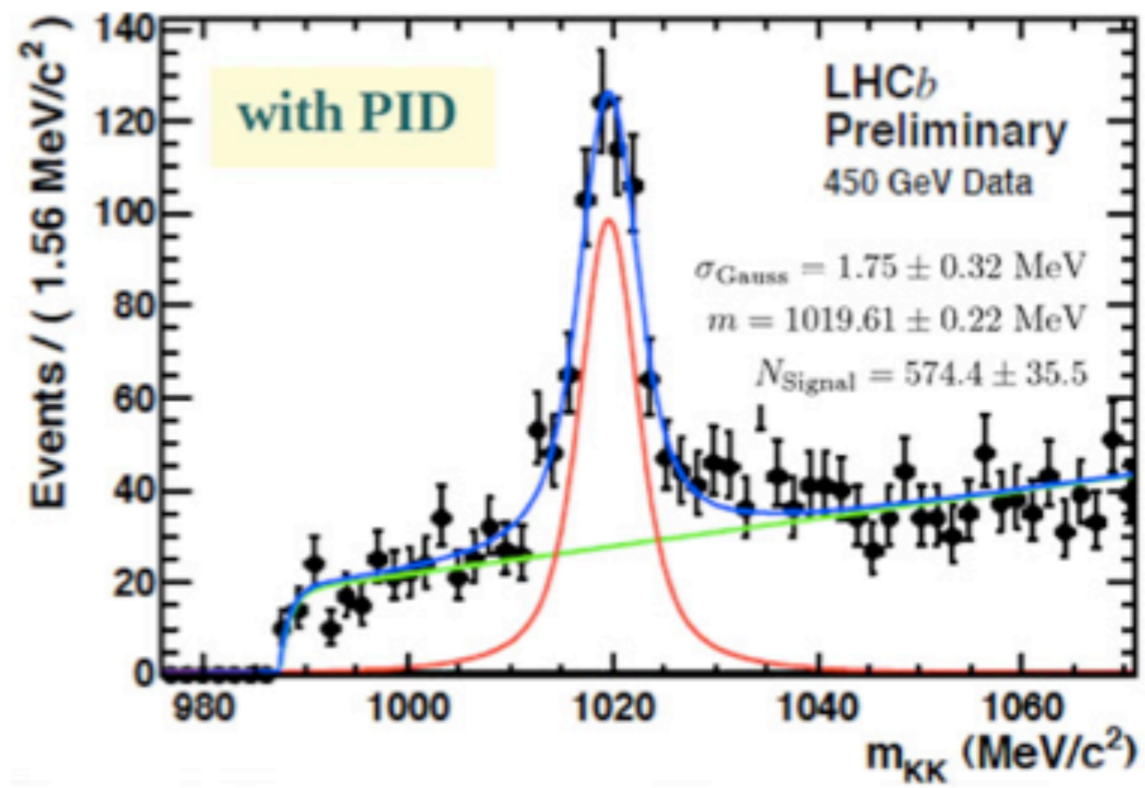


Detect Higgs boson through its decay to photons



Example of particle ID in flavor physics





many more examples where K/pion discrimination is important to study beauty and charm decays

Particle identification covers a wide range of techniques

- Exploit very different interaction of particles with matter
  (for instance calorimeter)
  - electron/photon/muon/hadron discrimination, neutrinos
- Measure mass of particle
  - Mass and charge enough to identify a particle
  - Once energy or momentum are measured, mass can be measured through measurement of beta (velocity) or gamma
  - mass from beta measurement works better a low energy
- Reconstruct decay of a particle to identify it
  - «identify» H by mass peak in H->gamma gamma
  - identify «long lived» particles by displaced decay vertex reconstruction

What is a «stable» particle ?

- Only few known particles are stable: photon, electron, proton, neutron(in nuclei), neutrinos
- Everything else decays but sometime are stable «enough» at the scale of the detector
- L = beta.gamma.c.tau
- Can a E=40 GeV muons (tau=2.2 10^{-6}s) in a collider experiment (size ~20m) be considered stable ?
- Can a E=1 GeV K0s (tau=8.9 10^{-11}s) in a LHC experiment be considered stable ? And a K0l (tau=5.10^{-8} s) ?
- In which cases can a charged pion (tau=2.6 10^{-8}s) be considered stable ? And a neutral pion ?
- Particle Identification depends on the experimental context and which particles are «directly» detected and which particle are «indirectly» detected (through their decay products)

What is a «stable» particle ?

- *Mean path length = beta.gamma.c.tau*

- Can a E=40 GeV muons (tau=2.2 10^{-6}s) in a collider experiment (size ~20m) be considered stable ?

  - *=> gamma ~380, L ~250 km*

- Can a E=1 GeV K0s (tau=8.9 10^{-11}s) in a LHC experiment be considered stable ? And a K0l (tau=5.10^{-8} s) ?

  - *=> gamma ~ 2, L ~5cm for K0s, L~30m (K0L)*

  - *Ks -> pi+pi- or pi0 pi0*

- In which cases can a charged pion (tau=2.6 10^{-8}s) be considered stable ? And a neutral pion ?

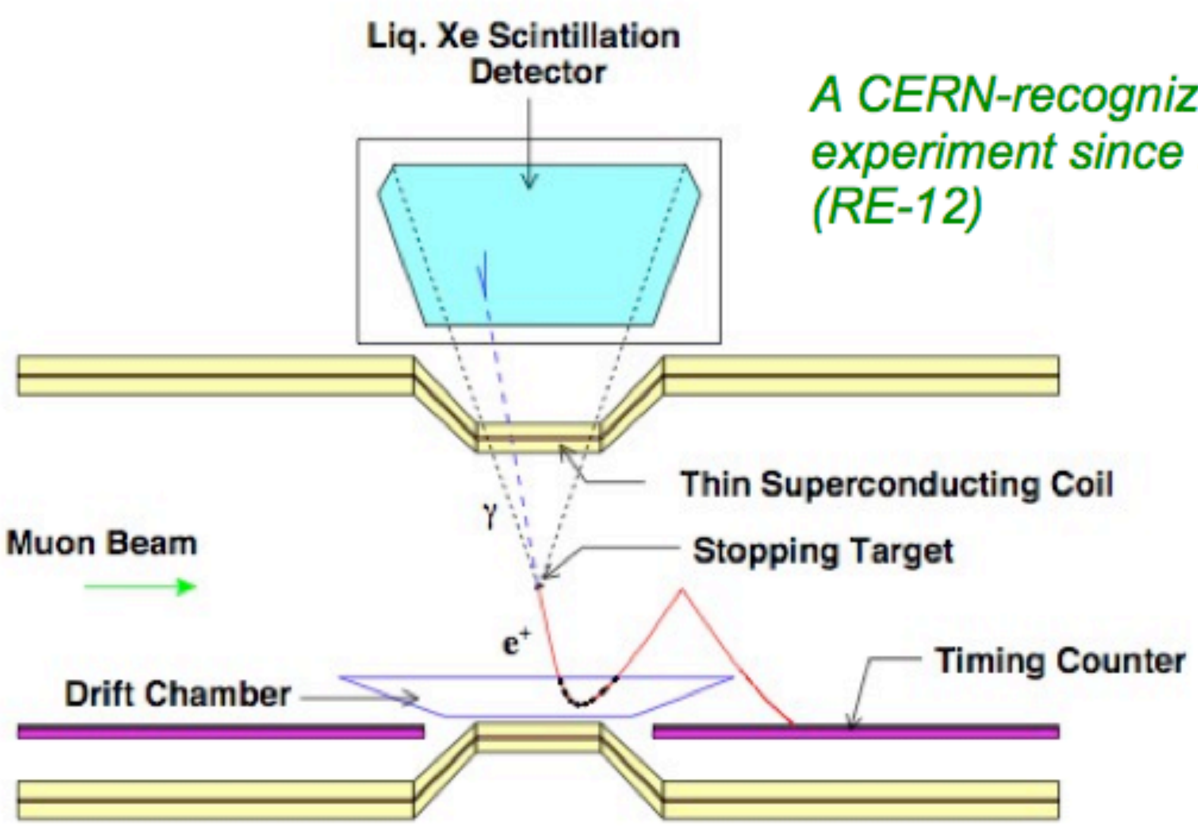
  - *L>~m if beta>~0.1 for charged pions. pi0 lifetime 8.10^{-17}s => ~never «stable»*

Example of an experiment looking for new ultra rare muon decay

The MEG experiment (arXiv:1303.2348)

- A search for $\mu \to e\,\gamma$ with the most intense DC muon beam of the world (3×10^7 μ/s @ PSI, Switzerland);
- Running since 2008.

A CERN-recognized experiment since 2005 (RE-12)



LXe calorimeter for photon detection

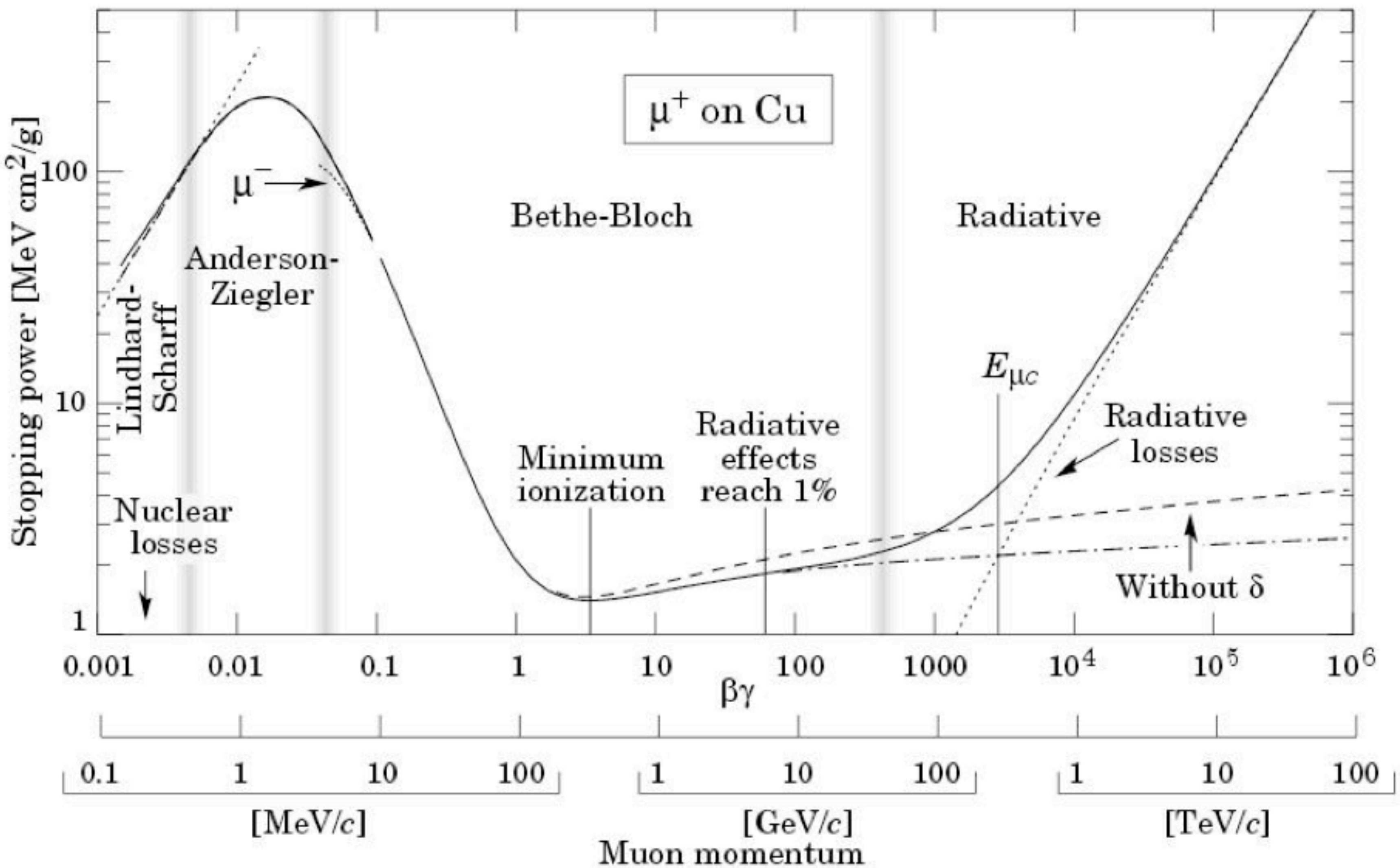
16 drift chambers for positron tracking

30 scintillating bars for positron timing and trigger (Timing Counter, TC)

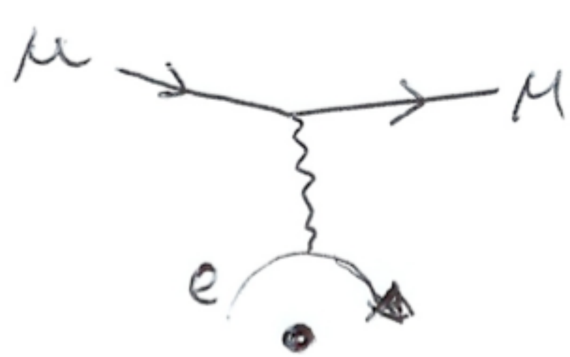
Exploiting different interactions with matter

- Mostly useful for e / muon / «hadron» discrimination
- In collider, high energy hadrons are not isolated but produced in «jets» from high energy quark and gluons
- Neutrinos are a special case

Muon energy loss

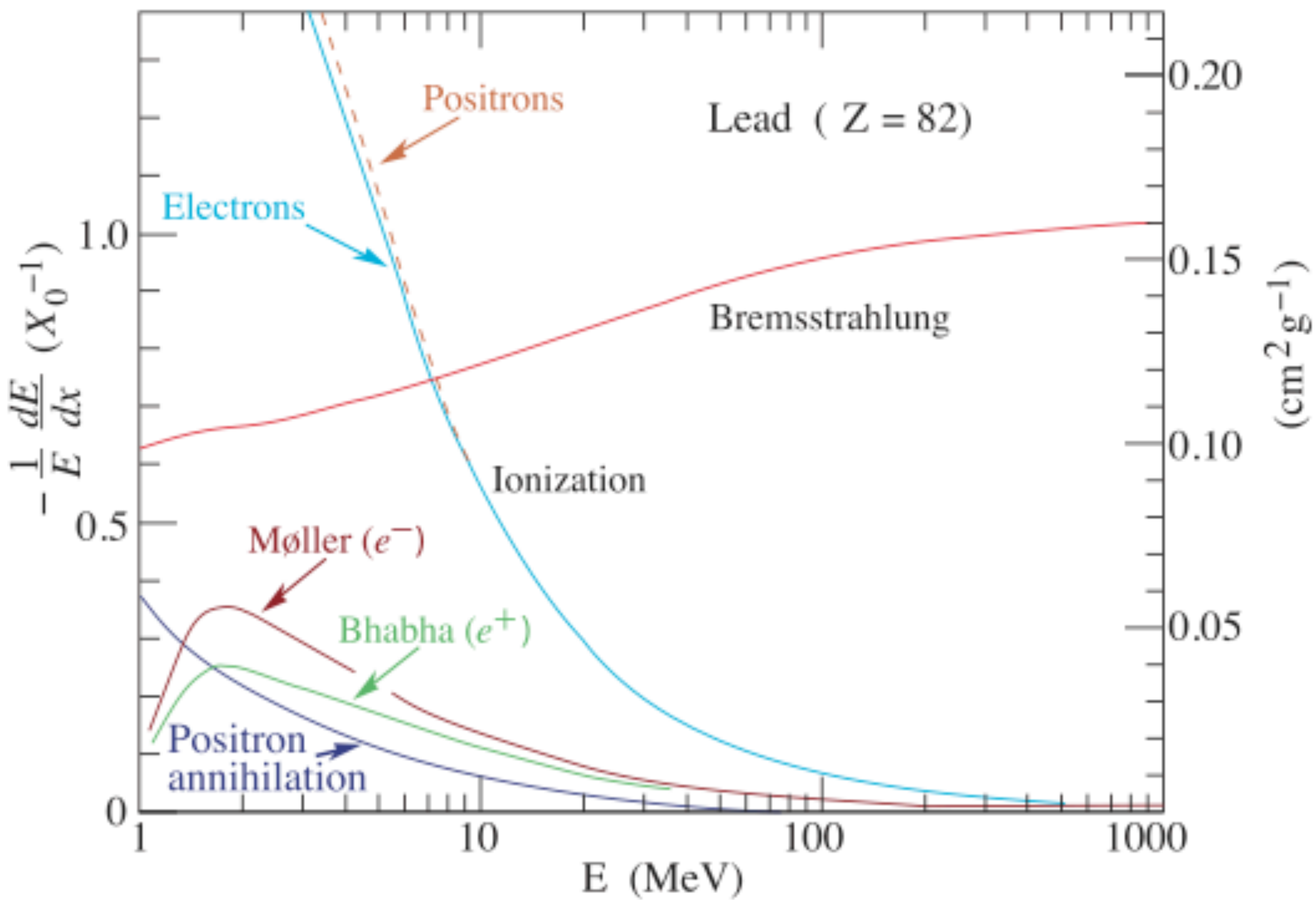


Ionization

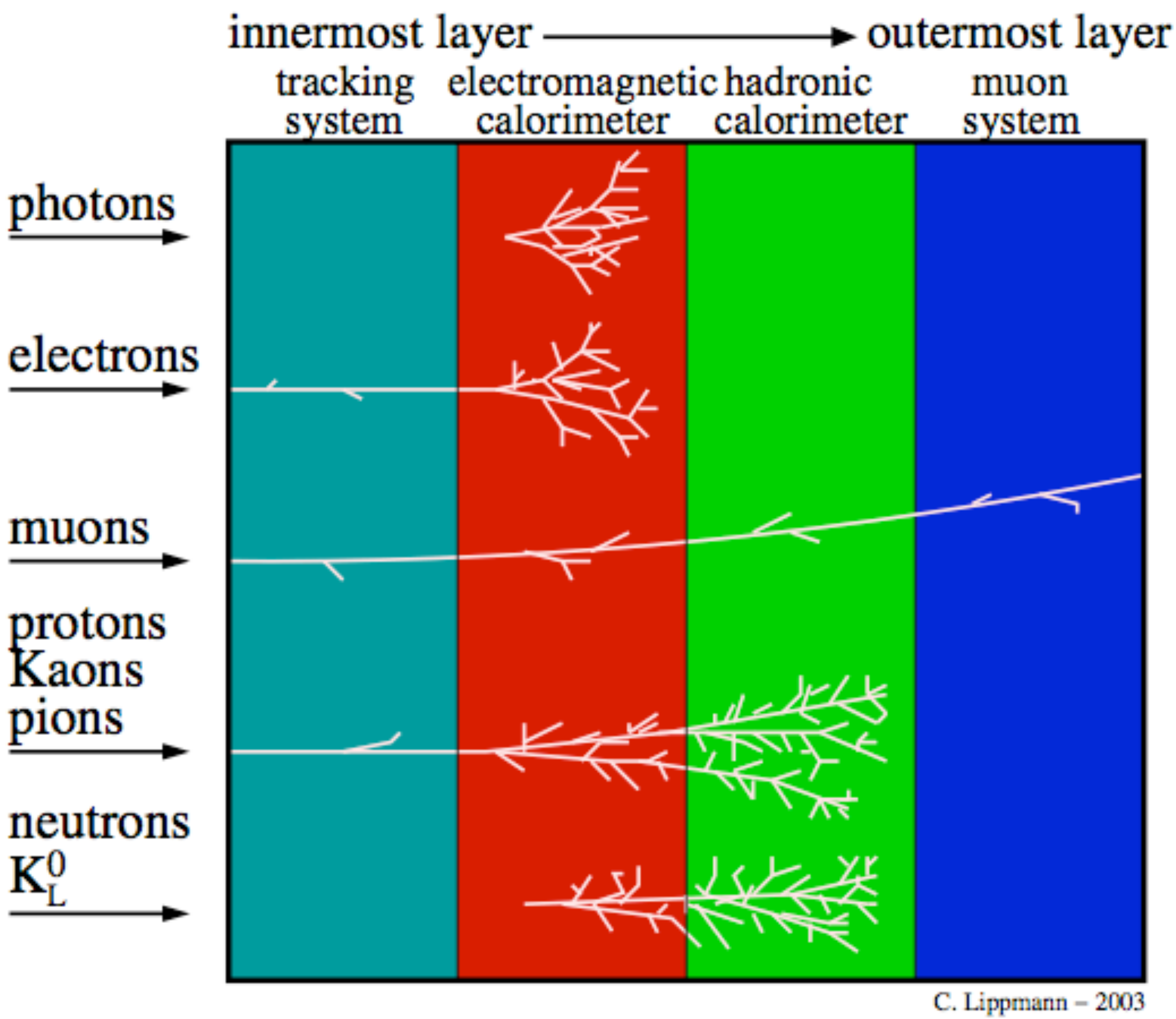


Bremsstrahlung

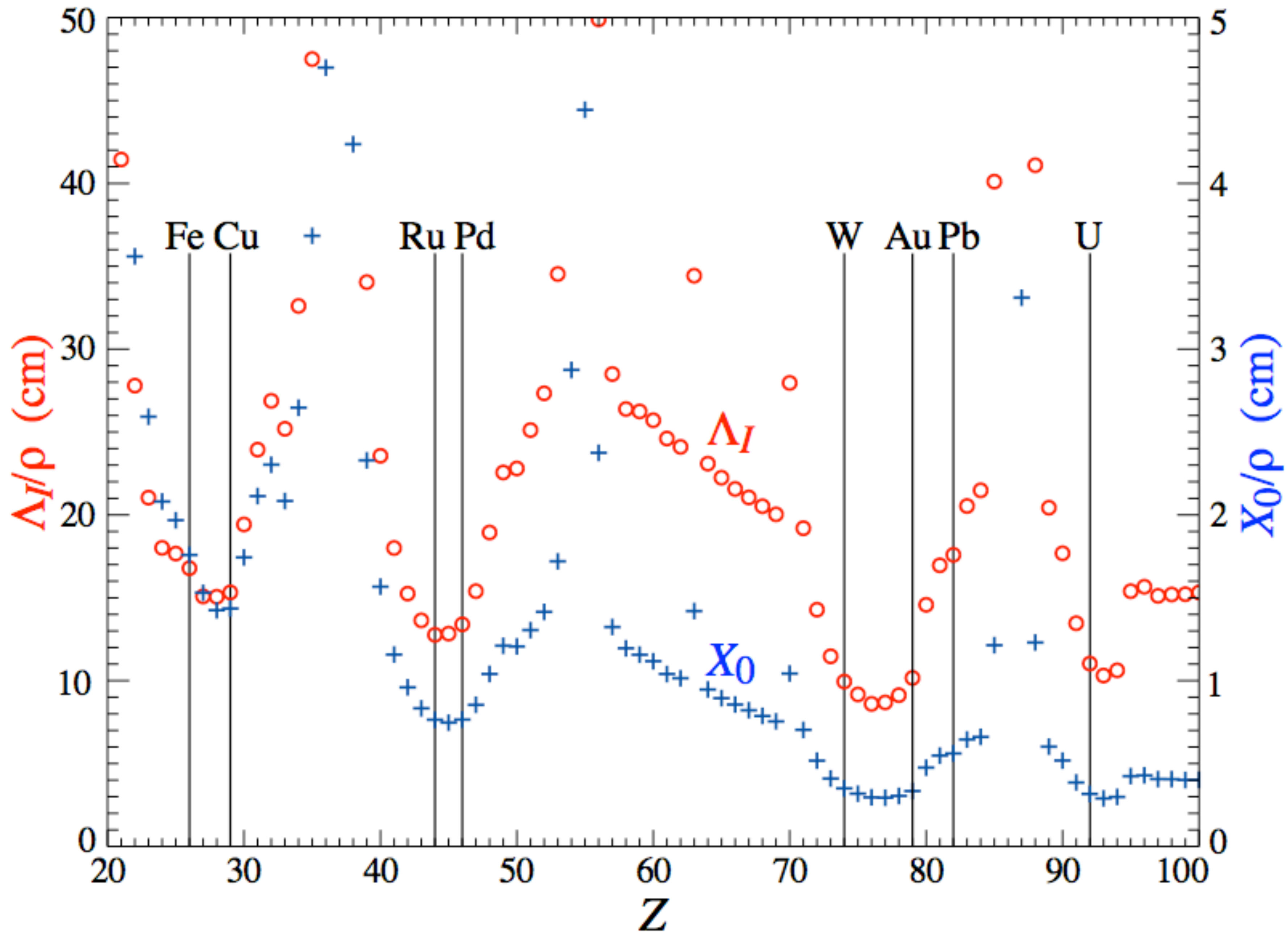
Electron energy loss



Sketch of particle interactions in detector

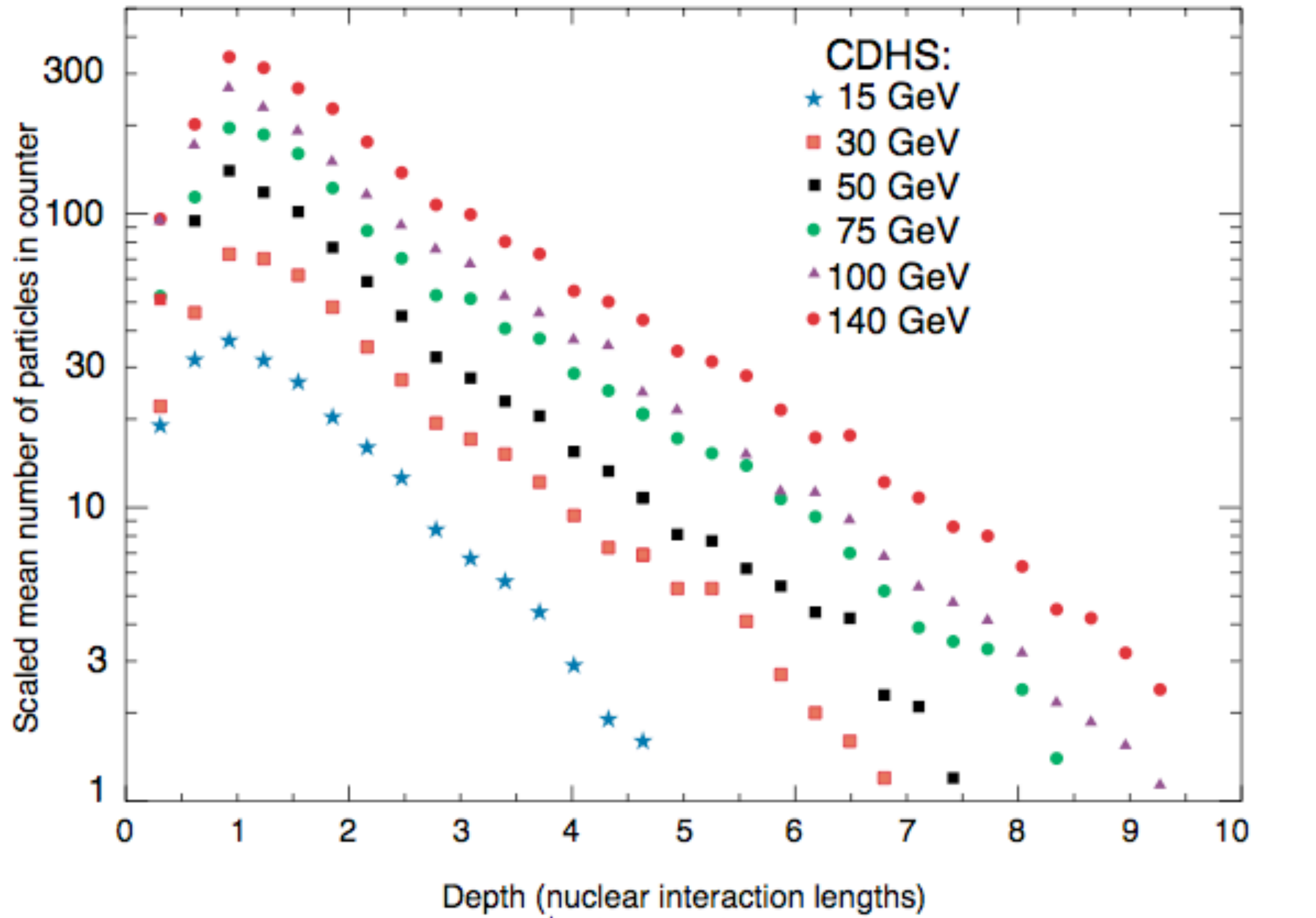


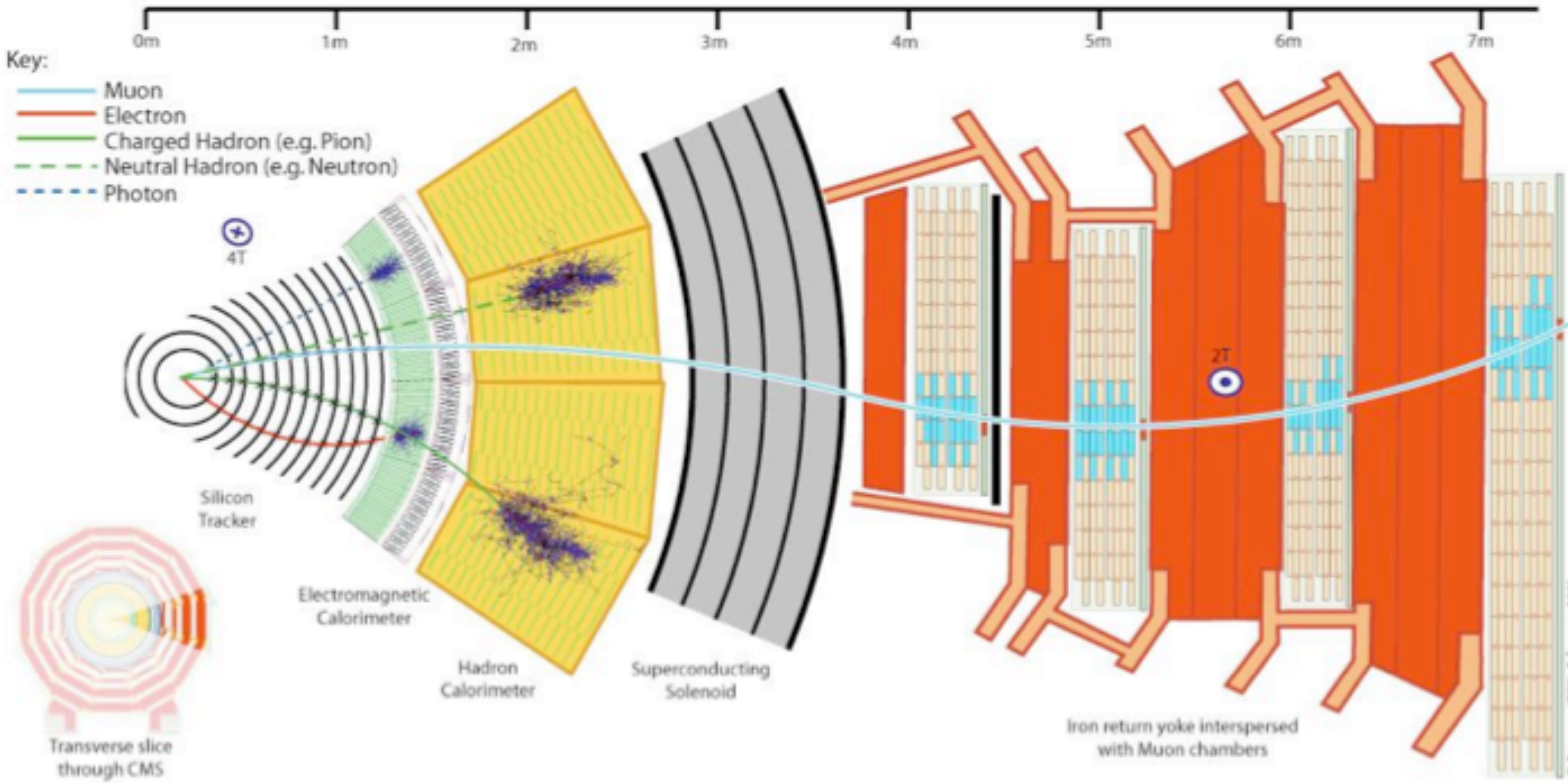
X0 = distance in which electron energy is reduced by 1/e by bremsstrahlung
Lambda_I = interaction length for hadronic interaction



$$X_0 = \frac{716.4 \cdot A}{Z(Z+1)\ln\frac{287}{\sqrt{Z}}}\ \mathrm{g \cdot cm^{-2}}$$[1].

How thick should a hadron calorimeter be ?

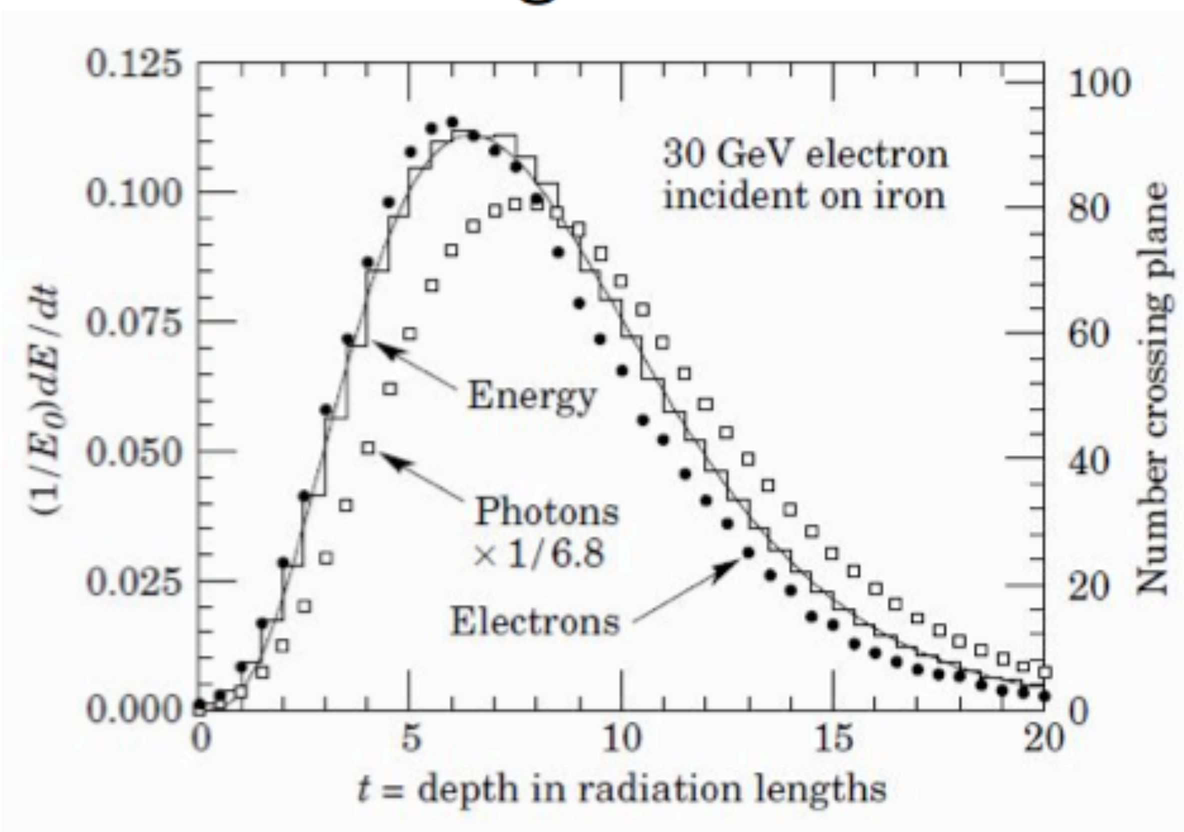




0m
1m
2m
3m
4m
5m
6m
7m
Key:
Muon
Electron
Charged Hadron (e.g. Pion)
Neutral Hadron (e.g. Neutron)
Photon
4T
2T
Silicon Tracker
Electromagnetic Calorimeter
Hadron Calorimeter
Superconducting Solenoid
Iron return yoke interspersed with Muon chambers
Transverse slice through CMS

Calorimeter showers initiated by e / photon

longitudinal



Difference electron-photon ?

Photon has to convert first
P(not convert) ~ exp(-7/9*x/x0)

lateral

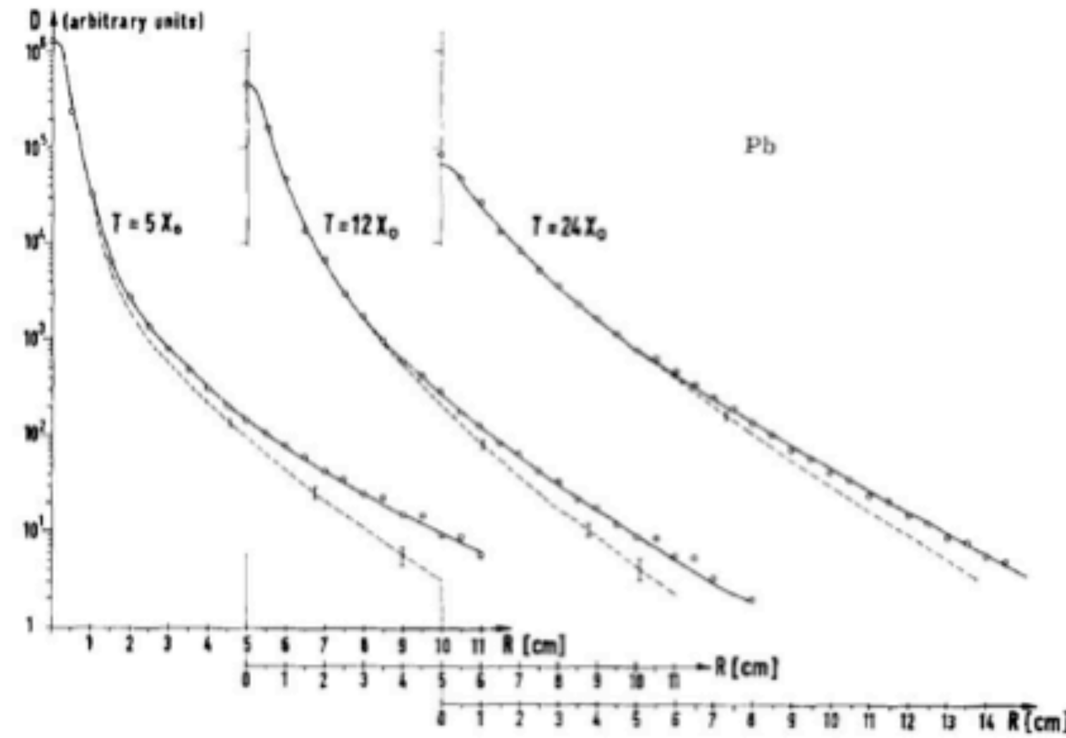


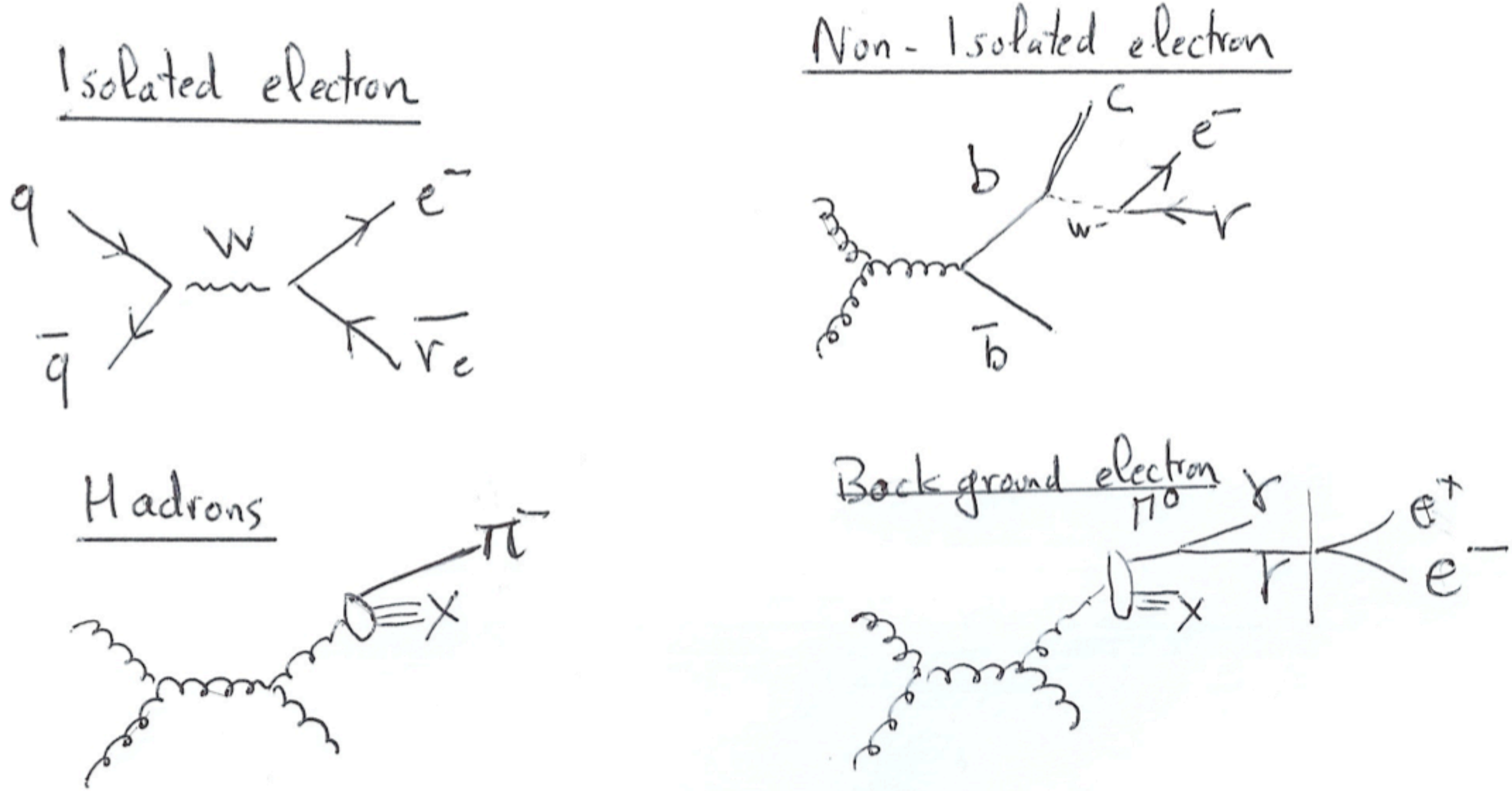
Fig. 4. Measured lateral distribution for lead (circles) in comparison with Monte-Carlo results (dotted line with error bars).

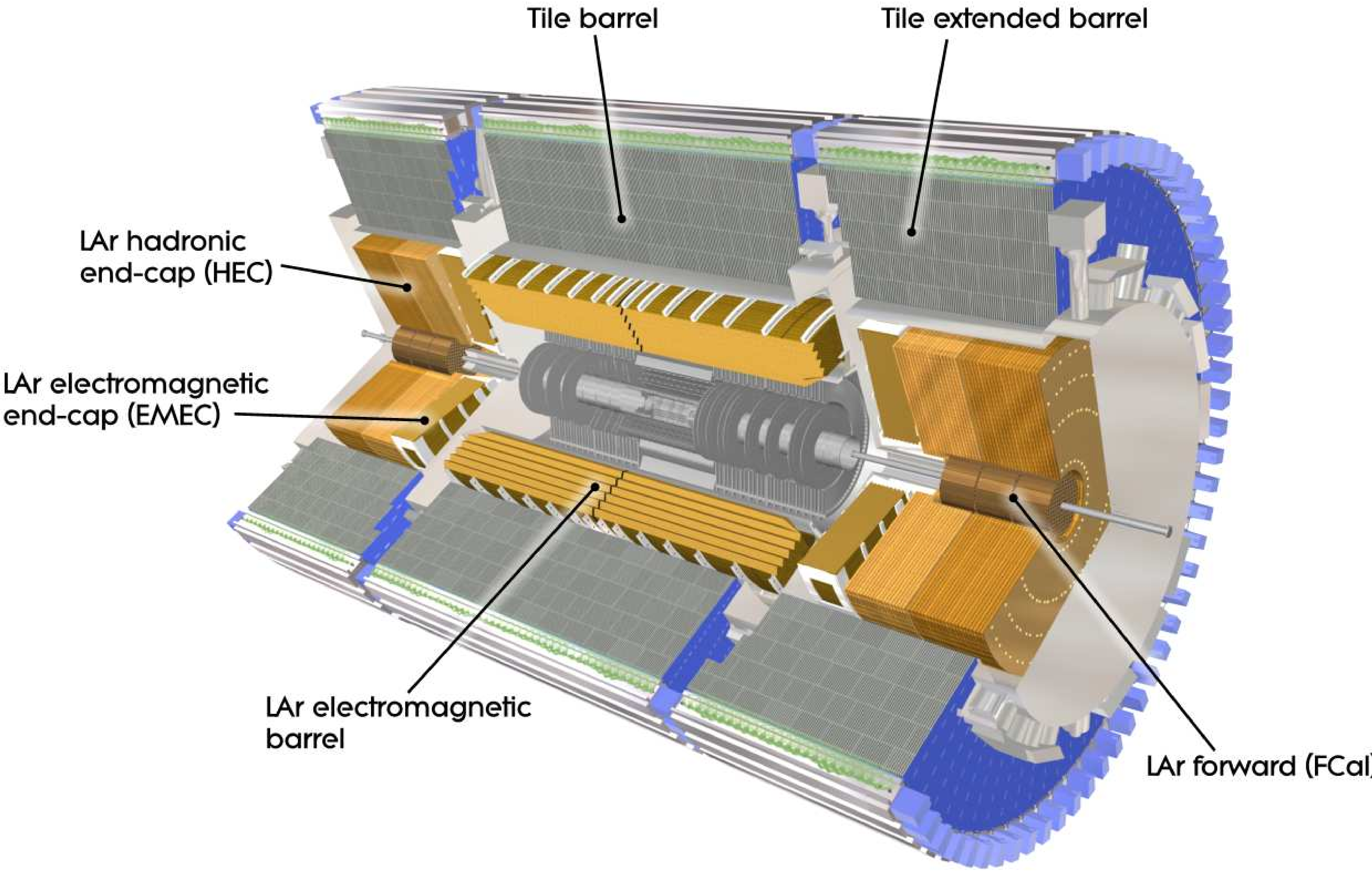
Moliere radius ~X0.(21MeV/Ec)
cylinder of ~2 Rm contains
~ 95% of energy

Electron identification in hadron colliders

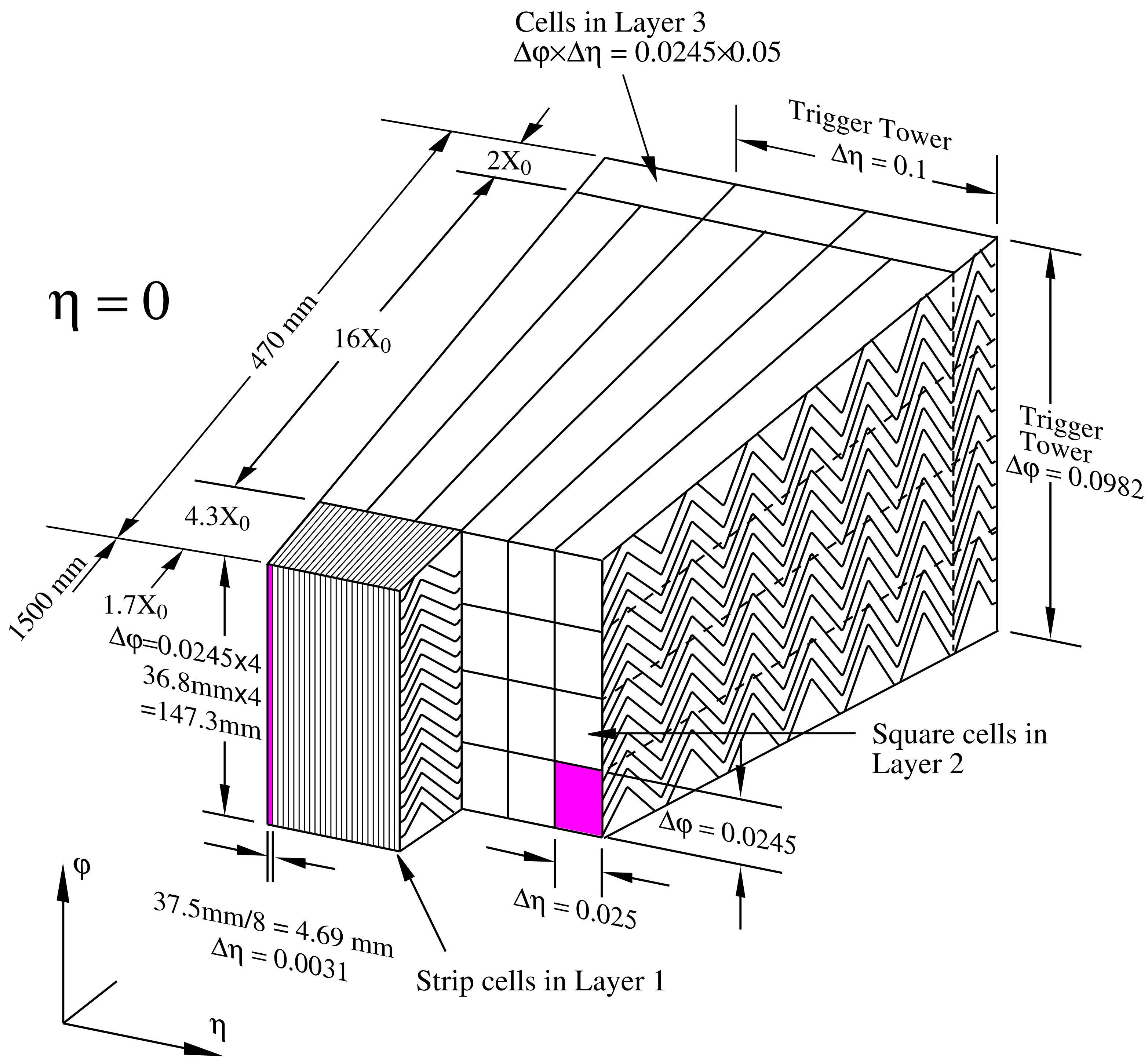
- High energy charged leptons are usually indication of «interesting» physics events, for instance decays of W or Z boson
- What are the backgrounds ?
- How to distinguish «good» electrons from them ?

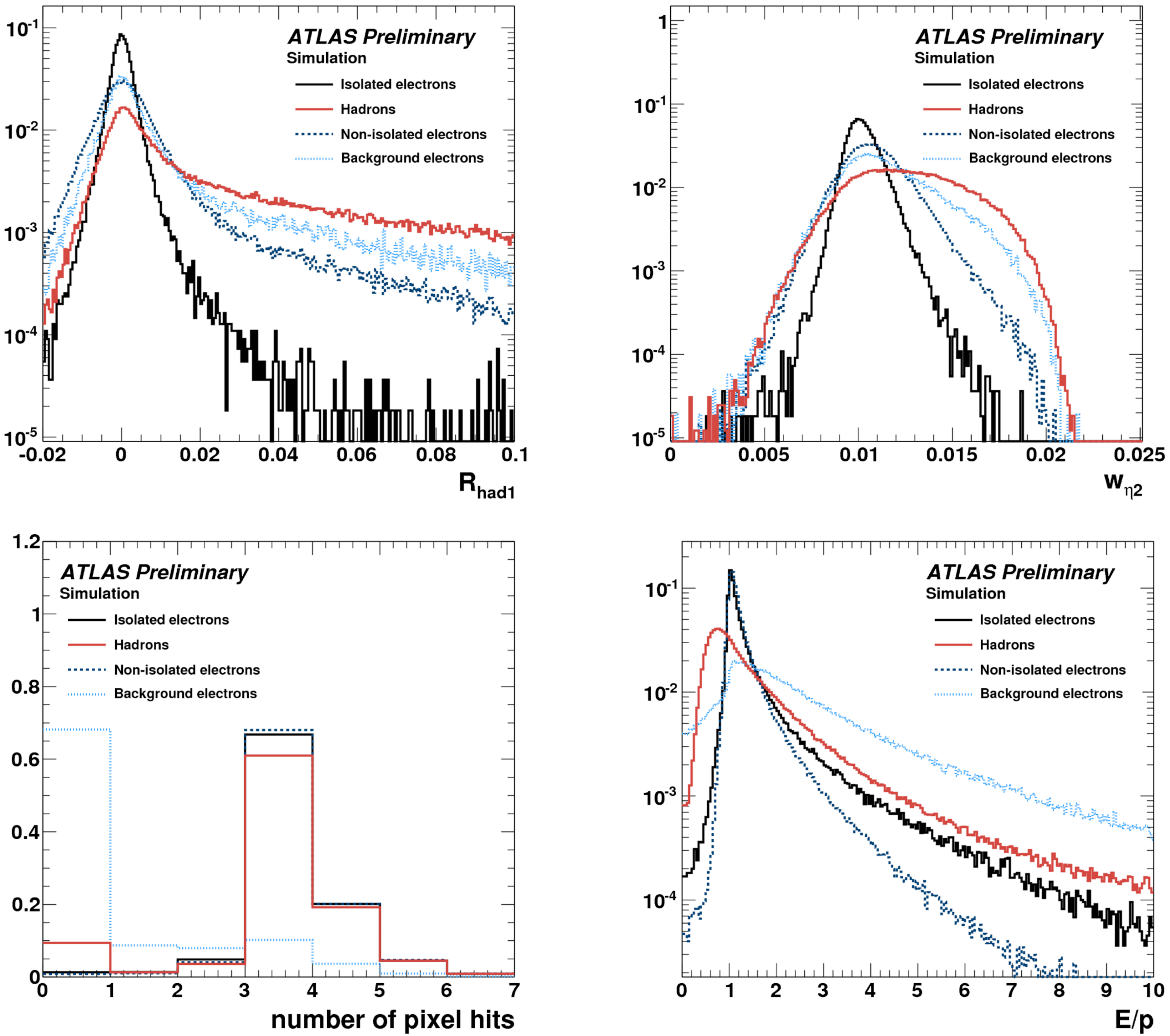
Description of different type of electron backgrounds

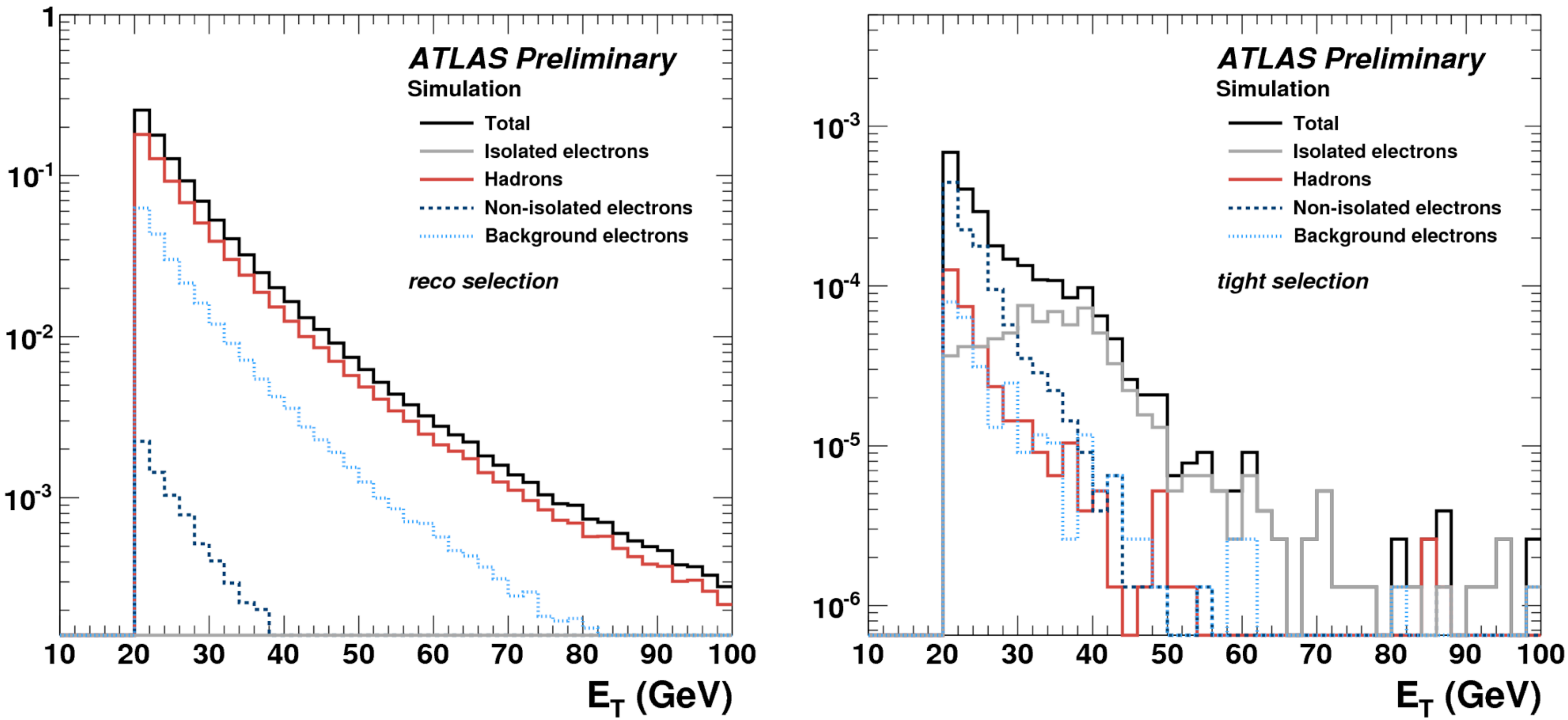



Tile barrel
Tile extended barrel
LAr hadronic end-cap (HEC)
LAr electromagnetic end-cap (EMEC)
LAr electromagnetic barrel
LAr forward (FCal)

Granularity of EM calorimeter to measure shower development



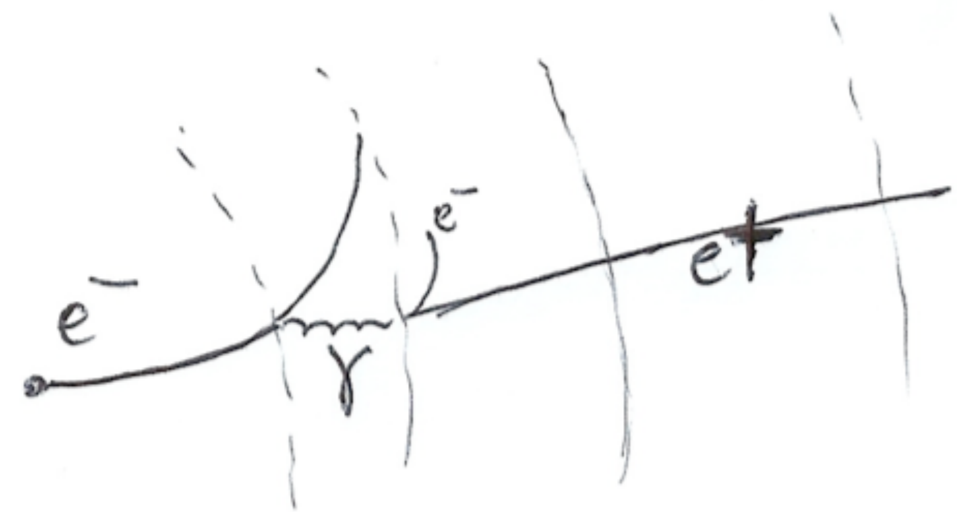

ATLAS Preliminary
Simulation
Isolated electrons
Hadrons
Non-isolated electrons
Background electrons
R_had1
w_η2
number of pixel hits
E/p

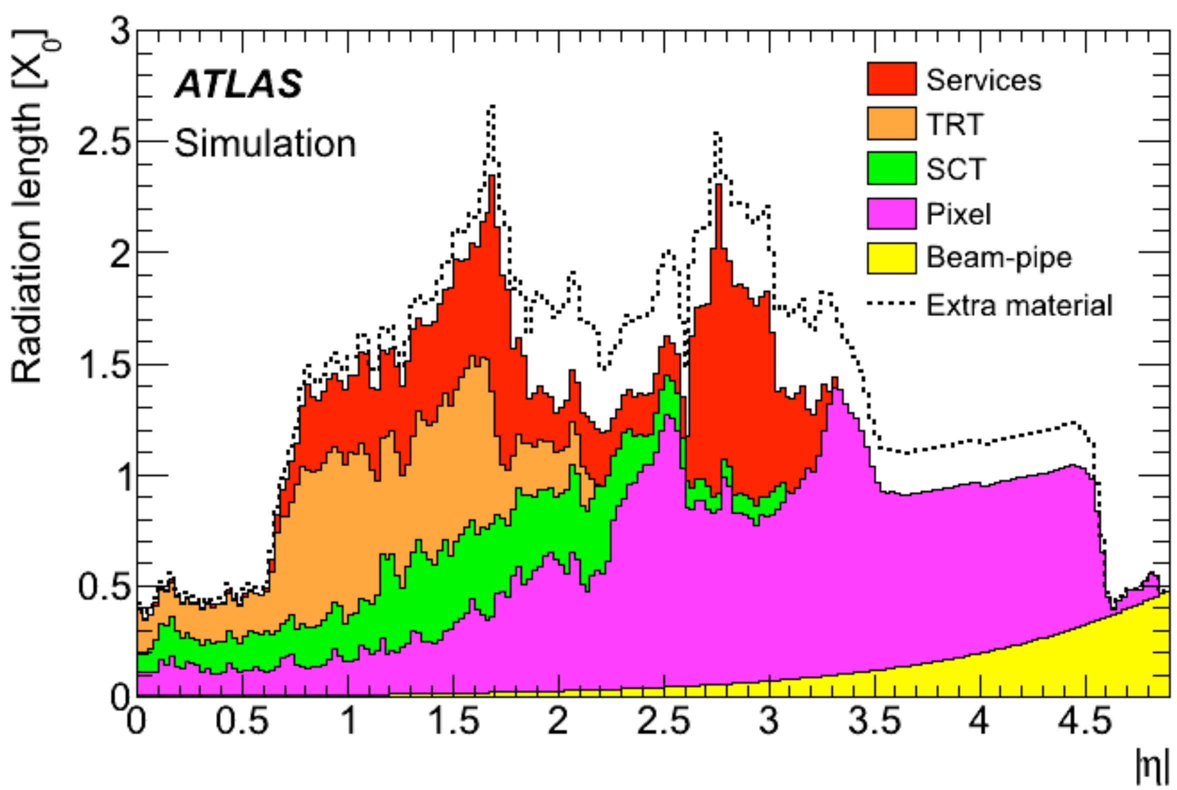


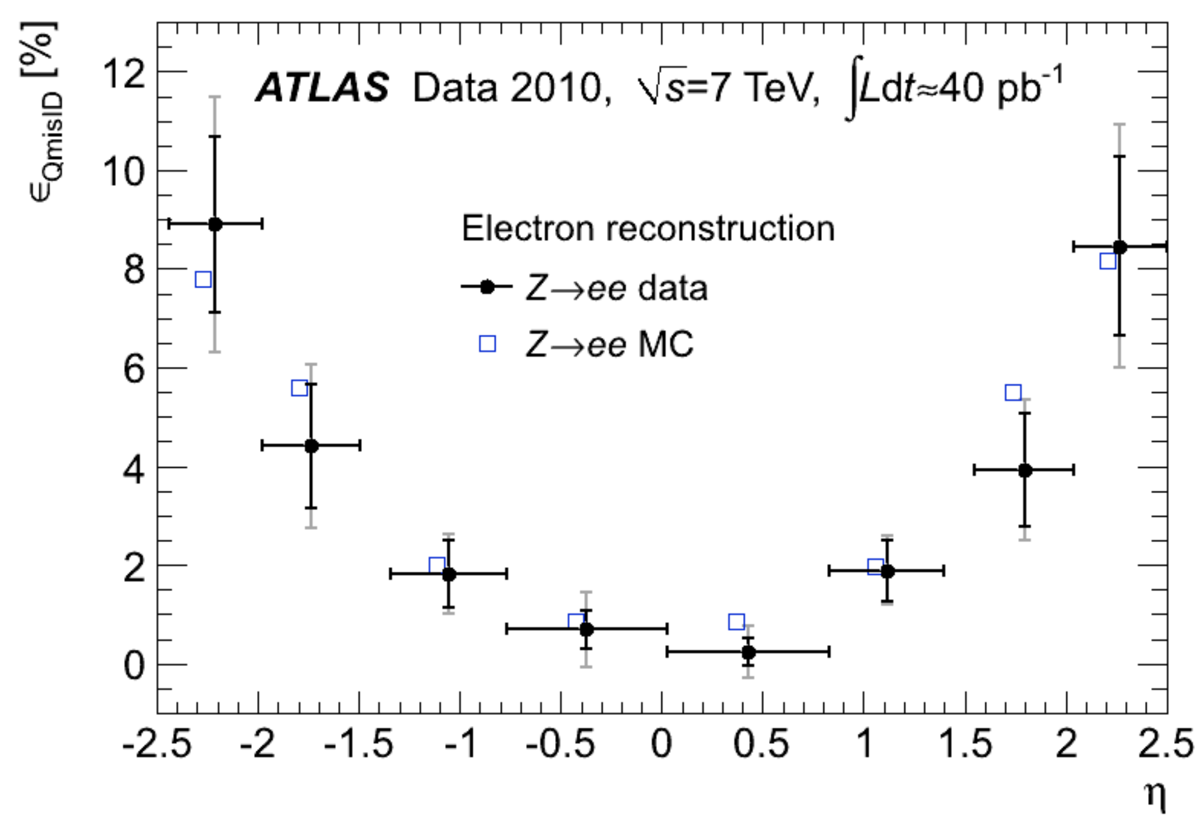
ATLAS Preliminary
Simulation
Total
Isolated electrons
Hadrons
Non-isolated electrons
Background electrons
reco selection
E_T (GeV)
ATLAS Preliminary
Simulation
Total
Isolated electrons
Hadrons
Non-isolated electrons
Background electrons
tight selection
E_T (GeV)

Electron charge

Calorimeter does not measure electron charge
Use track curvature in magnetic field for that
Main possibility of mistake for electron: Interaction with the inner detector material giving rise to bremstrahlung and conversions and not getting the «right» track



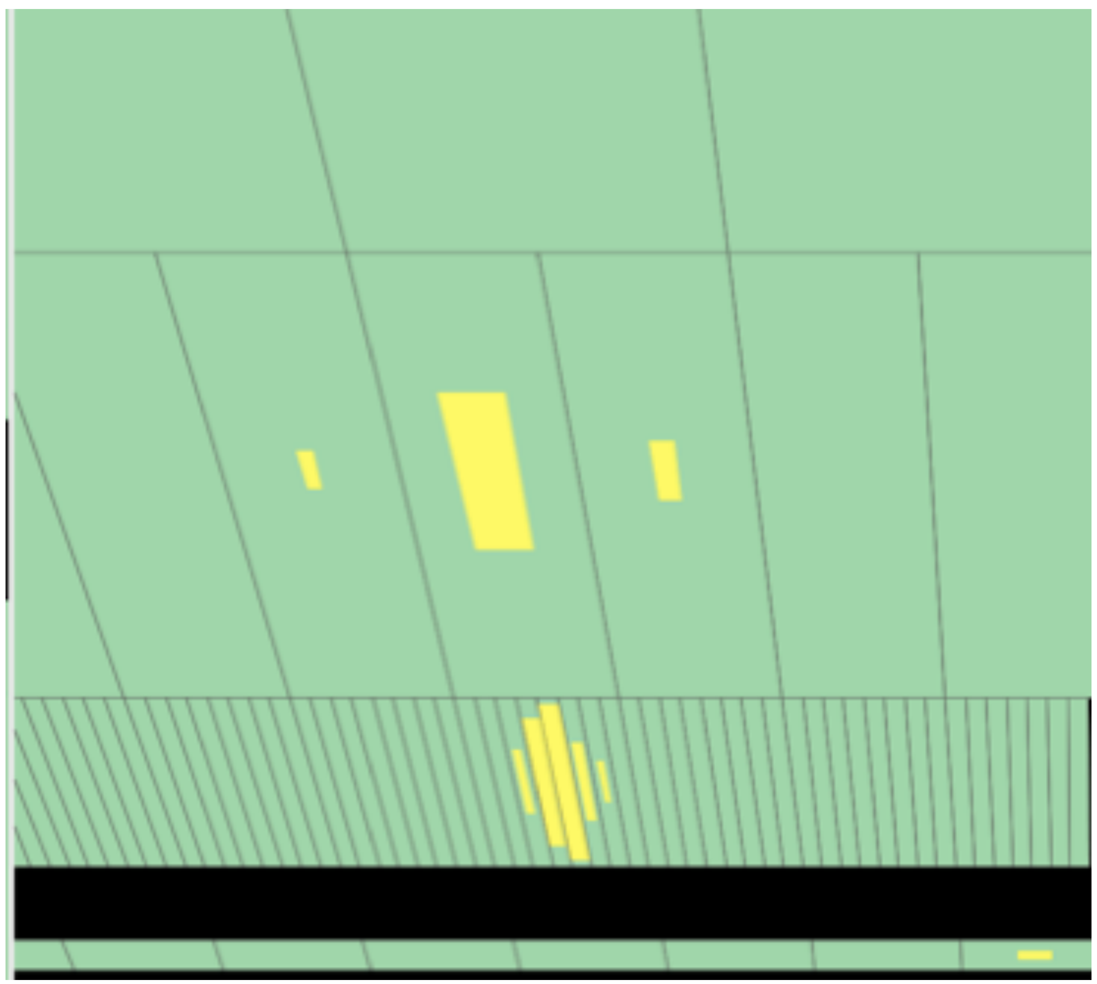


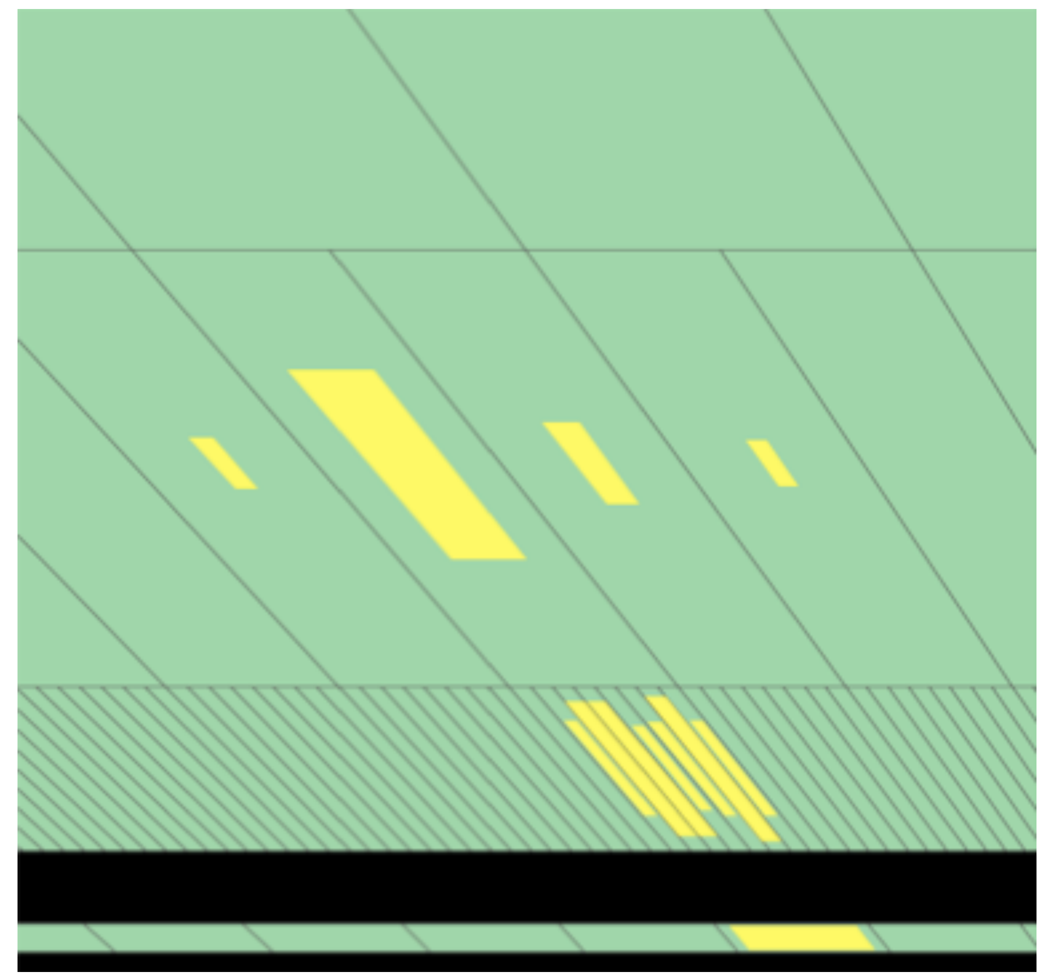


Photon identification in collider experiment

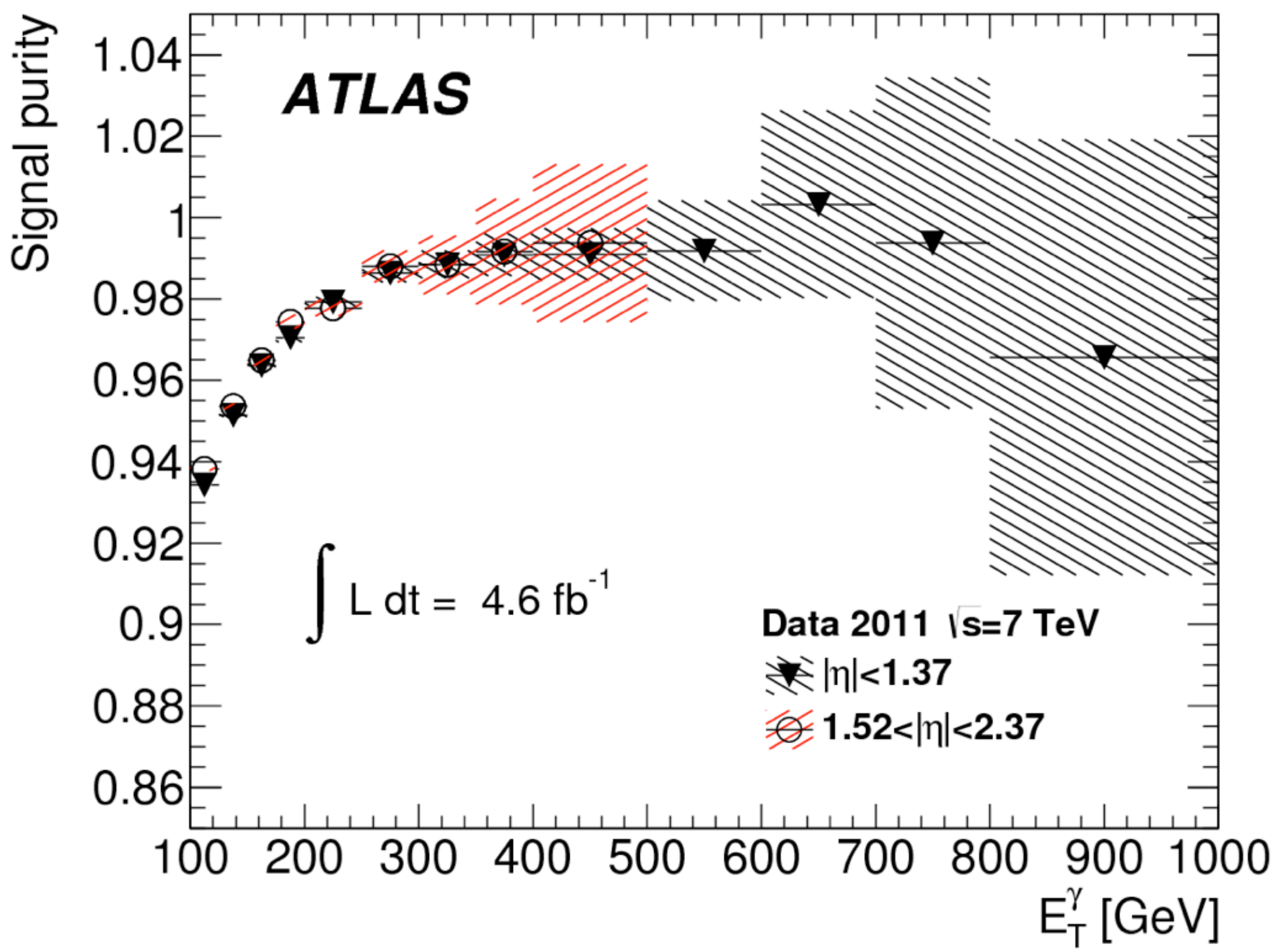
Background from high energy pi0->gamma gamma
What is the separation between the photons ?
What information can be exploited ?

theta_min ~2/gamma
~0.0067 at E=40 GeV
=> 1cm @ R=150cm

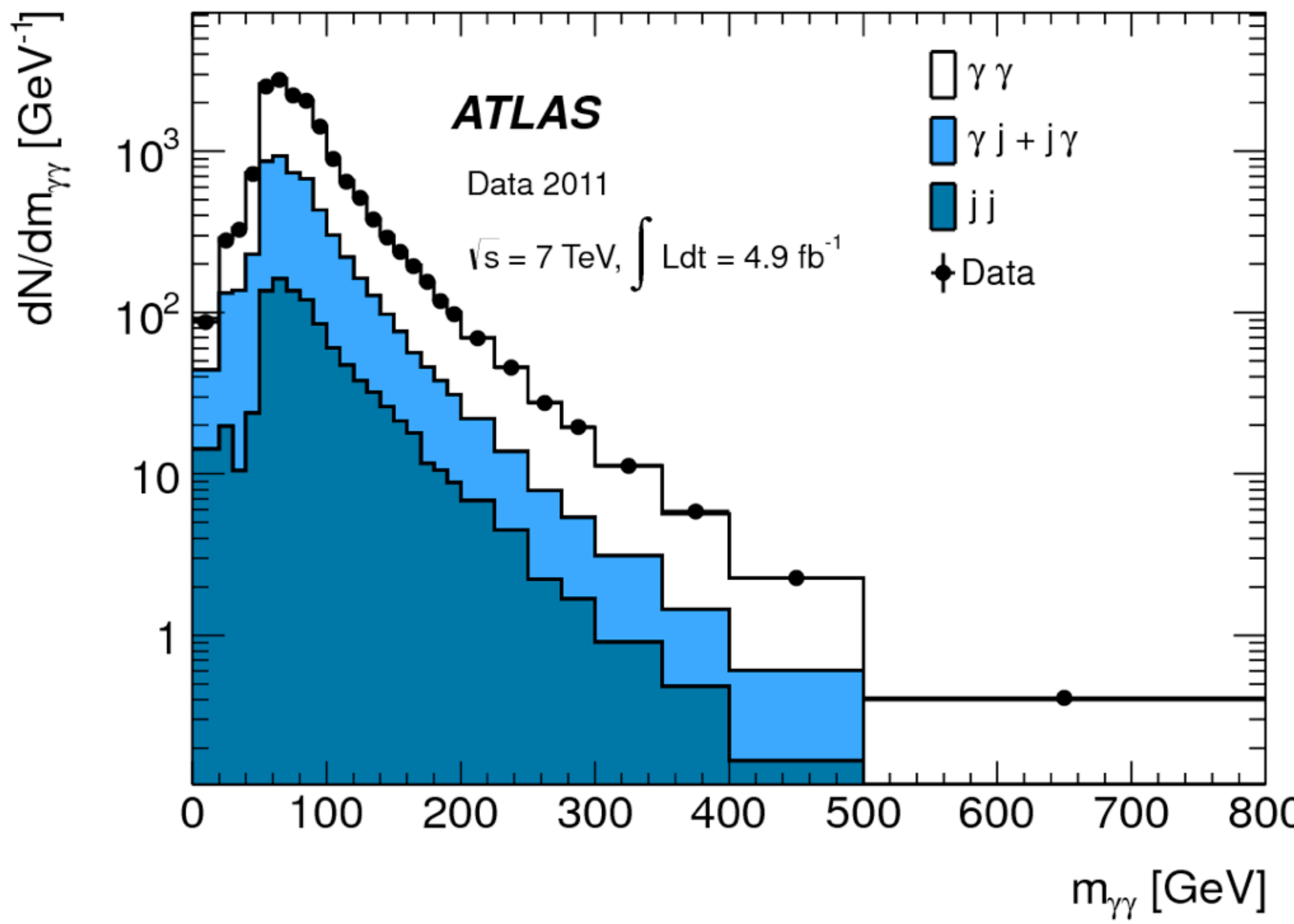




Example of photon identification performance in ATLAS



High energy inclusive photon
Purity >95%



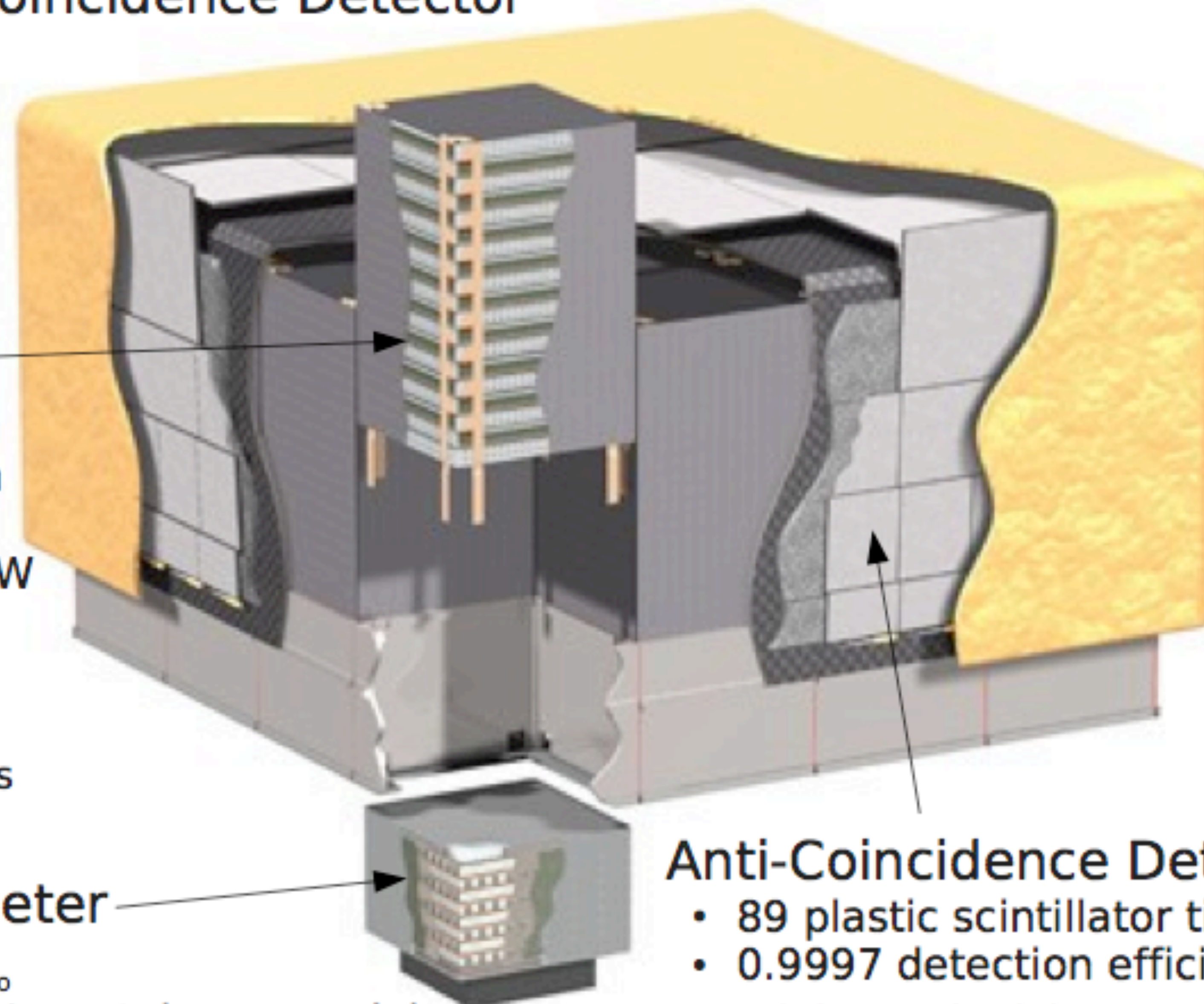
Di photon events at intermediate mass
Purity ~70-80%

Some of these techniques are also used in Space

- Fermi LAT : identify and measure ~50 MeV to ~300 GeV gamma rays with good angular resolution
- AMS : look for antimatter in space => particle identification and charge measurement

Fermi LAT

4x4 array of identical towers (tracker + calorimeter) surrounded by an Anti-Coincidence Detector



Tracker

- 18 layers (x-y) with silicon strip detectors + tungsten conversion foil
- 2 sections (depending on W thickness):
  - Thin (front) : 12x0.03X_o
  - Thick (back) : 4x0.18X_o
  - No W in the 2 bottom layers
- 1.4 X_o on axis

Calorimeter

- 8.6 X_o
- 96 CsI crystals per module

Anti-Coincidence Detector

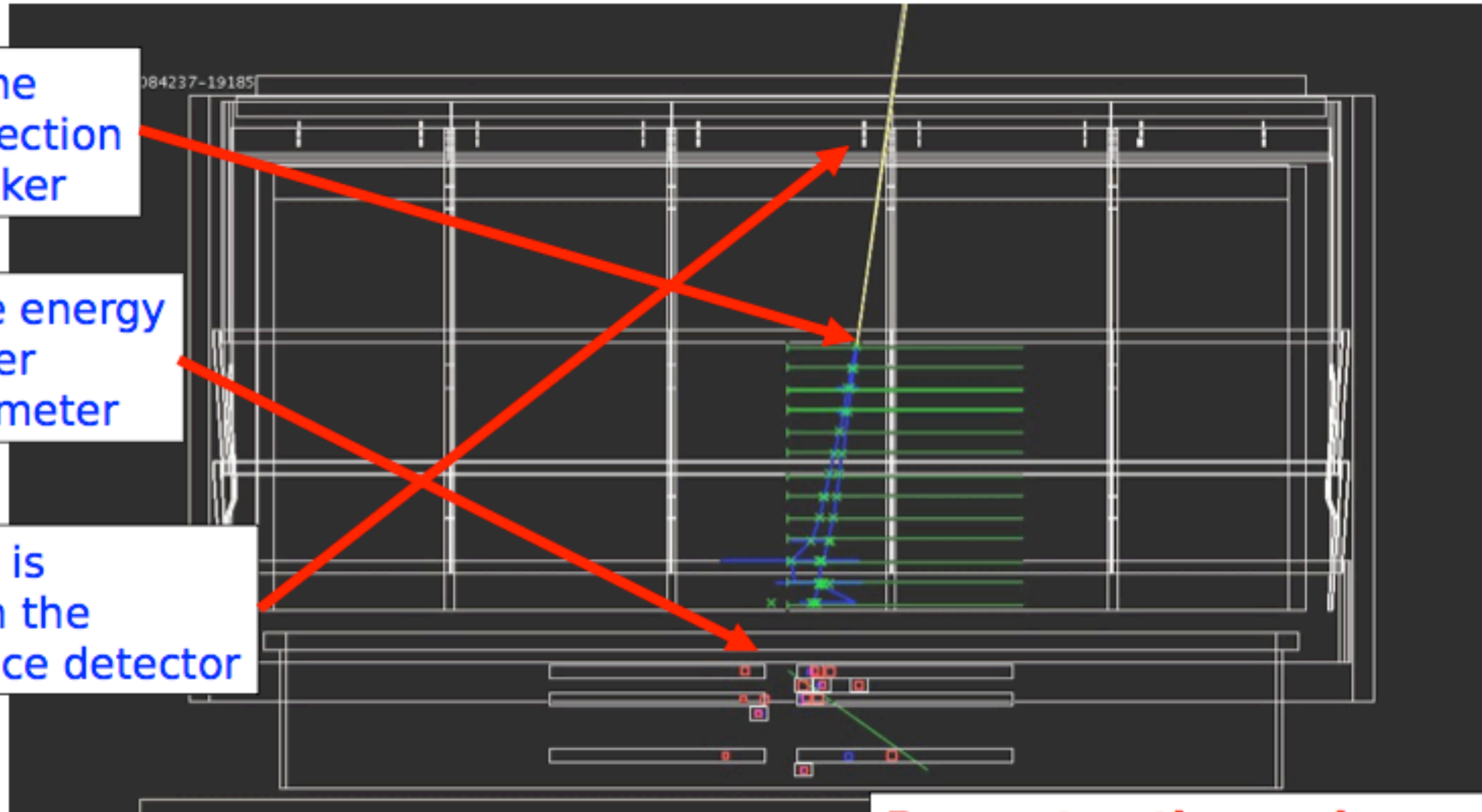
- 89 plastic scintillator tiles
- 0.9997 detection efficiency for minimum-ionizing particles

Ph. Bruel CALOR 2012, 4-8 June 2012 7

Determine the incoming direction with the tracker

Determine the energy with the tracker and the calorimeter

Check if there is some signal in the anti-coincidence detector



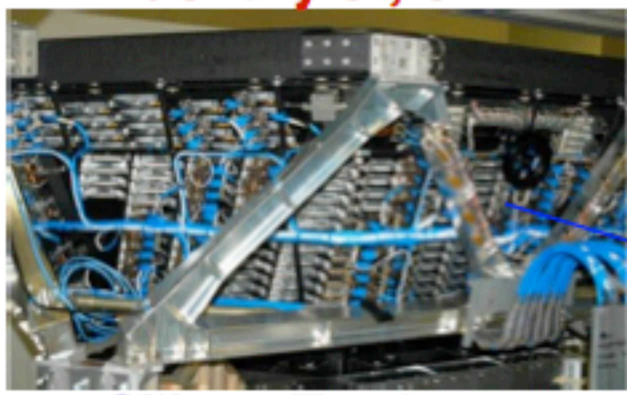
The calorimeter is used in the event selection : match between the track and the cluster (position, angle), cluster transverse size.

Reconstruction and selection are optimized using classification trees.

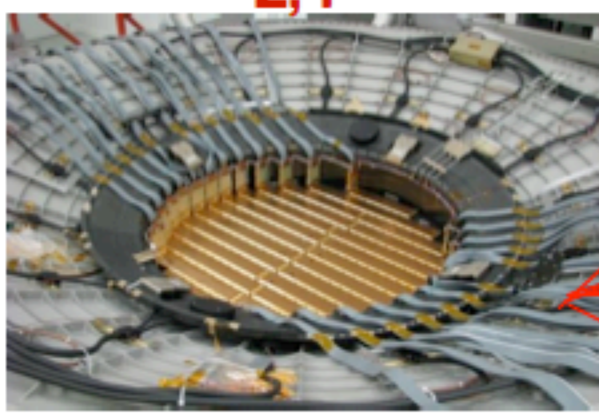
AMS: A TeV precision, multipurpose spectrometer

Particles and nuclei are defined by their charge (Z) and energy (E ~ P)

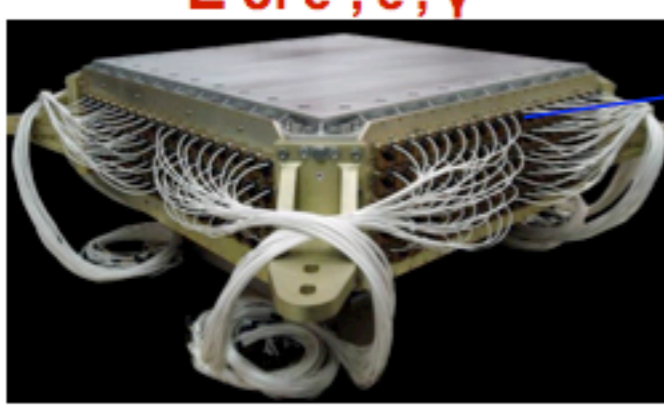
TRD
Identify e^+, e^-

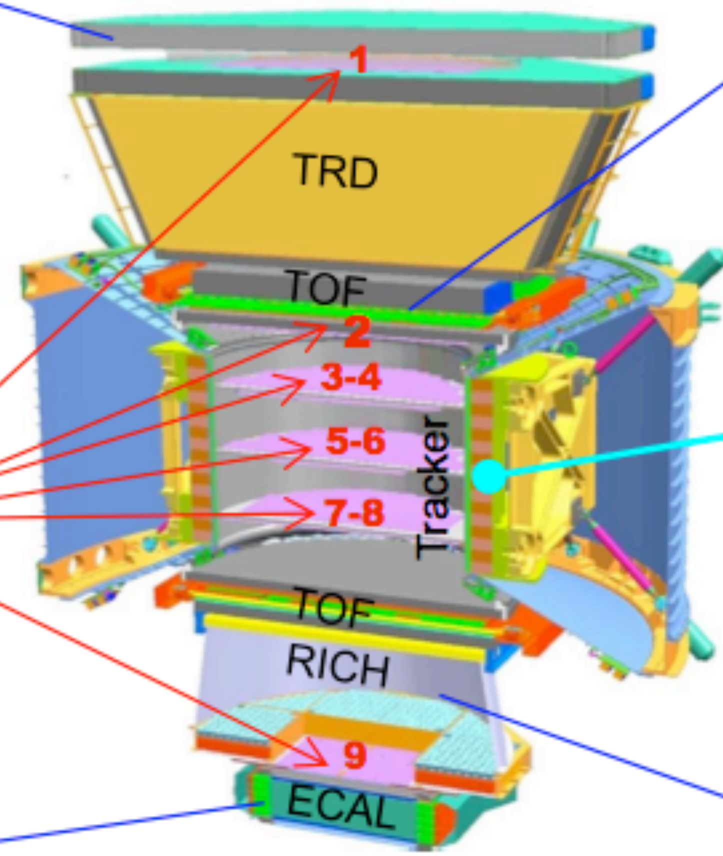


Silicon Tracker
Z, P



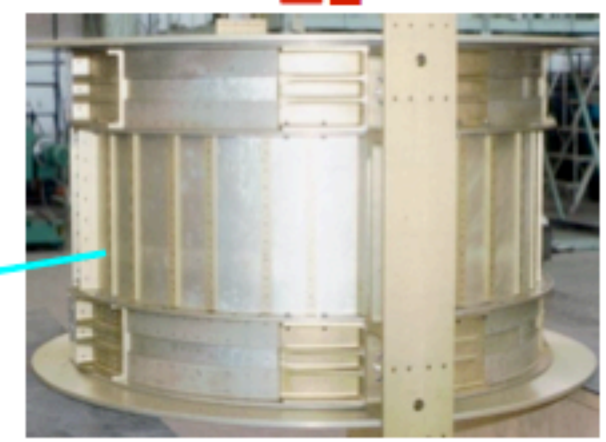
ECAL
E of e^+, e^-, γ



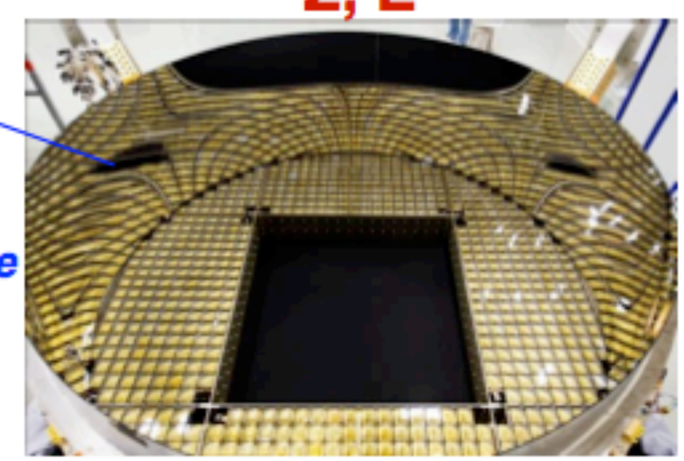


TOF
Z, E

Magnet
±Z



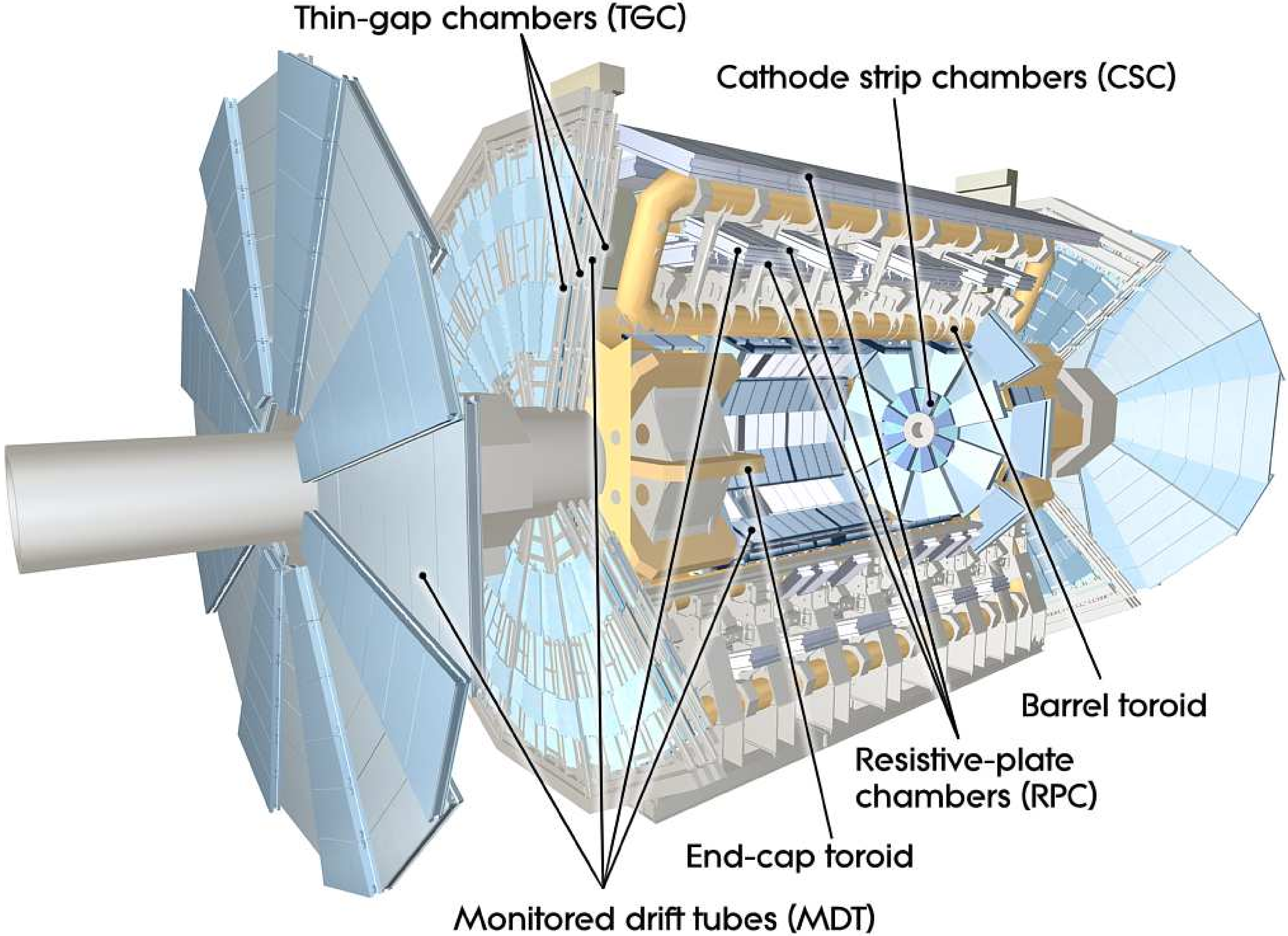
RICH
Z, E

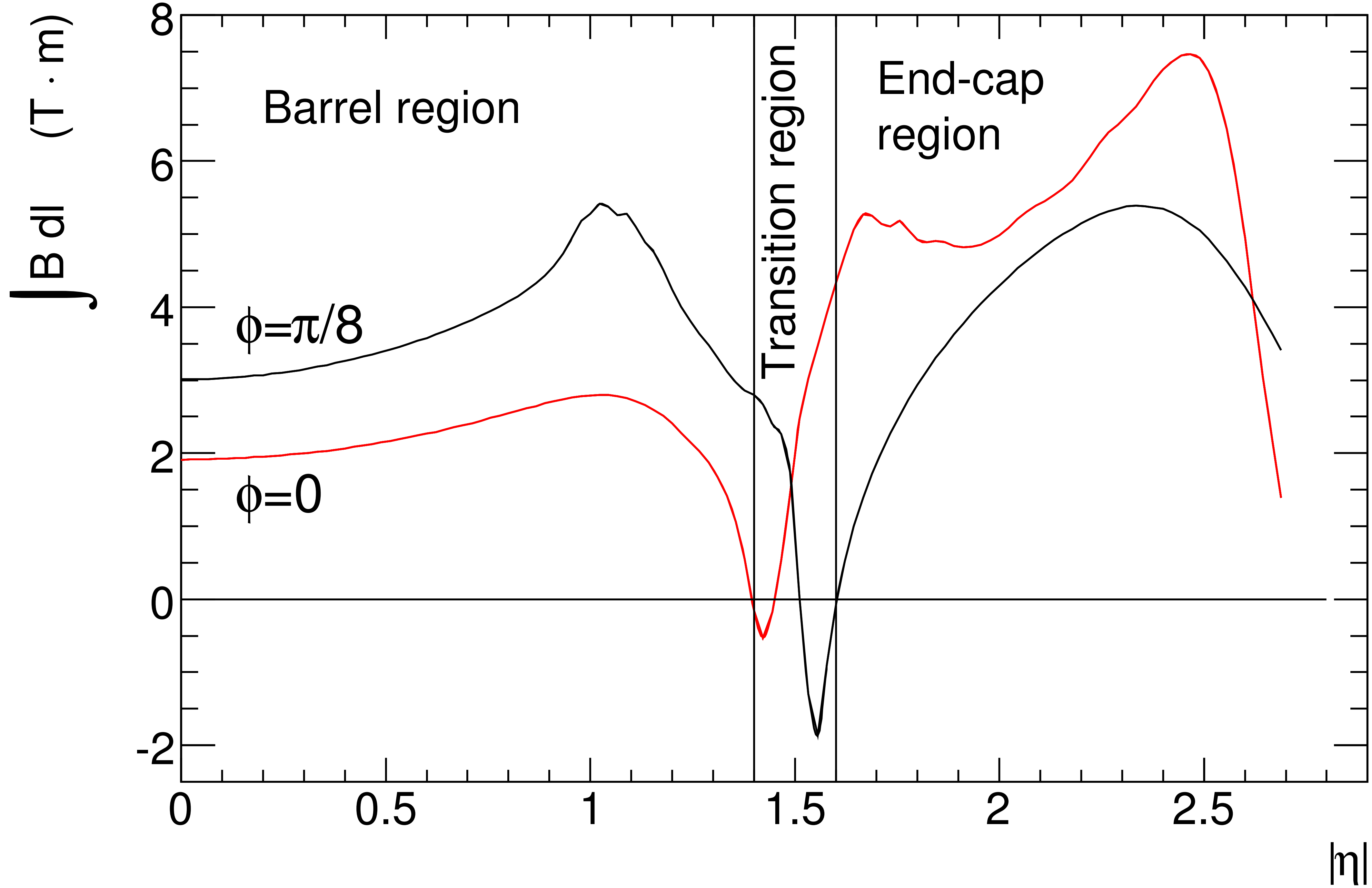


Z, P are measured independently by the Tracker, RICH, TOF and ECAL

Muon identification in hadron colliders

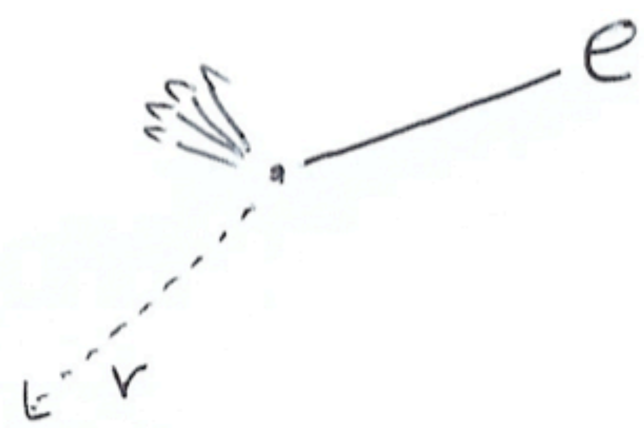
- Muons are usually clean signatures, less background than electrons
- Main sources of «muons»
  - punch through of hadronic showers
  - pi/k decays in the inner detector
  - Semileptonic B-hadron decays => «true» non-isolated muons
- Precise measurement of muons requires large magnetic detectors


Thin-gap chambers (TGC)
Cathode strip chambers (CSC)
Barrel toroid
Resistive-plate chambers (RPC)
End-cap toroid
Monitored drift tubes (MDT)


∫B dl (T · m)
Barrel region
Transition region
End-cap region
φ=π/8
φ=0
|η|
8
6
4
2
0
-2
0
0.5
1
1.5
2
2.5

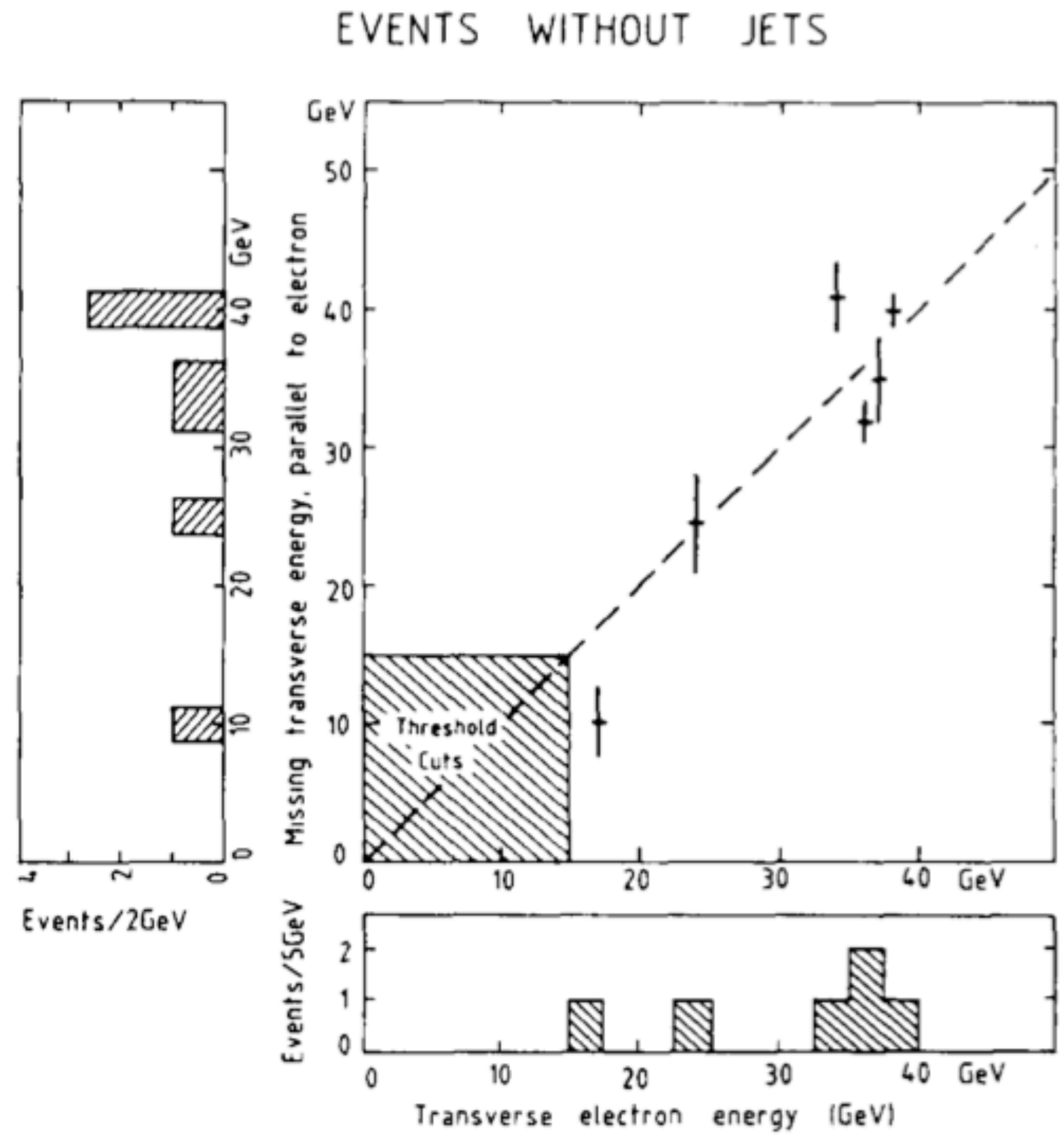
Neutrino «identification» in hadron colliders

- The probability of neutrino interaction in a collider experiment is ~null
- How to measure something that one does not detect ?

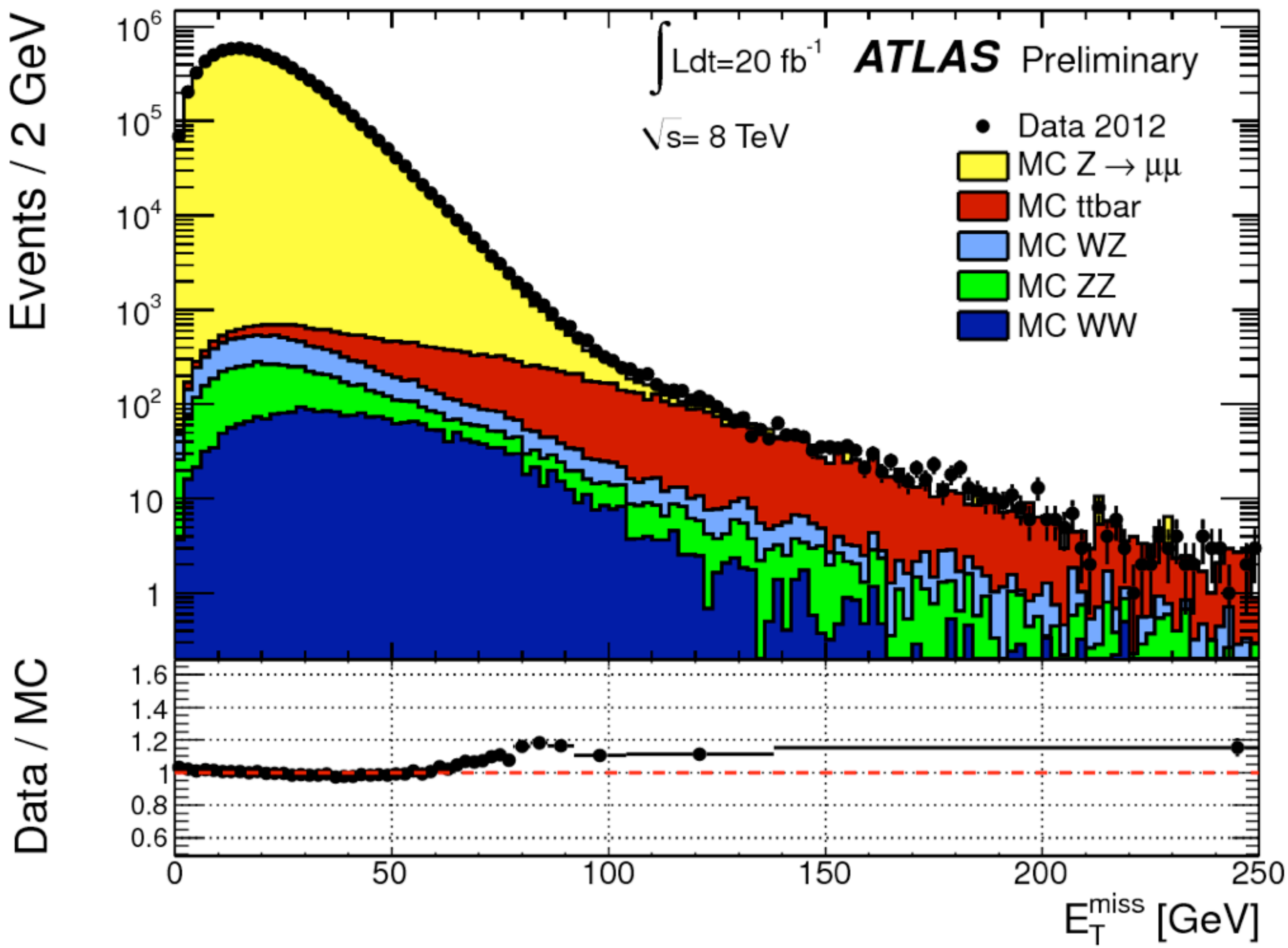


$$\vec{p}_T(\nu) = - \sum_i \vec{p}_T(\text{seen})$$

Missing transverse momentum for W boson discovery (1983)



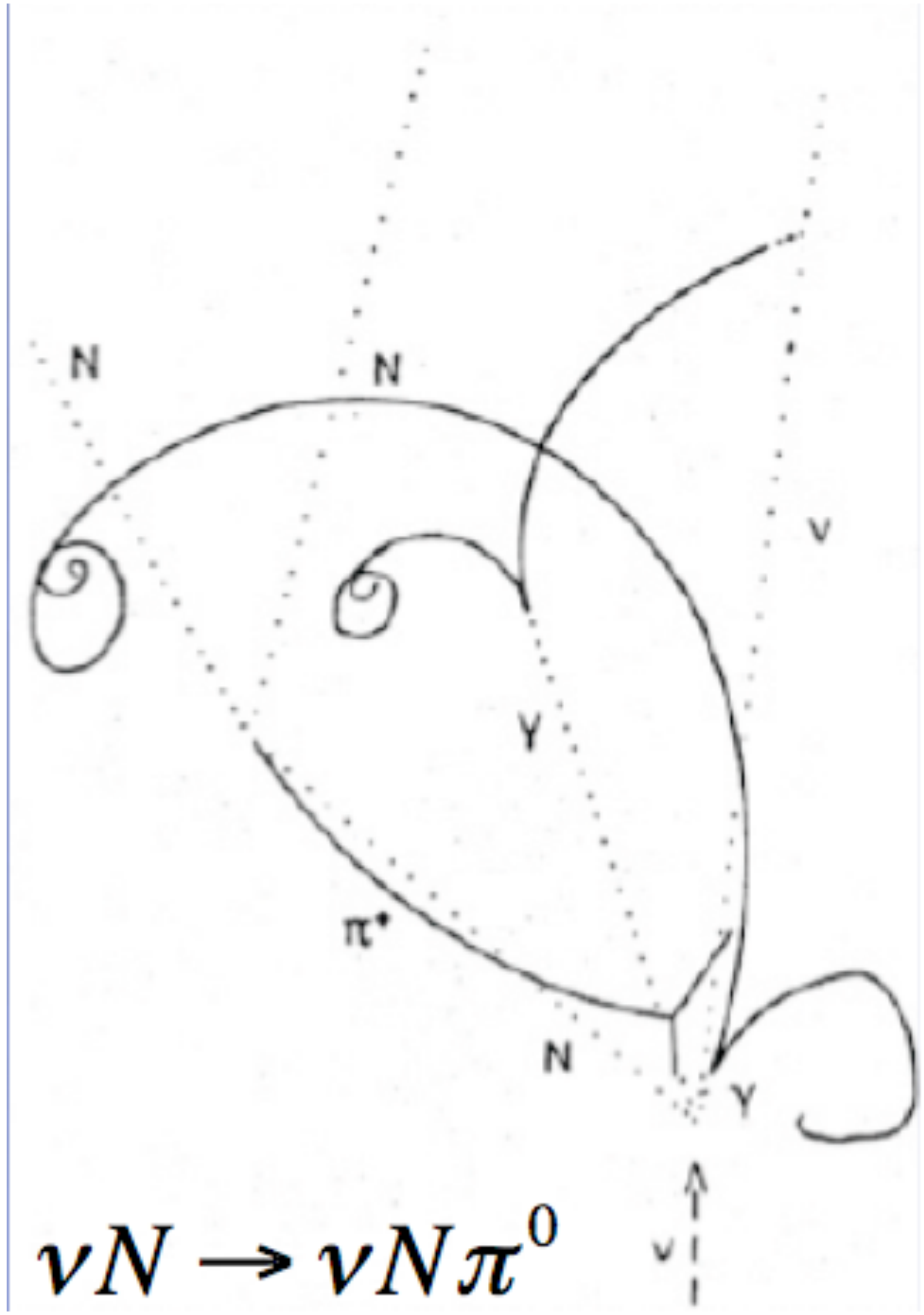
Missing transverse momentum in LHC under high pileup conditions



Direct detection of neutrinos

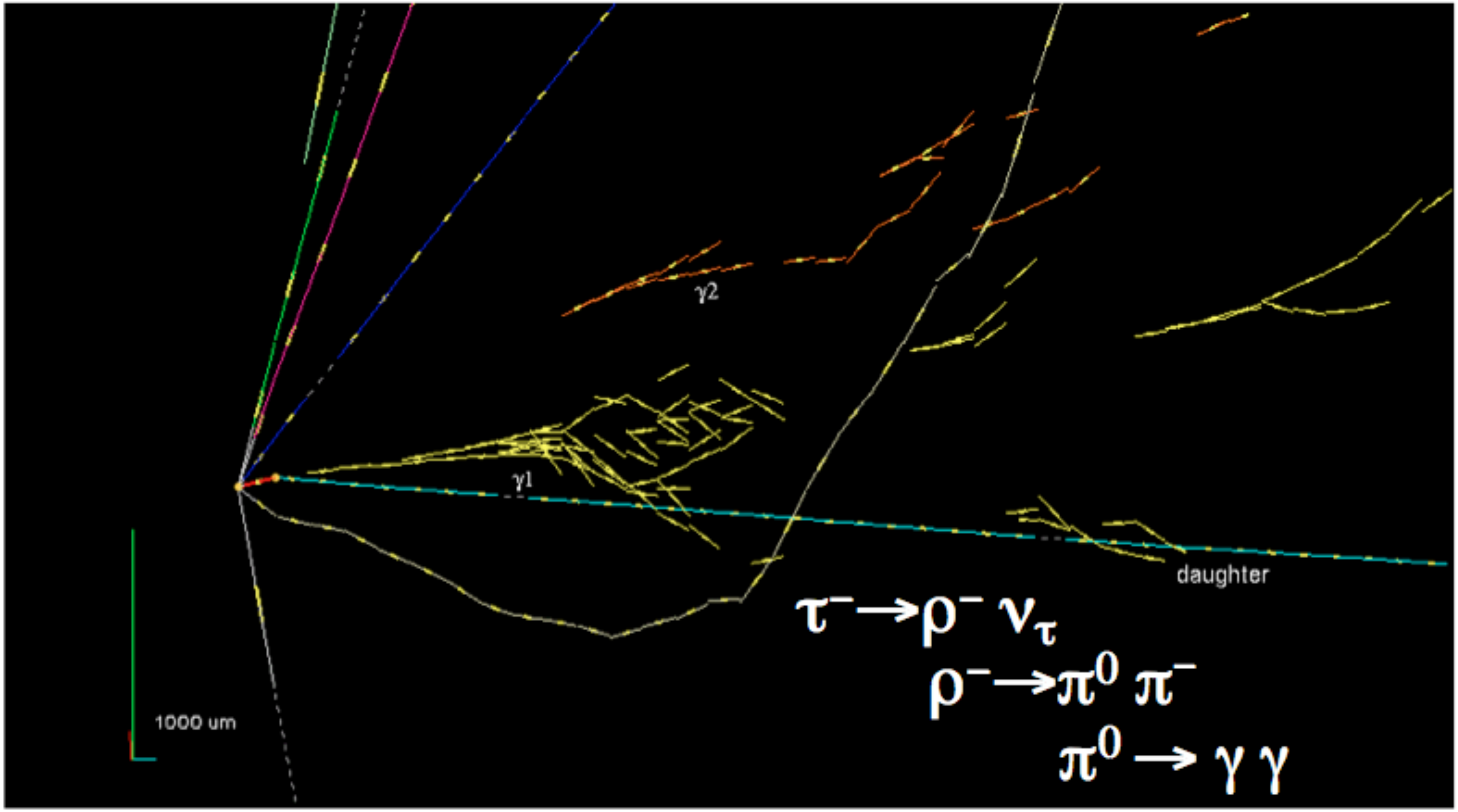
- High flux of incoming neutrinos (for instance neutrino beams)
- High mass detector
- => can observed neutrino interactions
  - Charged currents: produce e,mu or tau depending on neutrino flavor at the interaction
  - Neutral currents: ~universal for all (non-sterile) neutrinos
- Neutrino cross-section increases with energy
  - at O(> PeV) energy, earth becomes opaque to neutrinos

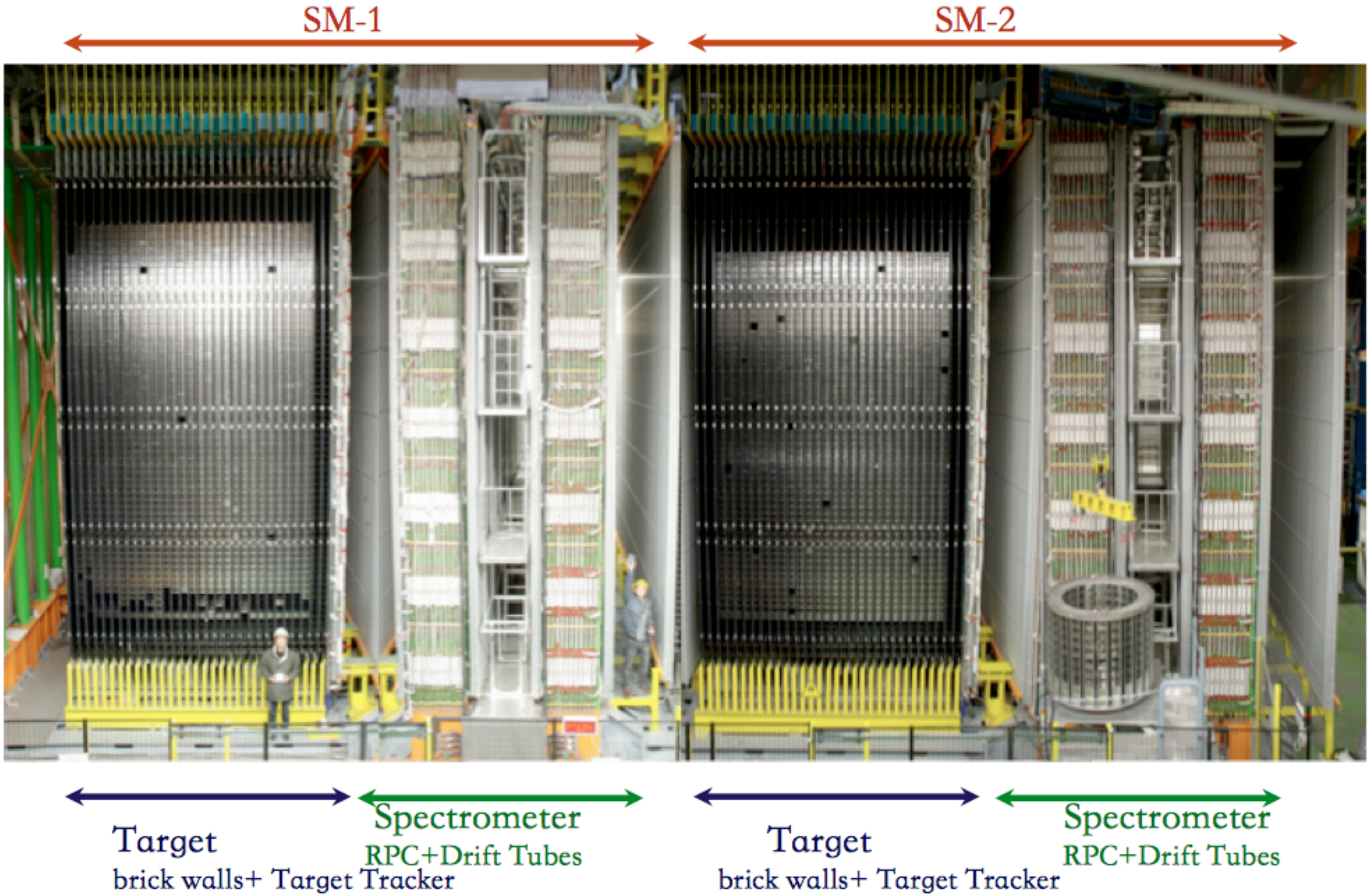
What is this event ?



$\nu N \rightarrow \nu N \pi^0$

and this one ?


γ2
γ1
daughter
1000 um
$\tau^- \rightarrow \rho^- \nu_\tau$
$\rho^- \rightarrow \pi^0 \pi^-$
$\pi^0 \rightarrow \gamma\gamma$


SM-1
SM-2
Target
brick walls+ Target Tracker
Spectrometer
RPC+Drift Tubes
Target
brick walls+ Target Tracker
Spectrometer
RPC+Drift Tubes

Measure beta or gamma of particle

- Direct measurement of velocity («time of flight»
    - v = d/t
- Measurement of beta.gamma through ionization energy loss
- Measurement of beta through Cherenkov radiation
- Measurement of gamma through Transition radiation

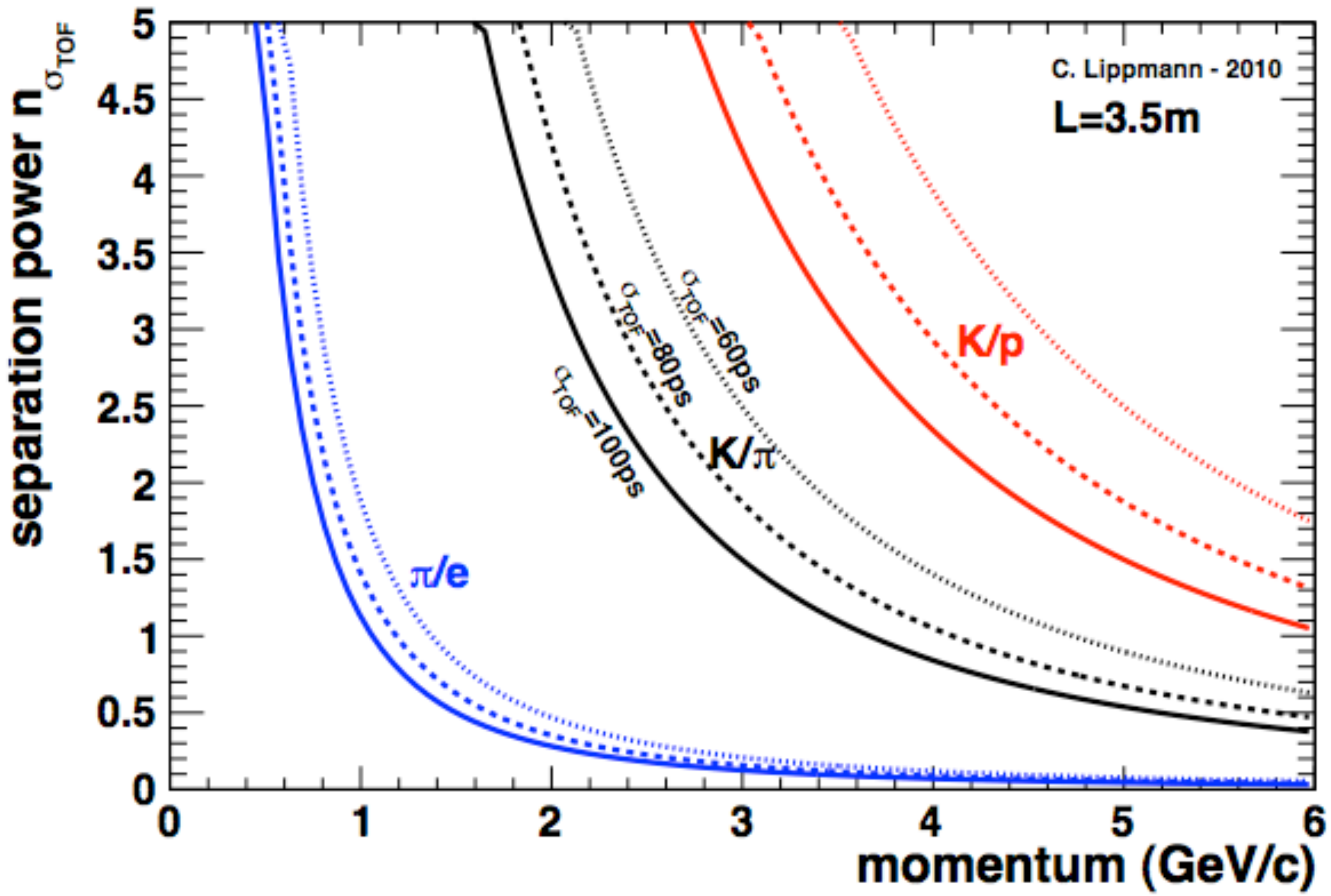
time of flight

$$\beta = \frac{v}{c} = \frac{L}{t \cdot c}$$

$$m = \frac{p}{c}\sqrt{\frac{c^2 t^2}{L^2} - 1} \qquad \frac{dm}{m} = \frac{dp}{p} + \gamma^2\left(\frac{dt}{t} + \frac{dL}{L}\right)$$

For 2 hypothesis m_A, m_B:

$$|t_A - t_B| = \frac{L\,c}{2p^2}\,|m_A^2 - m_B^2|$$



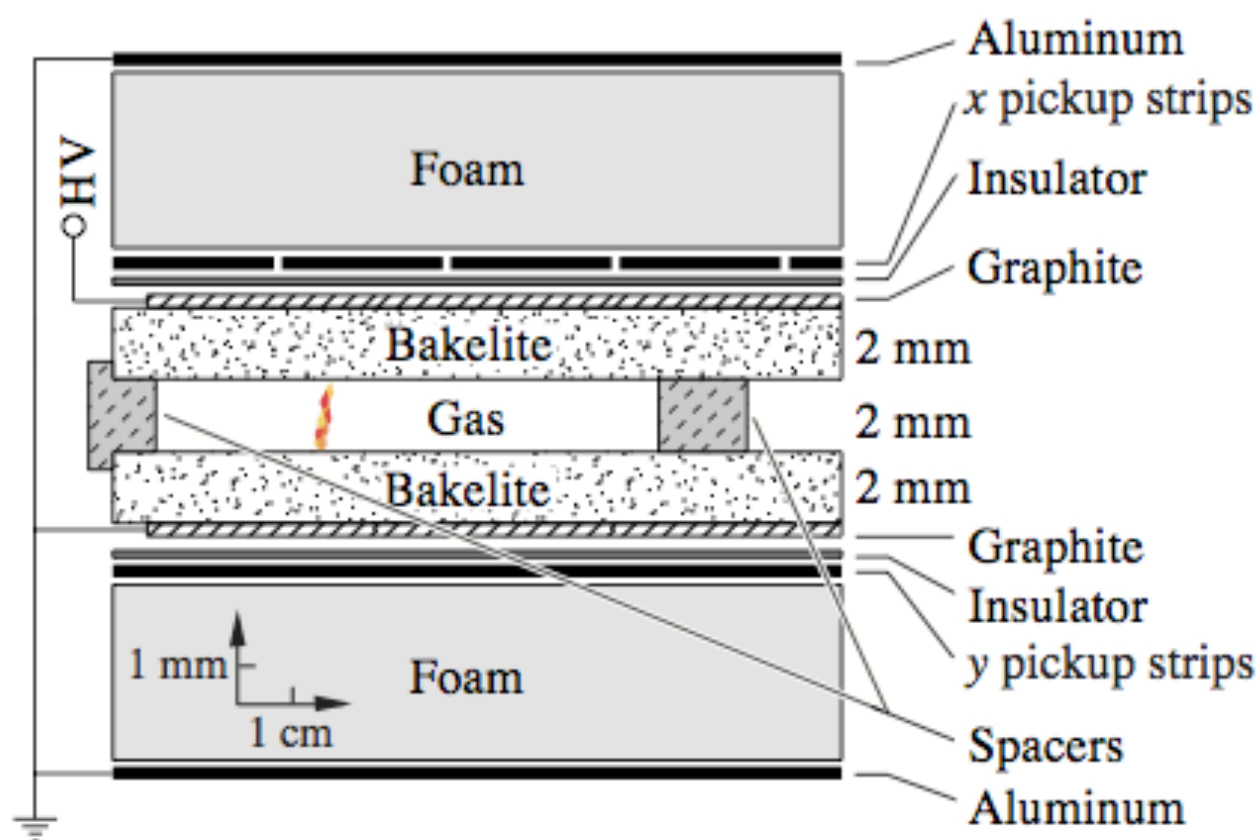
Dedicated detectors for time measurement can reach < 100 ps accuracy even on large system

At LHC, the collision time has an intrinsic jitter of ~140 ps (bunch length)

Need dedicated measurement to remove this contribution from time resolution

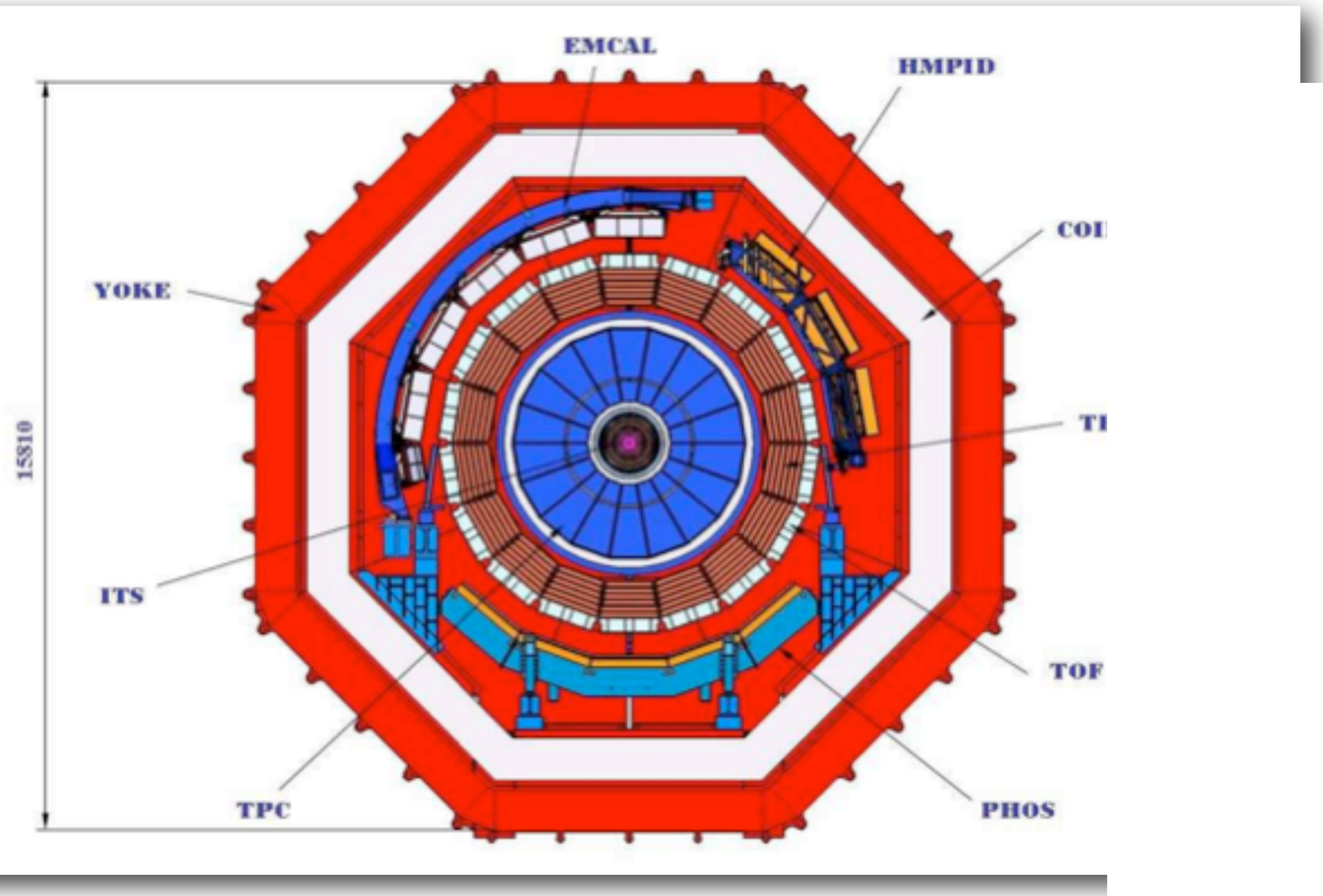
Most commonly used detectors for timing were based on scintillation
(can also use other techniques like calorimetry, etc..)

Gaseous ionization detectors like RPC developed to cover large area in a cost-effective way

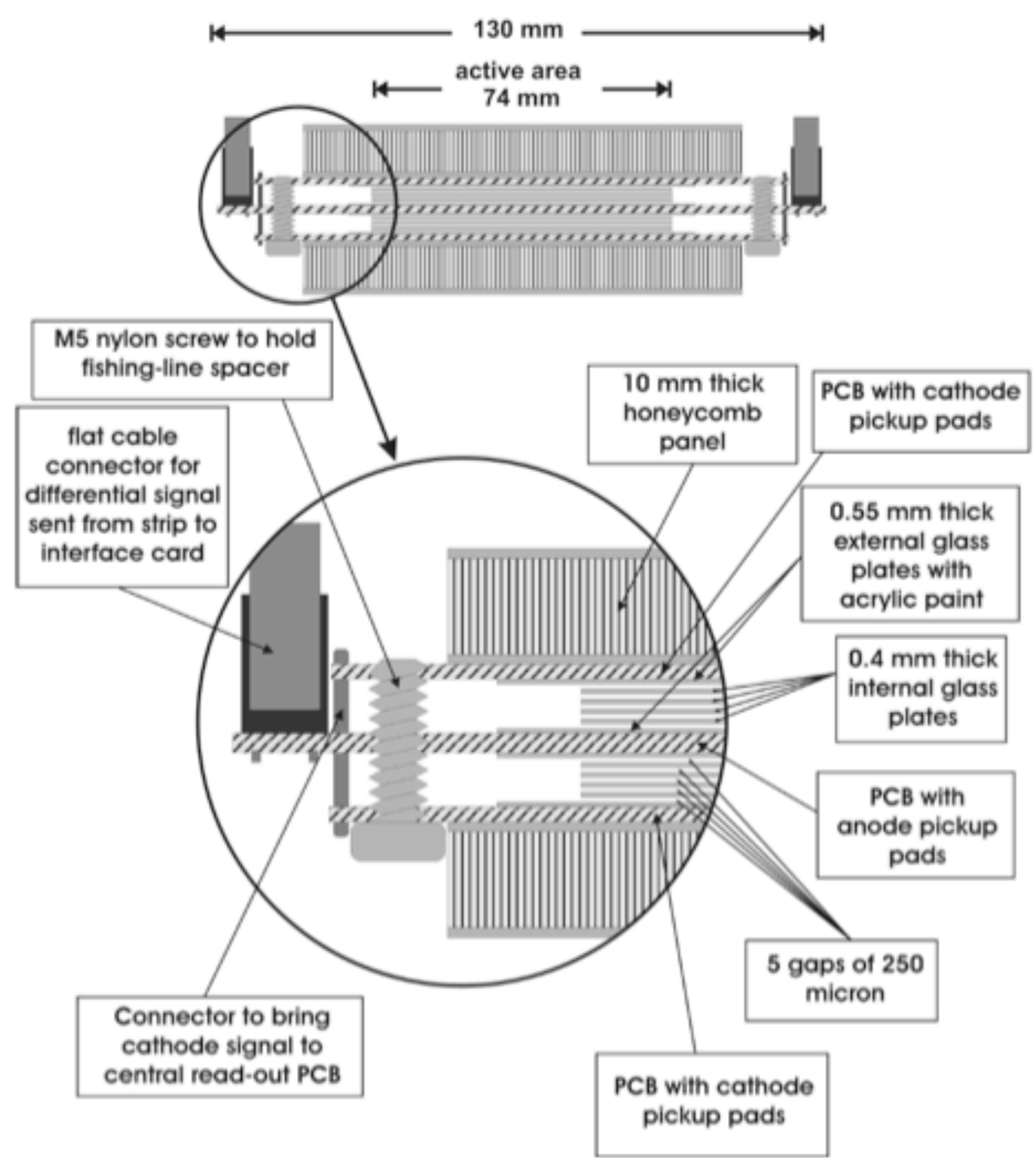


Strong uniform electric field => avalanche starts immediately after primary ionization
Can reach intrinsic time resolution of ~ 50 ps for multigap RPC
Rate limitation O(kHz/cm2)

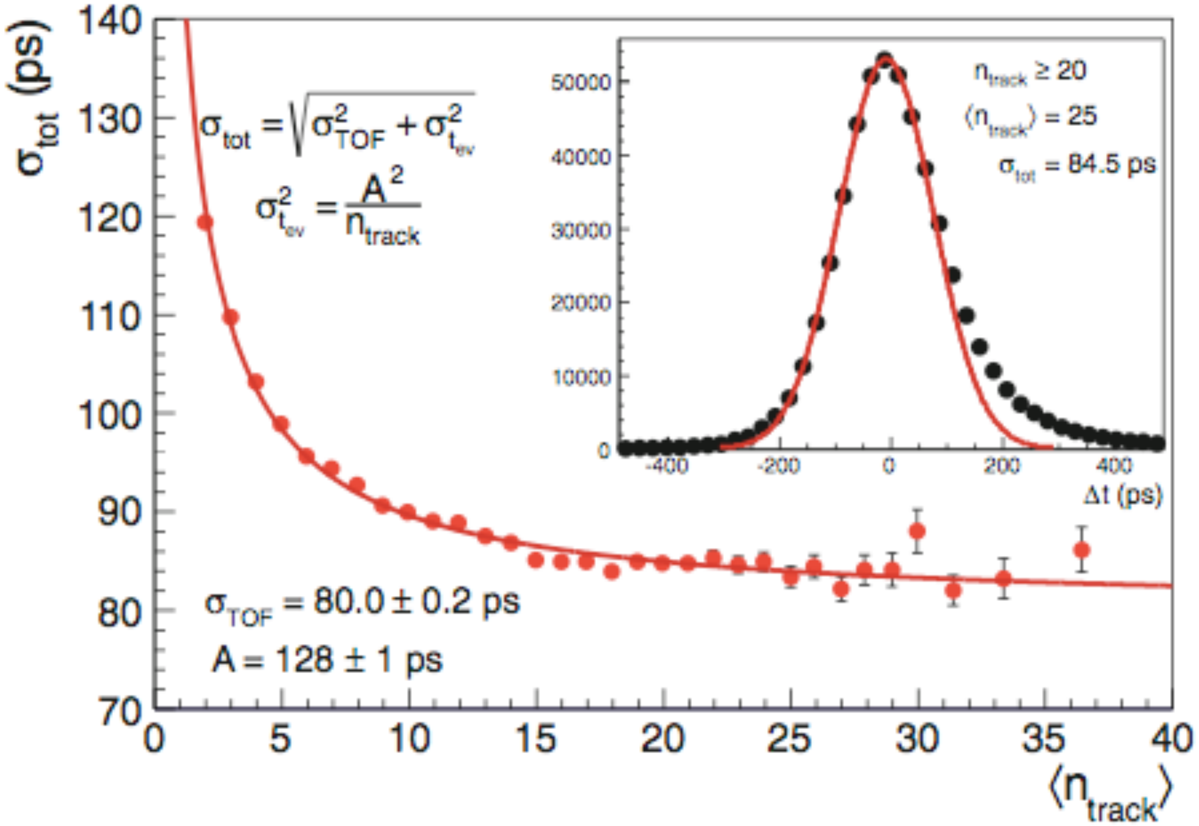
ALICE time of flight based on MRPC $\sim 10^5$ channels



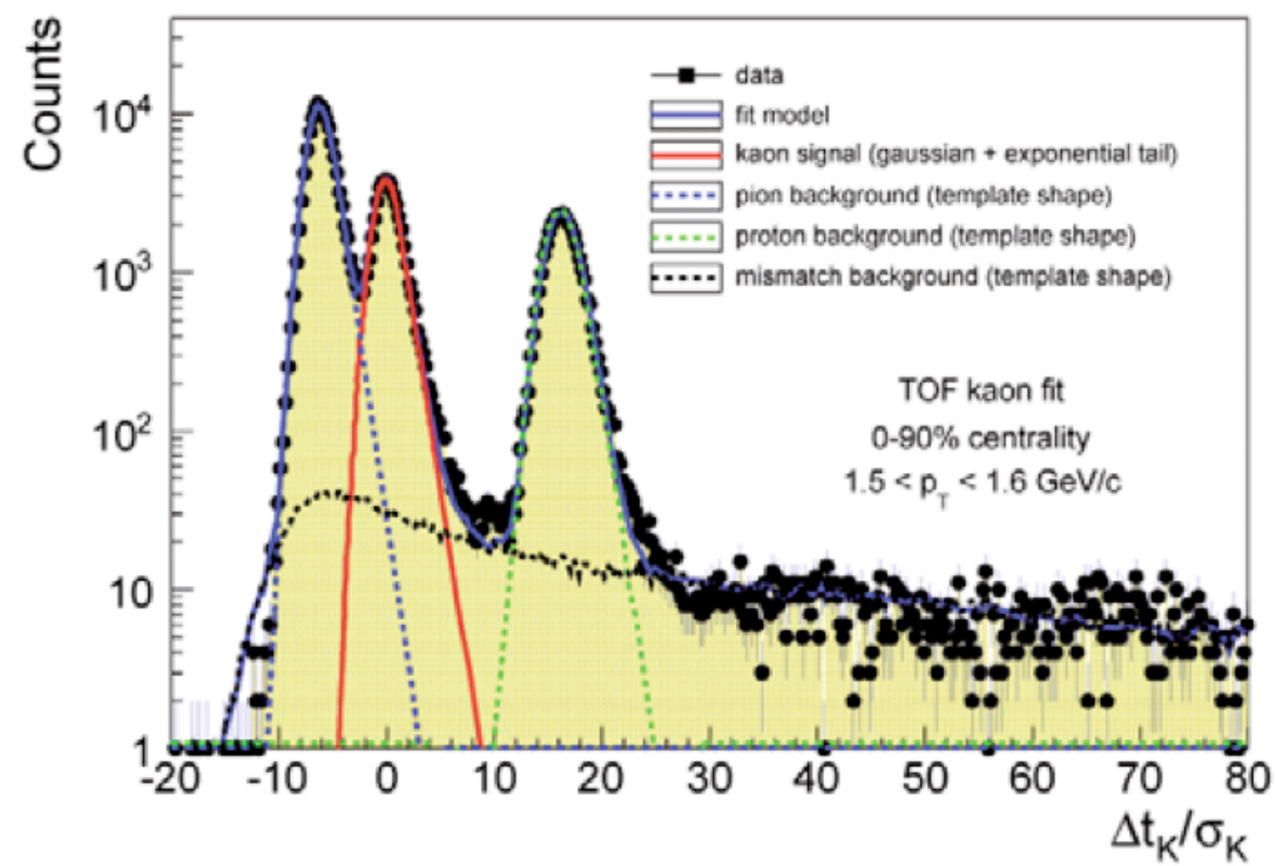
TOF @R=3.7m from interaction point

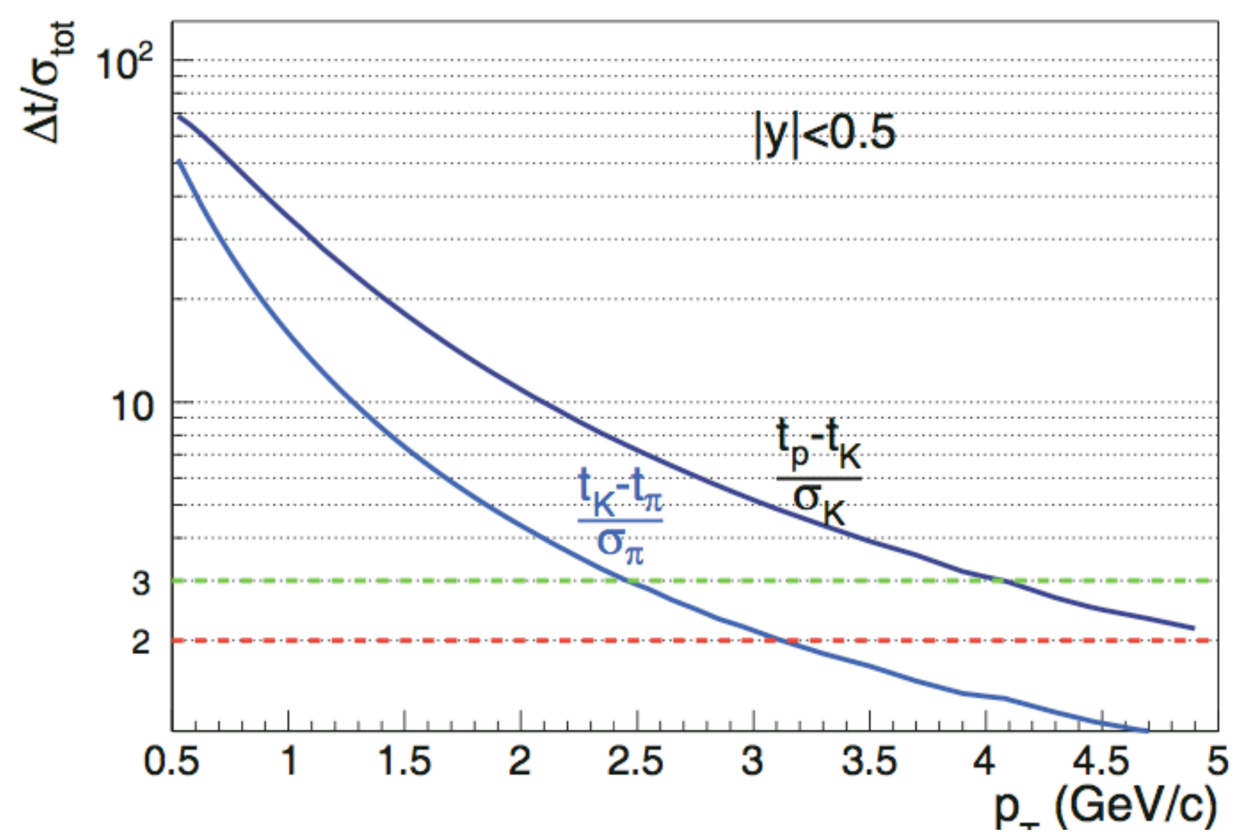


Measured time resolution
in ALICE

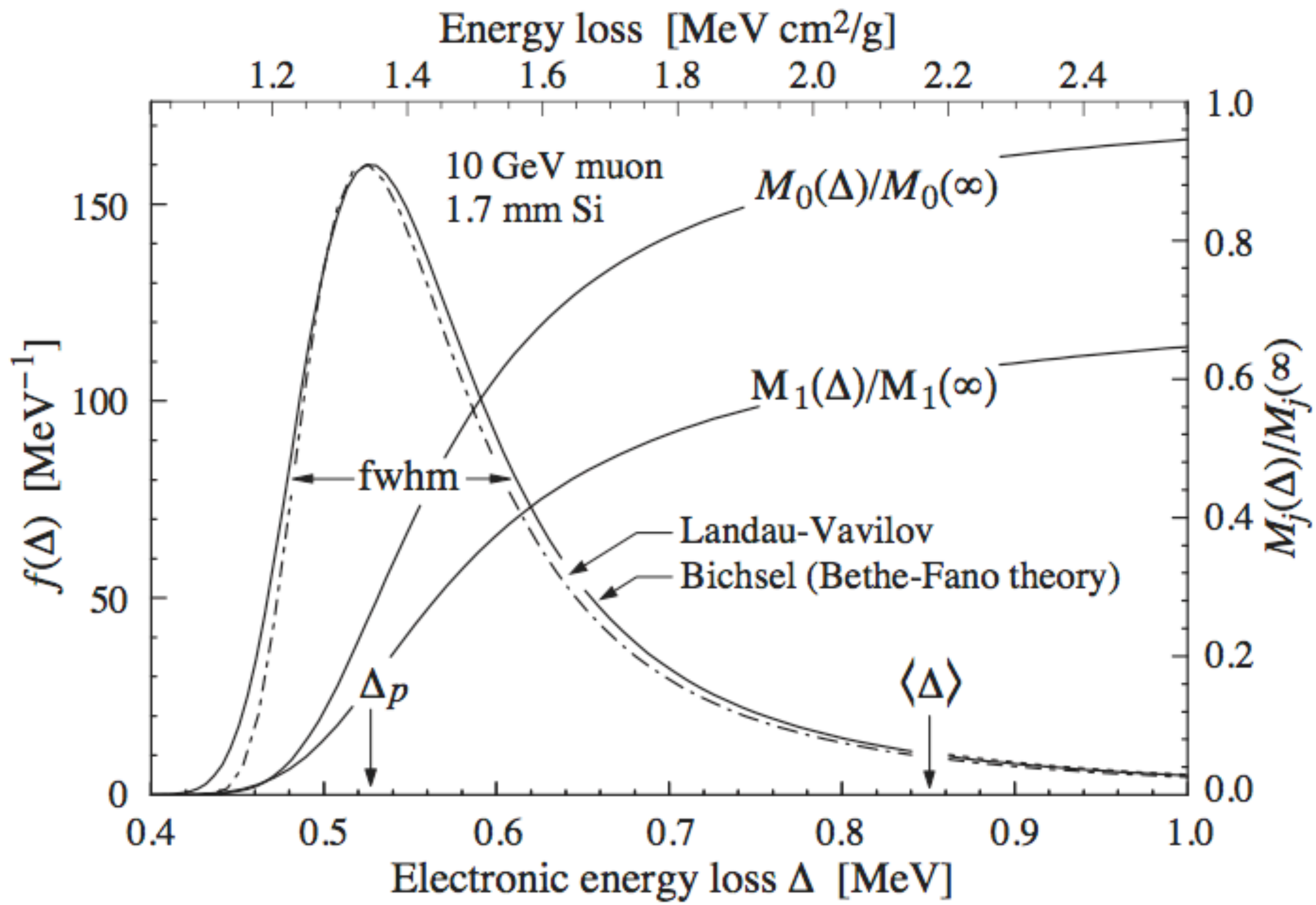


Intrinsic resolution +
time jitters (electronics,
clock)+ channel to
channel variation +
residual time slewing
effects





Ionization measurement



Formula for restricted energy loss

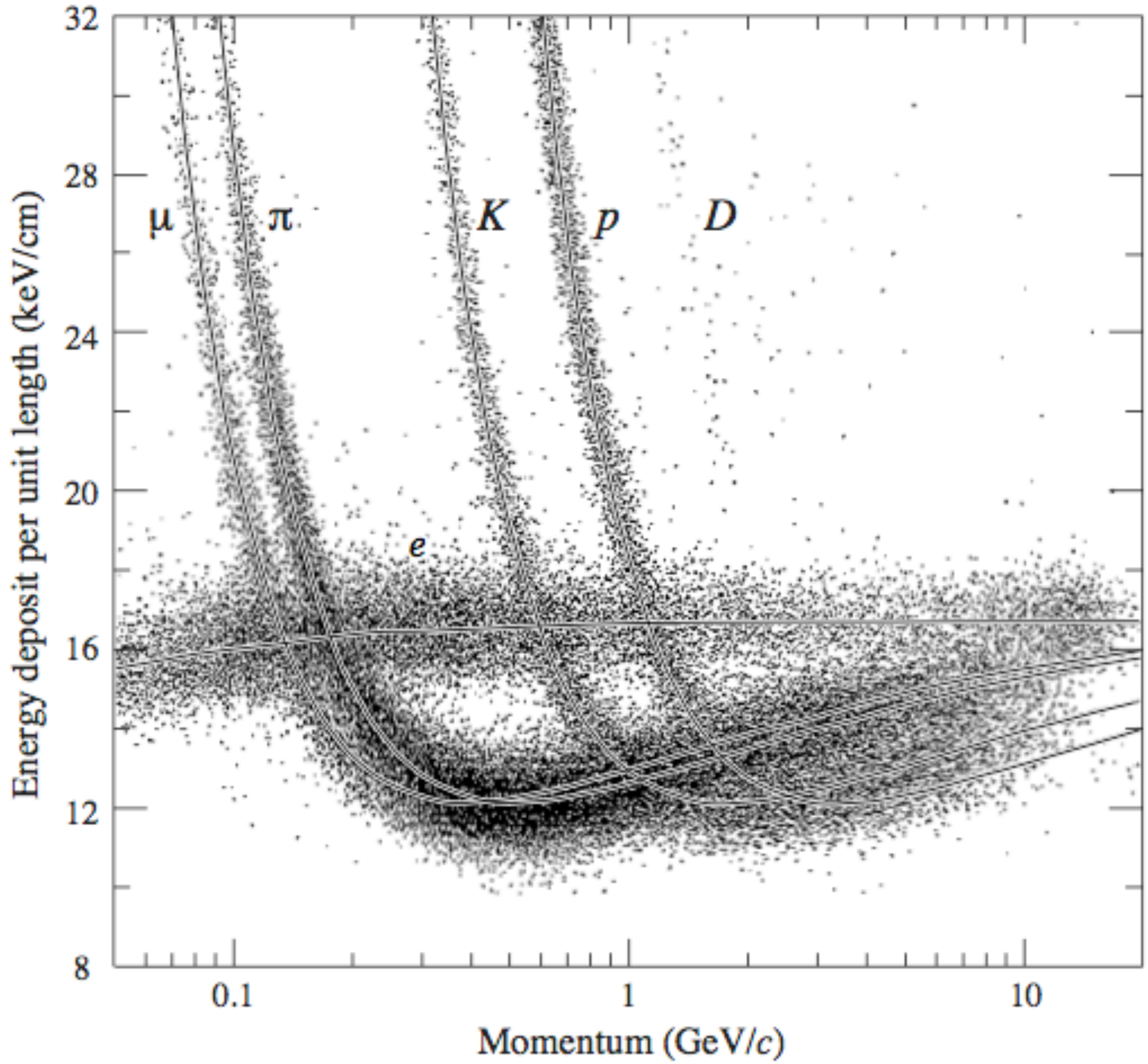
$$\left\langle \frac{dE}{dx} \right\rangle \propto \frac{z^2}{\beta^2} \left(\log \frac{\sqrt{2 m_e c^2 E_{cut}}\, \beta\gamma}{I} - \frac{\beta^2}{2} - \frac{\delta}{2} \right)$$

I = effective excitation energy

δ = density correction effect

E_{cut} = upper limit for energy transfer in single collision

Ionisation measurement in a TPC



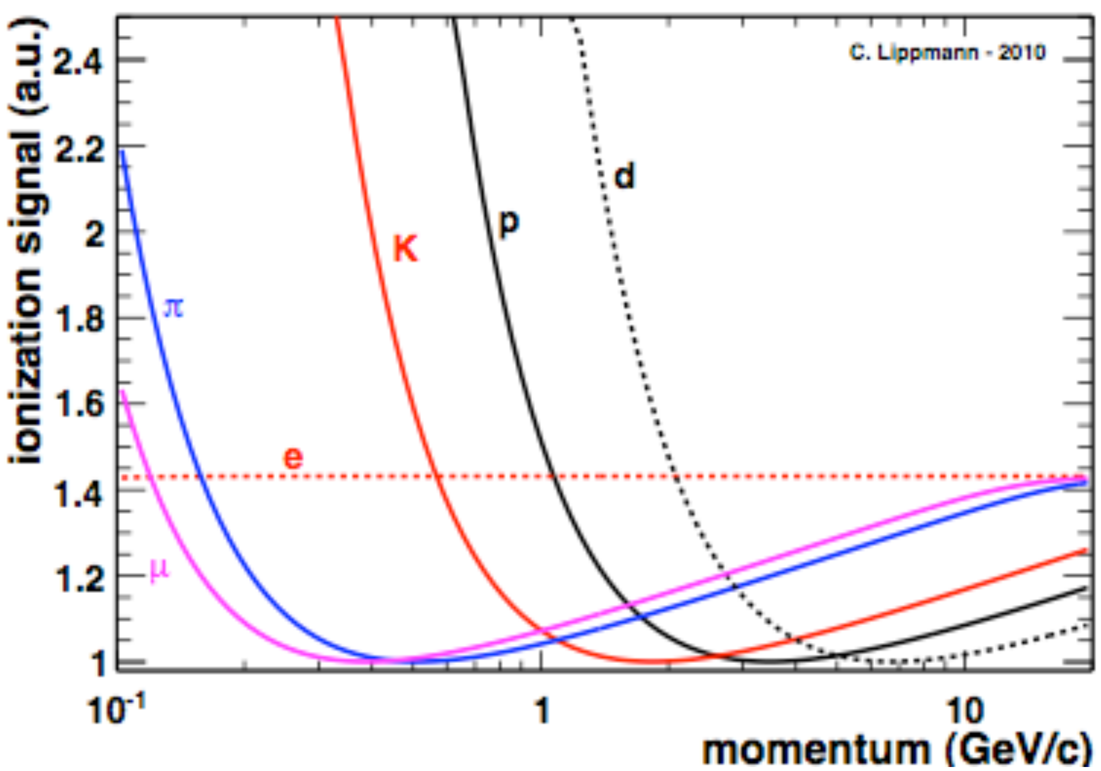
Can use gaseous or solid state counter to measure ionisation

Provide signal pulse height ~ N(electrons liberated in ionization) and measurement of track length
=> allows one to compute dE/dx

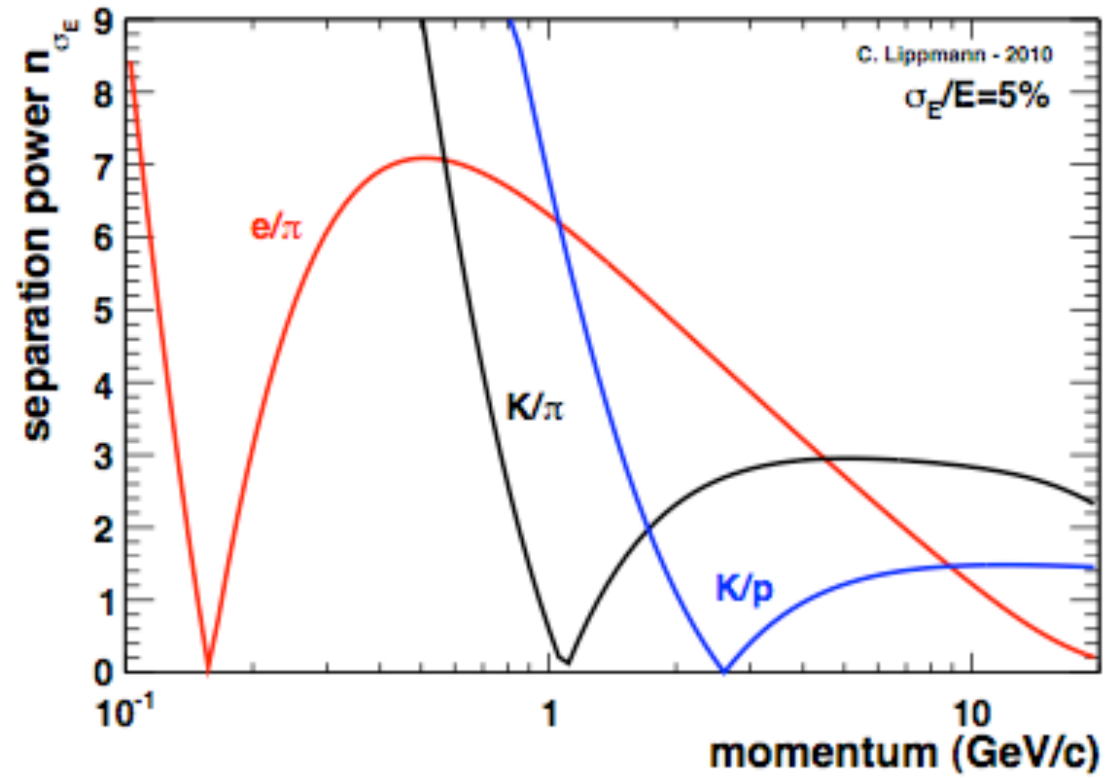
Average several measurements with a truncated mean to reduce tail impact

Typical other errors affecting measurement:
- energy calibration of the detector
- detector conditions (for instance gas pressure)
- detector geometry and track orientation (affects track length)
- overlapping tracks in dense environment
- etc..

Typical ionization signals vs p (gaseous detector)
(for Si detector, plateau only slightly above minimum => less separation at high energy)



Separation assuming 5% resolution

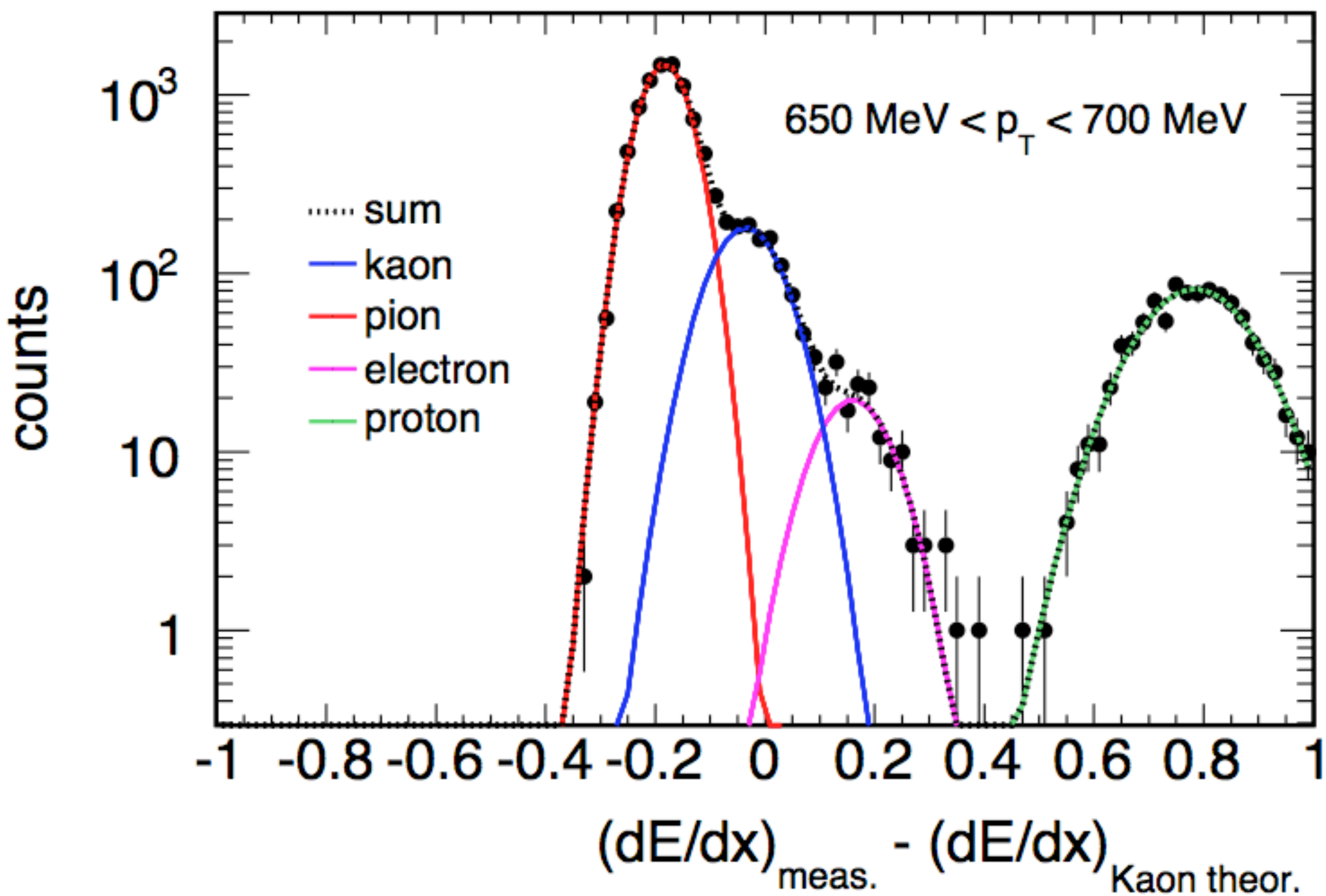


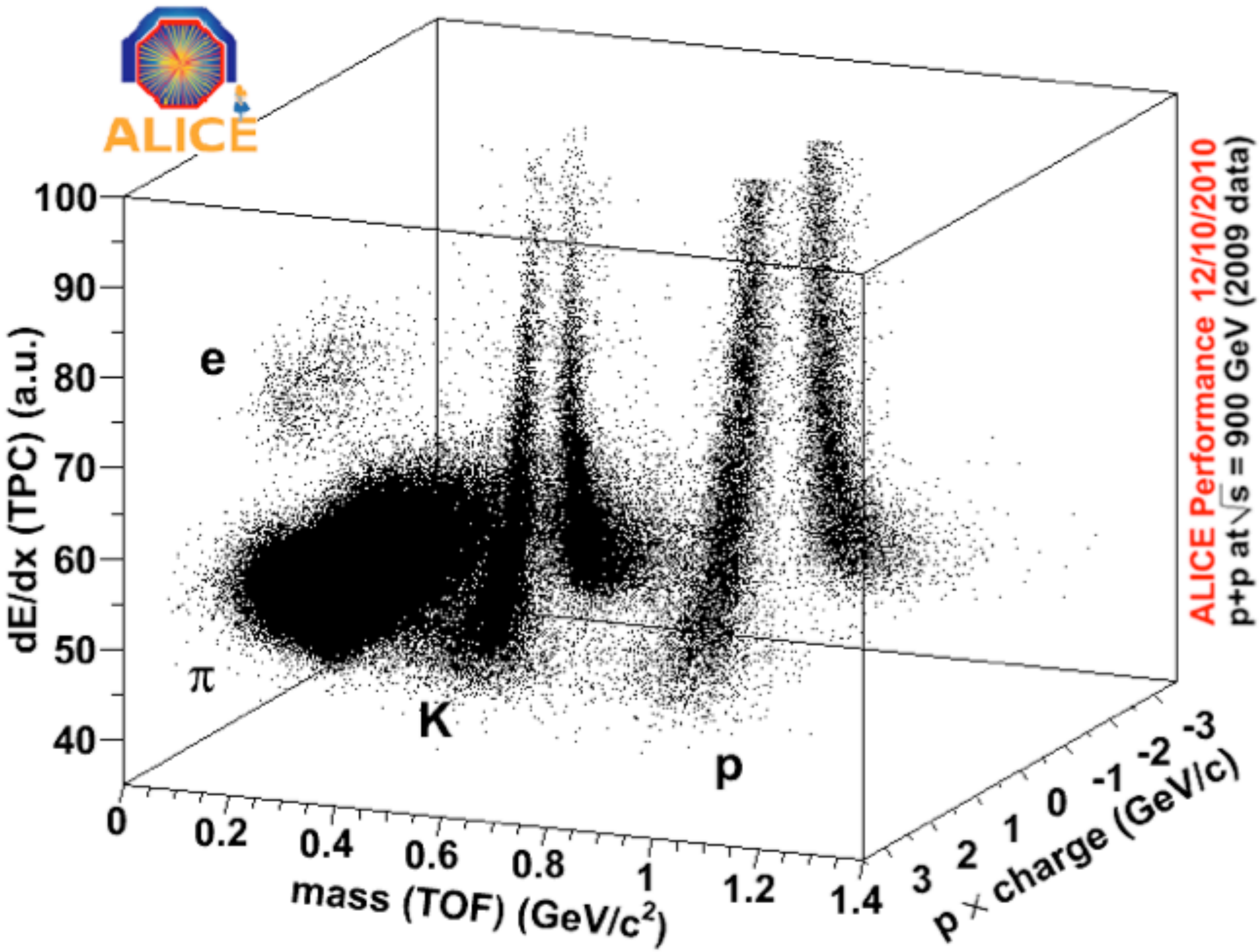
Empirical scaling formula for resolution in gaseous detector:

$$\sigma_E = 0.41\, N_R^{-0.43} (xP)^{-0.32} \,.$$

Nr = number of measurements
x = thickness of sampling layers
(x.Nr = total detector thickness)
P = pressure

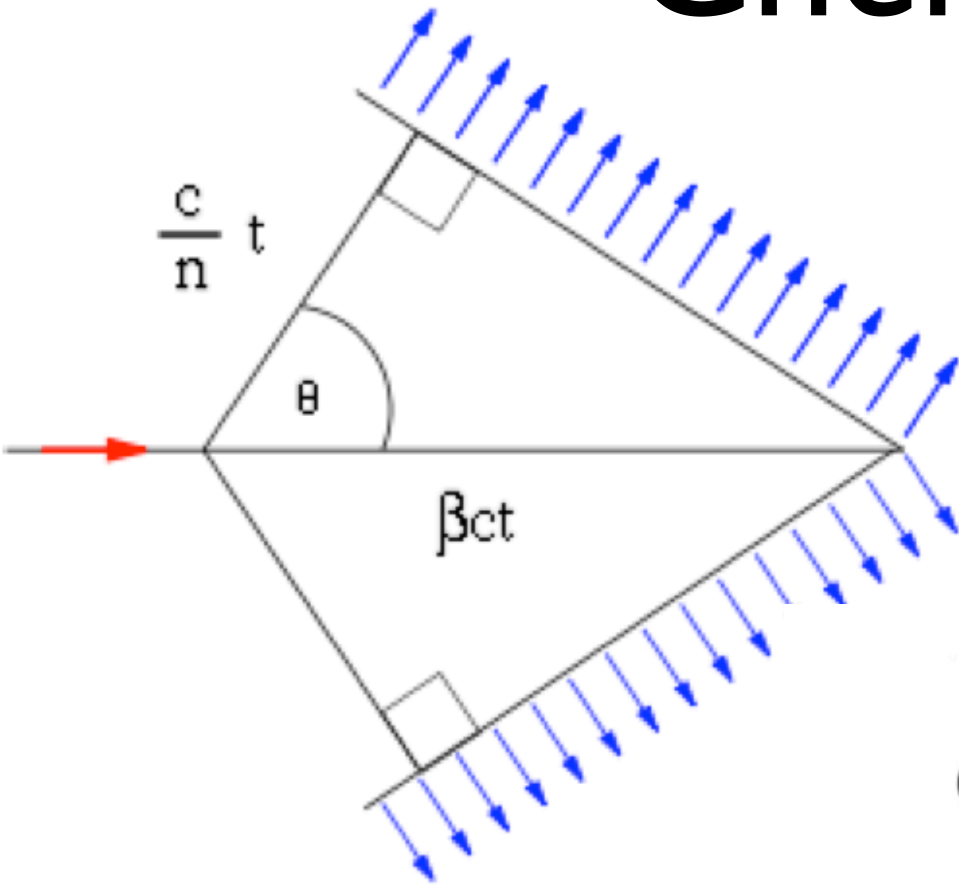
ALICE TPC detector
reaches ~5% dEdx resolution





ALICE
ALICE Performance 12/10/2010
p+p at √s = 900 GeV (2009 data)
dE/dx (TPC) (a.u.)
100
90
80
70
60
50
40
e
π
K
p
0
0.2
0.4
0.6
0.8
1
1.2
1.4
mass (TOF) (GeV/c²)
3 2 1 0 -1 -2 -3
p × charge (GeV/c)

Cherenkov radiation



$$\cos\theta_c = \frac{1}{\beta n} \qquad (n = n(E_{photon}))$$

Threshold velocity $\beta_t = \frac{1}{n}$

$$\frac{d^2 N_\gamma}{dE\,dx} = \frac{\alpha z^2}{\hbar c}\sin^2\theta_c \qquad \text{or} \qquad \frac{d^2N}{d\lambda\,dx} \propto \frac{1}{\lambda^2}$$

$$\approx 370 \sin^2\theta_c(E)\ \text{eV}^{-1}\text{cm}^{-1} \text{ for } z=1$$

different type of Cherenkov detectors

- threshold Cherenkov detectors: yes/no decision depending if particle is above/below threshold beta=1/n
  - main issue is optimising photon detection and minimising noise
- Imaging Cherenkov detectors

$$\frac{\sigma_\beta}{\beta} = \tan\theta_C \; \sigma(\theta_C)$$

with $$\sigma(\theta_C) = \frac{\langle \sigma(\theta_i) \rangle}{\sqrt{N_{p.e}}} \oplus C$$

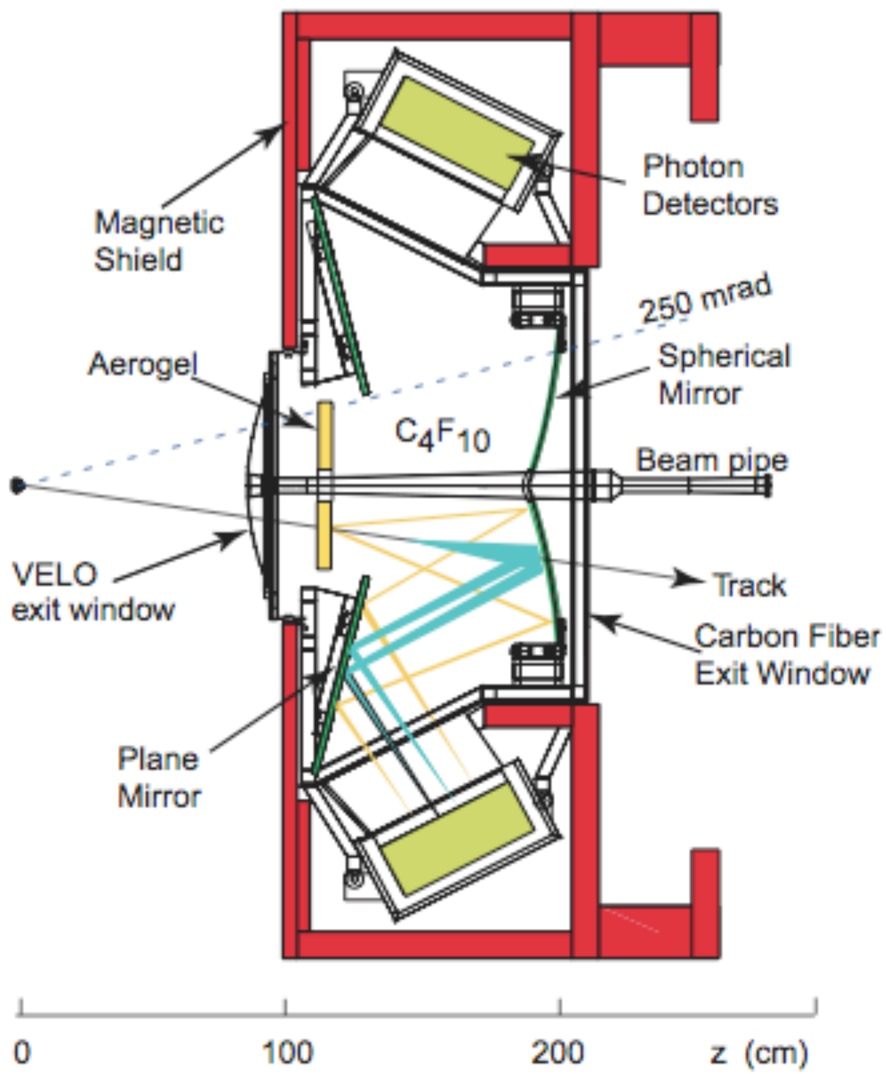
$\langle \sigma(\theta_i) \rangle$ = average single photo electron resolution (optics, detector geometry, ...)

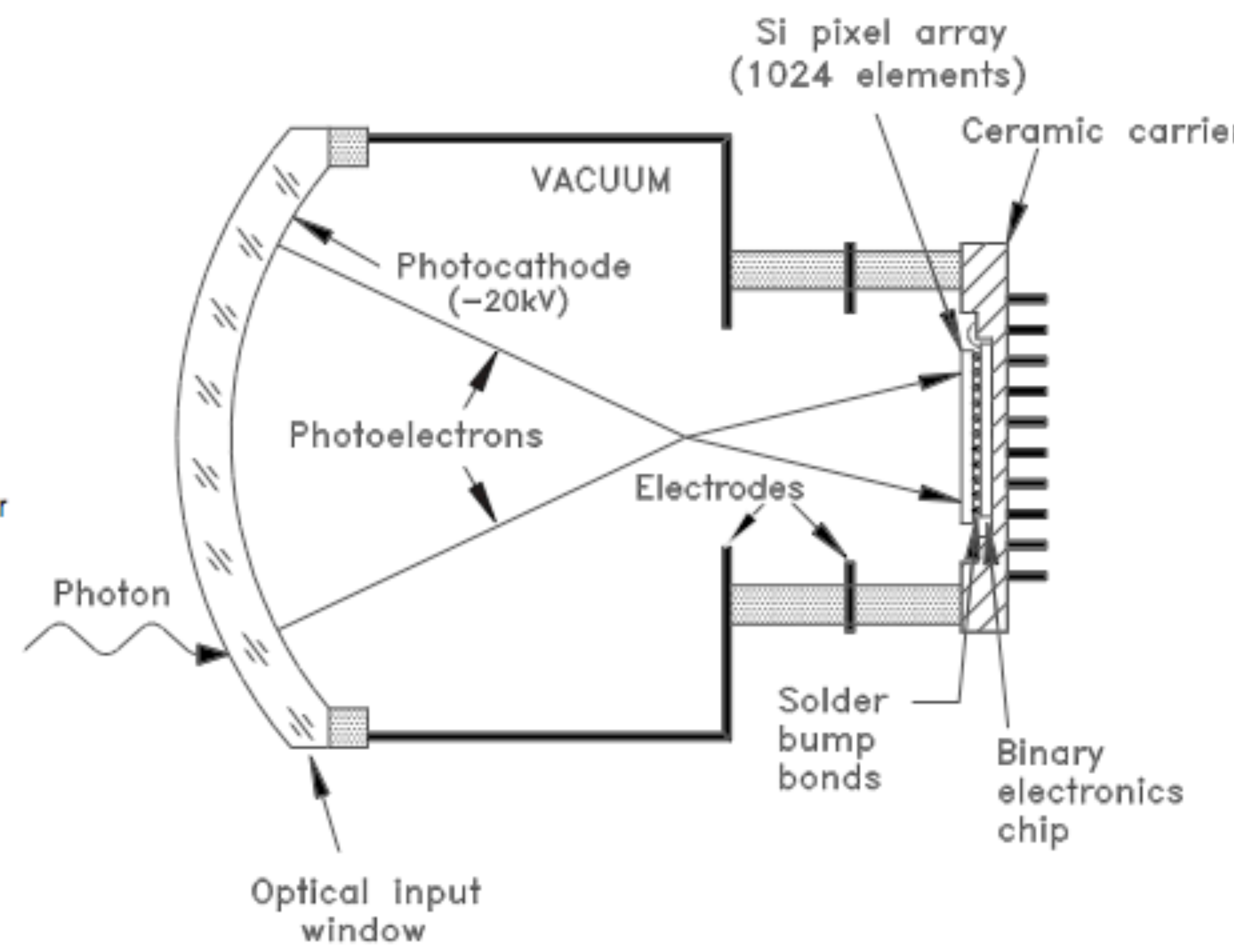
$N_{p.e}$ = number of photo electron detected

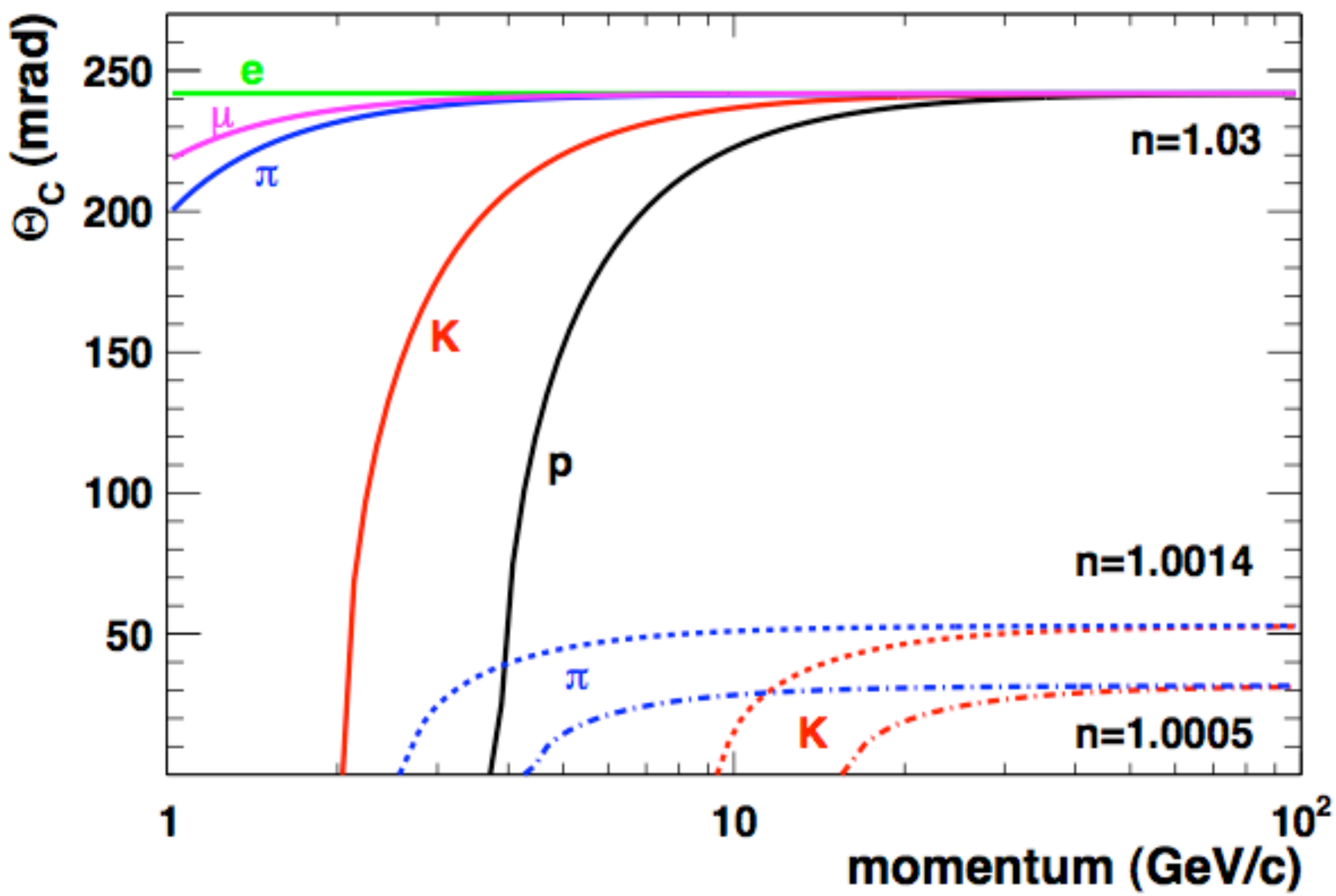
C : alignement, multiple scattering, ambiguities background, etc..

$$N_{sigma} \simeq \frac{|m_1^2 - m_2^2|}{2p^2 \sigma(\theta_c) \sqrt{n^2-1}}$$

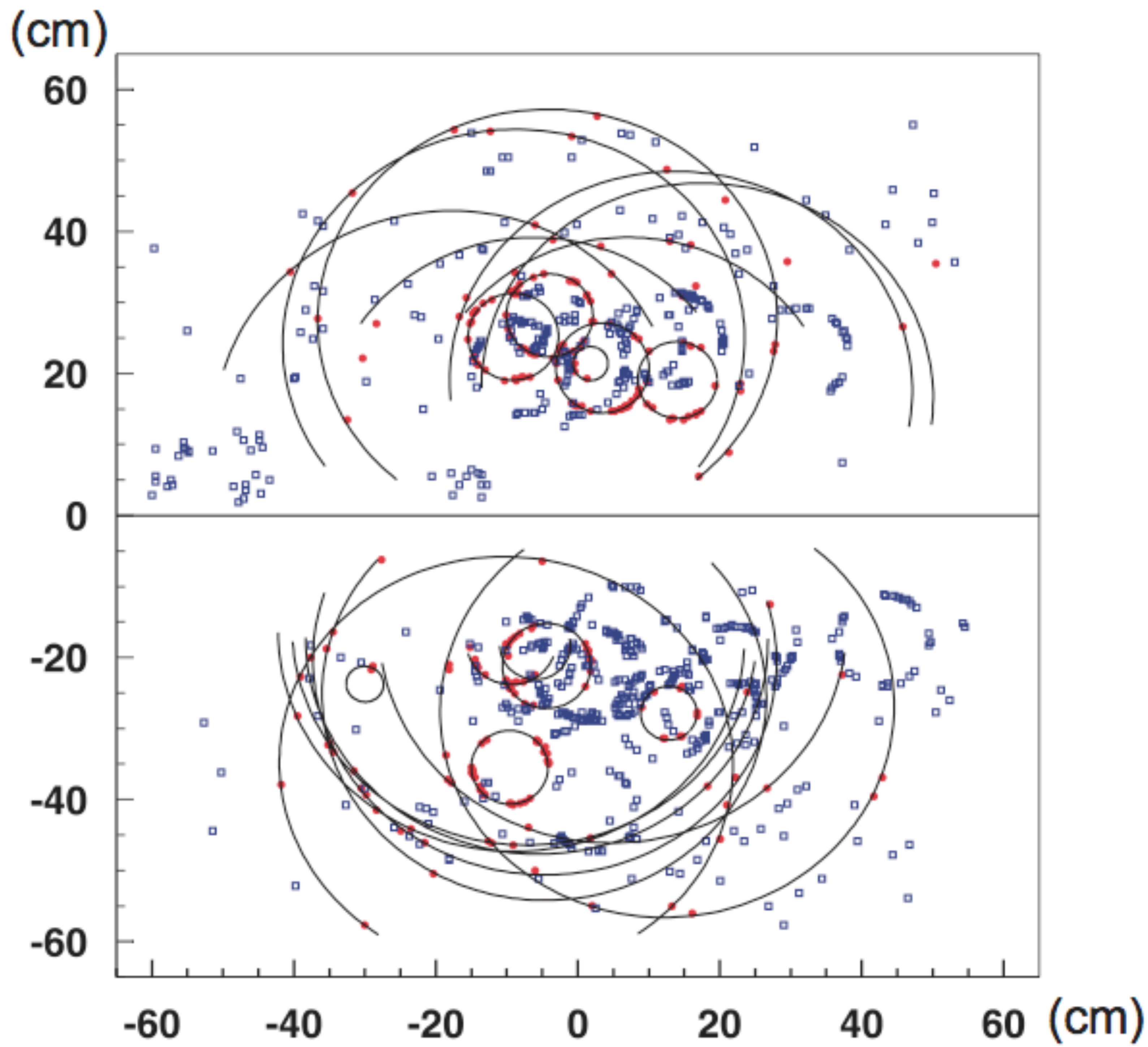
Cherenkov imaging detector
LHCb example







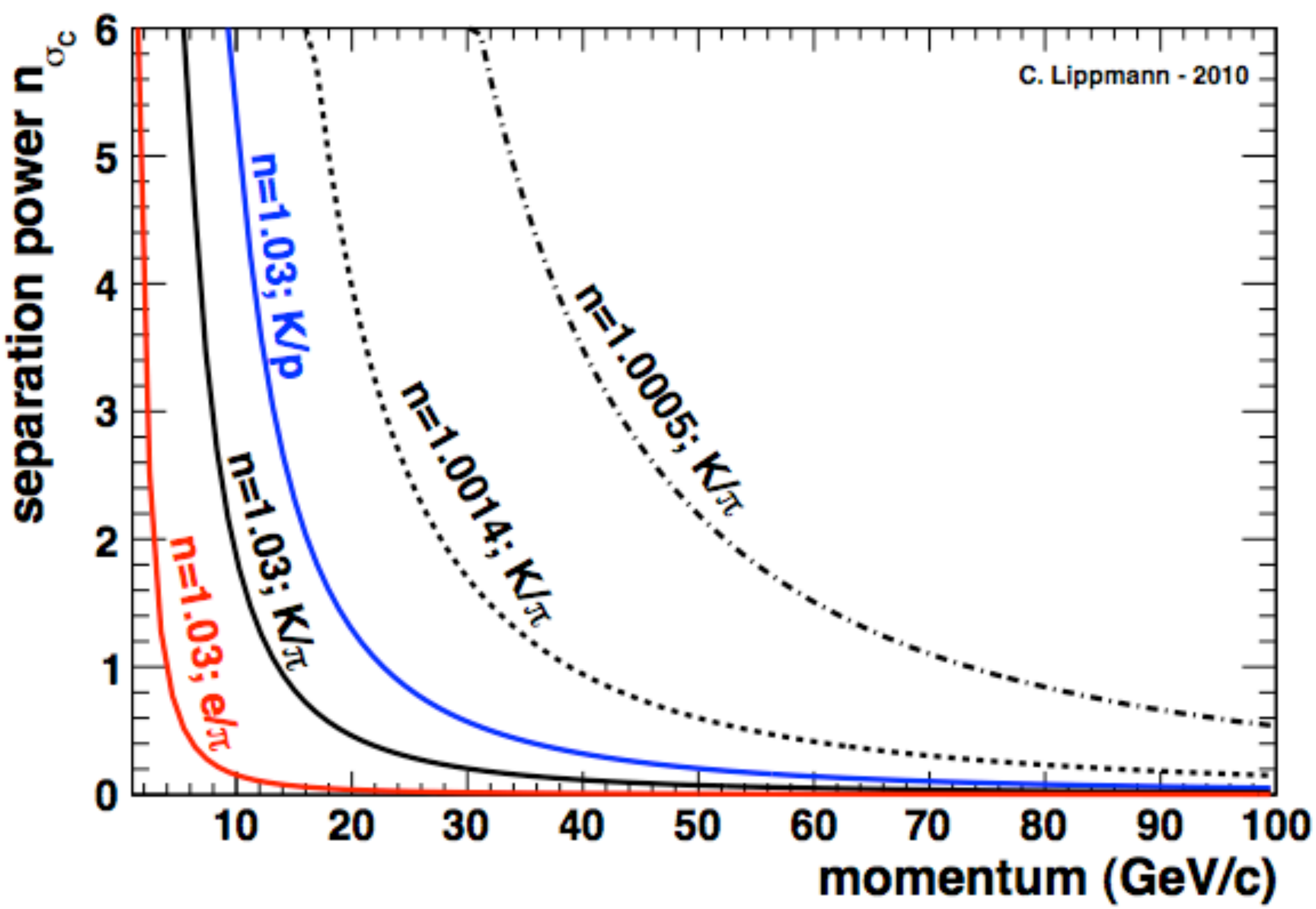
Θc (mrad)
250
200
150
100
50
e
μ
π
K
p
n=1.03
n=1.0014
π
K
n=1.0005
1
10
10^2
momentum (GeV/c)



(cm)
60
40
20
0
-20
-40
-60
-60
-40
-20
0
20
40
60
(cm)

Table 3: Some parameters of the LHCb RICH detectors. The measured single photoelectron angular resolutions [87] are for the preliminary alignment available from the first data sample with p–p collisions at $\sqrt{s} = 7$ TeV.

		RICH1		RICH2
		Silica aerogel	C_4F_{10}	CF_4
Momentum range [GeV/c]		≤ 10	$10 \lesssim p \lesssim 60$	$16 \lesssim p \lesssim 100$
Angular acceptance [mrad]	vertical	± 25 to ± 250		± 15 to ± 100
	horizontal	± 25 to ± 300		± 15 to ± 120
Radiator length [cm]		5	95	180
Refractive index n		1.03 (1.037)	1.0014	1.0005
Maximum Cherenkov angle [mrad]		242 (268)	53	32
Expected photon yield at $\beta \approx 1$		6.7	30.3	21.9
σ_{Θ_i} [mrad]	expected	2.6	1.57	0.67
	measured	~7.5	2.18	0.91



Need good software to reconstruct the Cherenkov cones for each charged particle

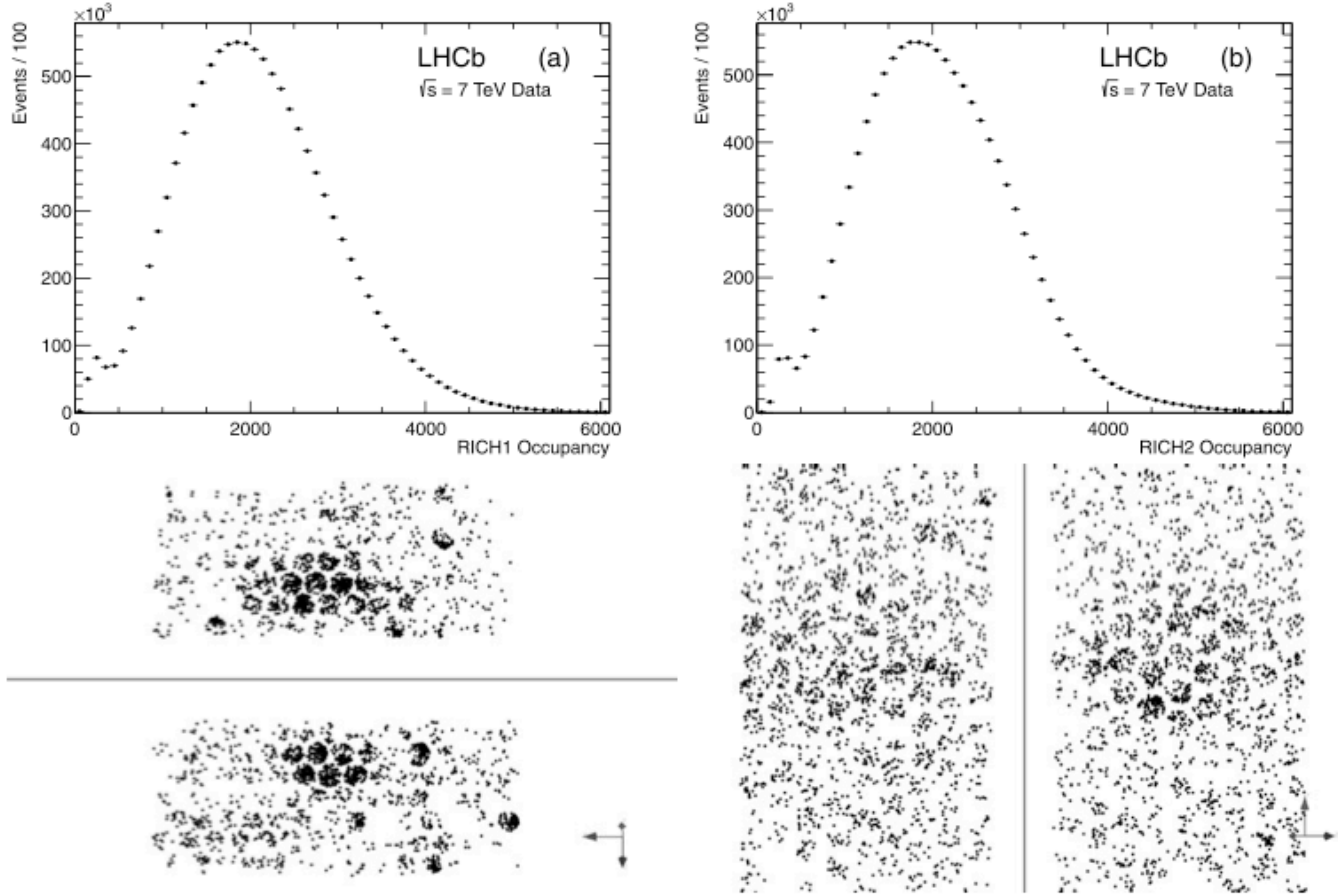
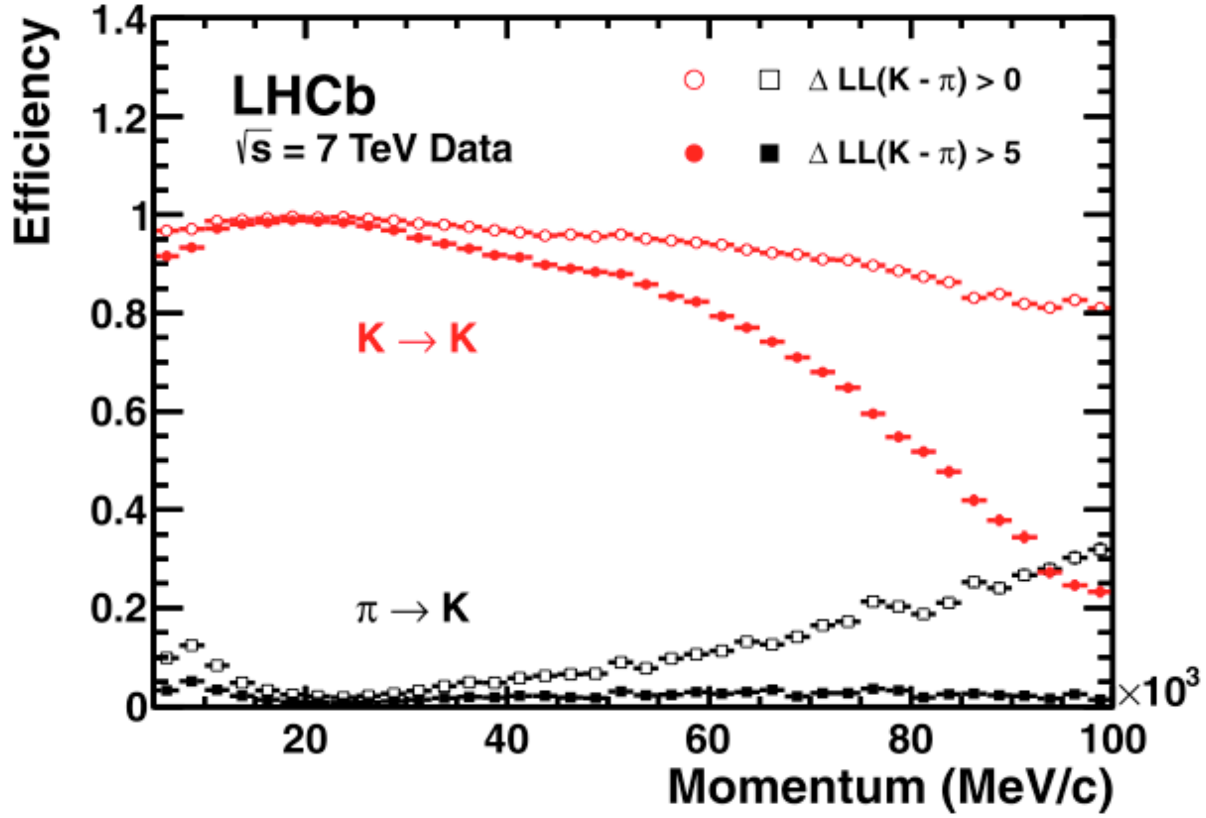


Fig. 13 Distribution of the number of pixel hits per event in (**a**) RICH 1 and (**b**) RICH 2. An example of a typical LHCb event as seen by the RICH detectors, is shown below the distributions. The *upper/lower* HPD panels in RICH 1 and the *left/right* panels in RICH 2 are shown separately



pi/kaon separation using RICH in LHCb

Impact on physics analysis

Eur. Phys. J. C (2013) 73:2431

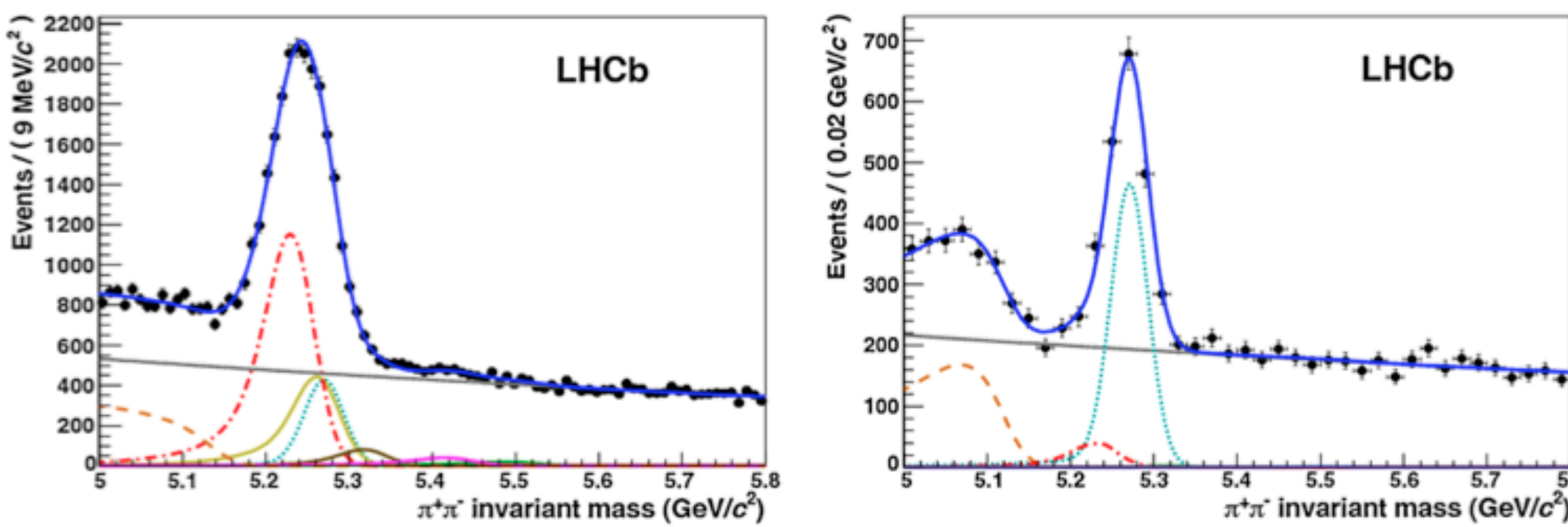
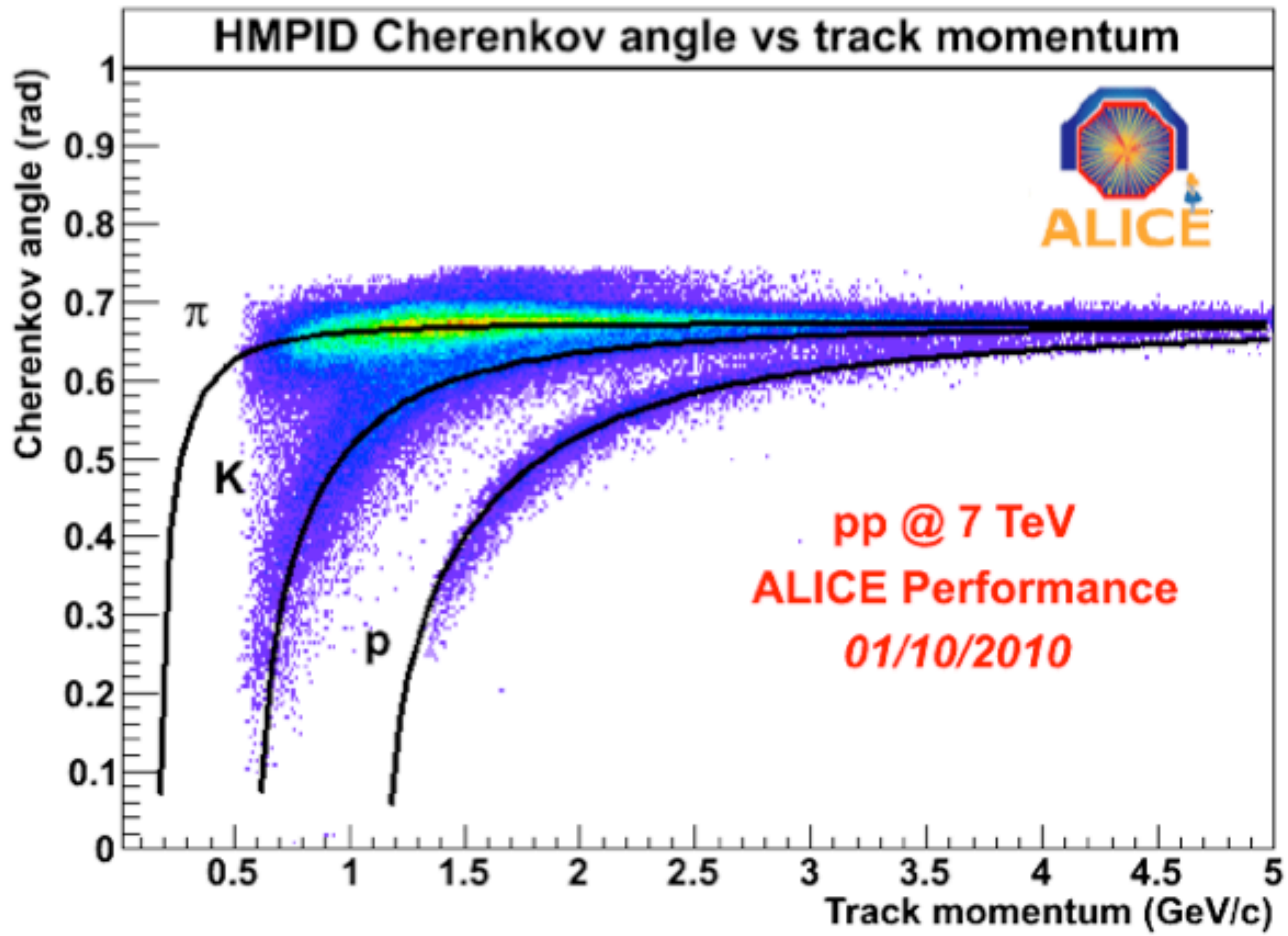


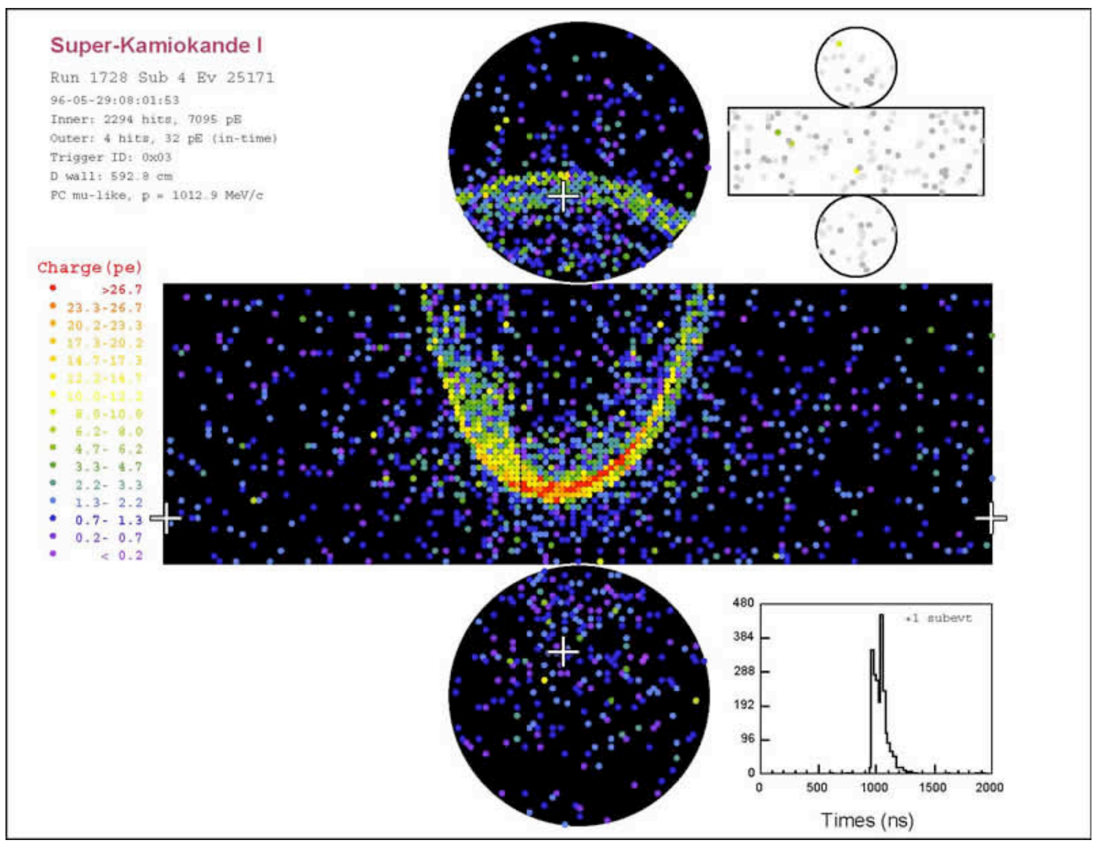
Fig. 2 Invariant mass distribution for $B \to h^+h^-$ decays [6] in the LHCb data before the use of the RICH information (*left*), and after applying RICH particle identification (*right*). The signal under study is the decay $B^0 \to \pi^+\pi^-$, represented by the turquoise *dotted line*. The contributions from different *b*-hadron decay modes ($B^0 \to K\pi$ *red dashed-dotted line*, $B^0 \to$ 3-body *orange dashed-dashed line*, $B_s \to KK$ *yellow line*, $B_s \to K\pi$ *brown line*, $\Lambda_b \to pK$ *purple line*, $\Lambda_b \to p\pi$ *green line*), are eliminated by positive identification of pions, kaons and protons and only the signal and two background contributions remain visible in the plot on the right. The *grey solid line* is the combinatorial background (Color figure online)

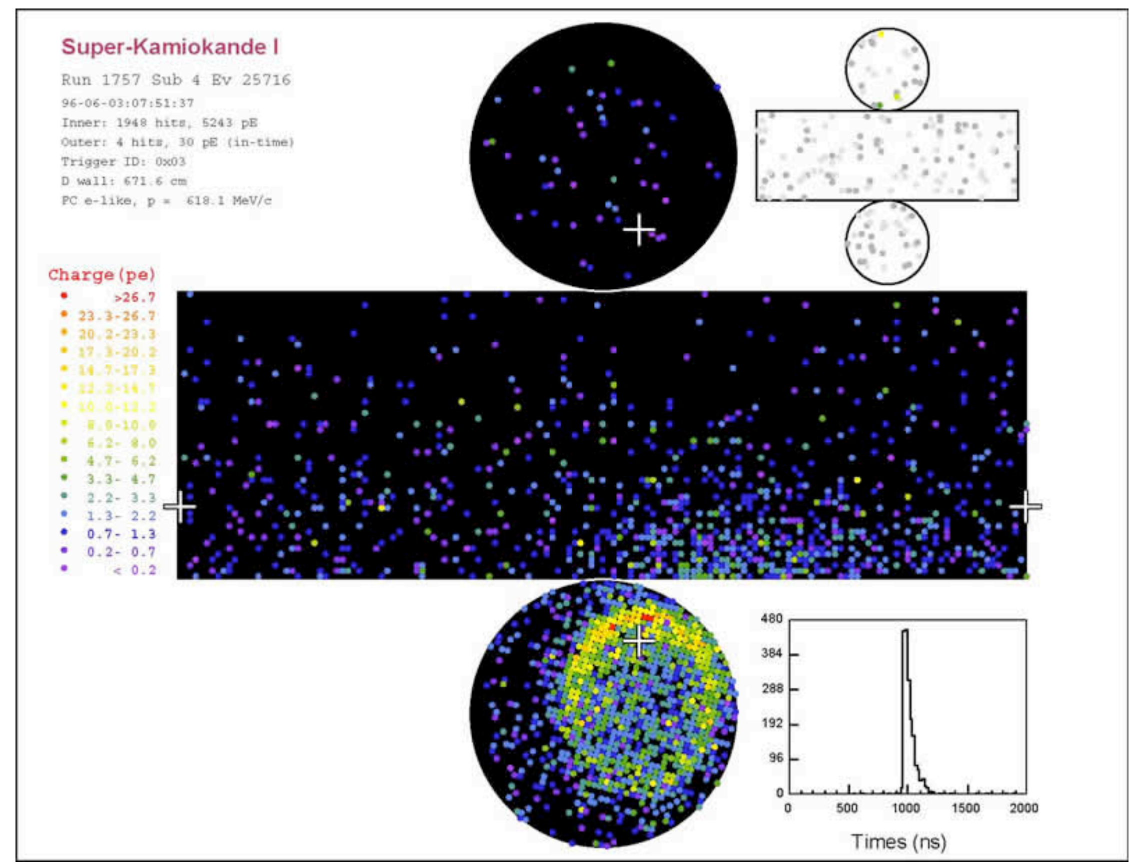
Cherenkov detector in ALICE



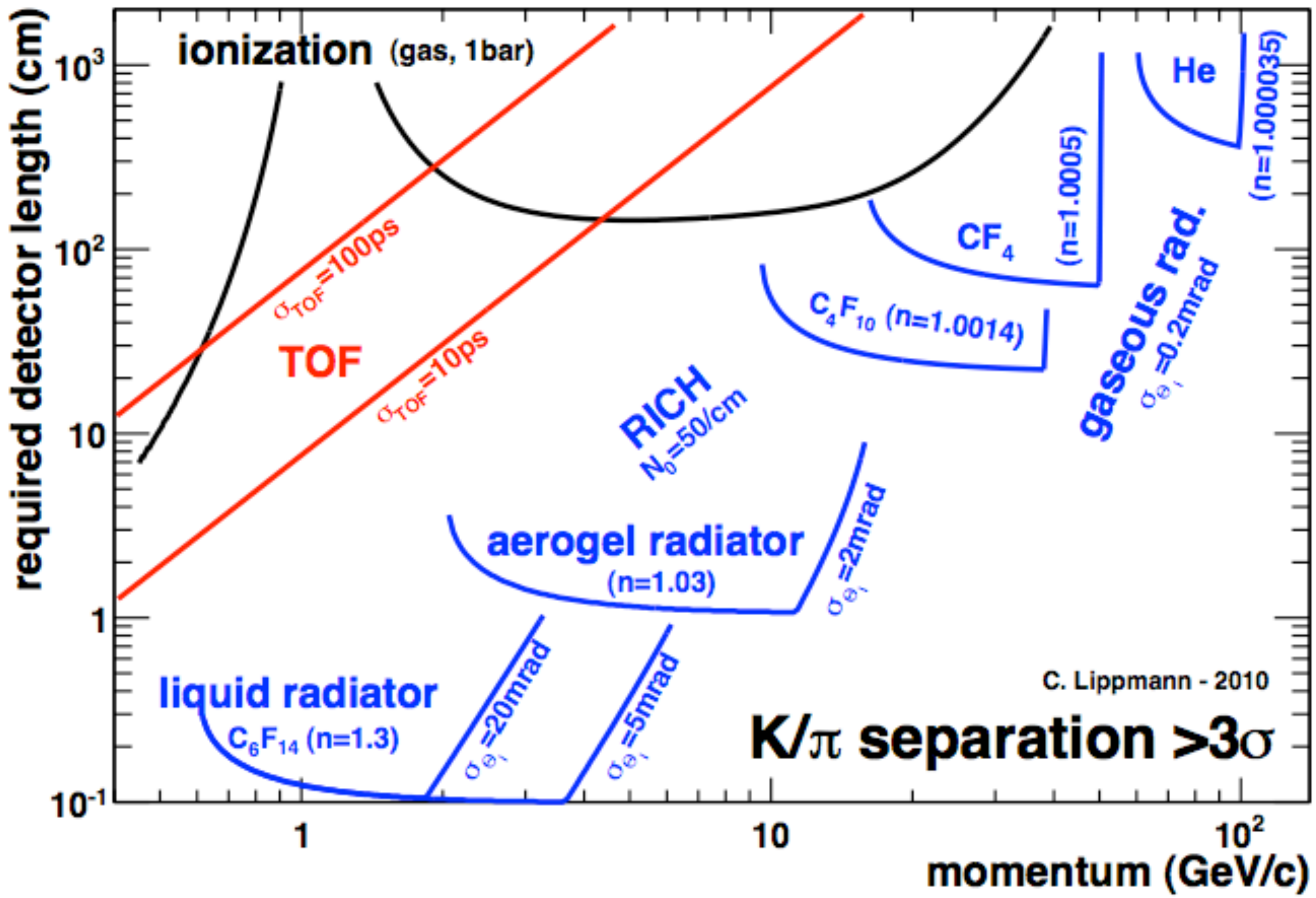
Application of Cherenkov for neutrino detector

neutrino interaction in water produces muon or electron
which are above Cherenkov threshold
Light is detected by photo multipliers around the water tank





Comparison of different techniques



Transition radiation

When charge z.e crosses boundary vacuum/medium

$$I = \frac{1}{3} \alpha z^2 \gamma \hbar \omega_p$$

$$\hbar \omega_p = \sqrt{4\pi N_e r_e^3 \frac{m_e c^2}{\alpha}} = \sqrt{\rho/(g/cm^3) \langle \frac{Z}{A} \rangle} \times 28.81 eV$$

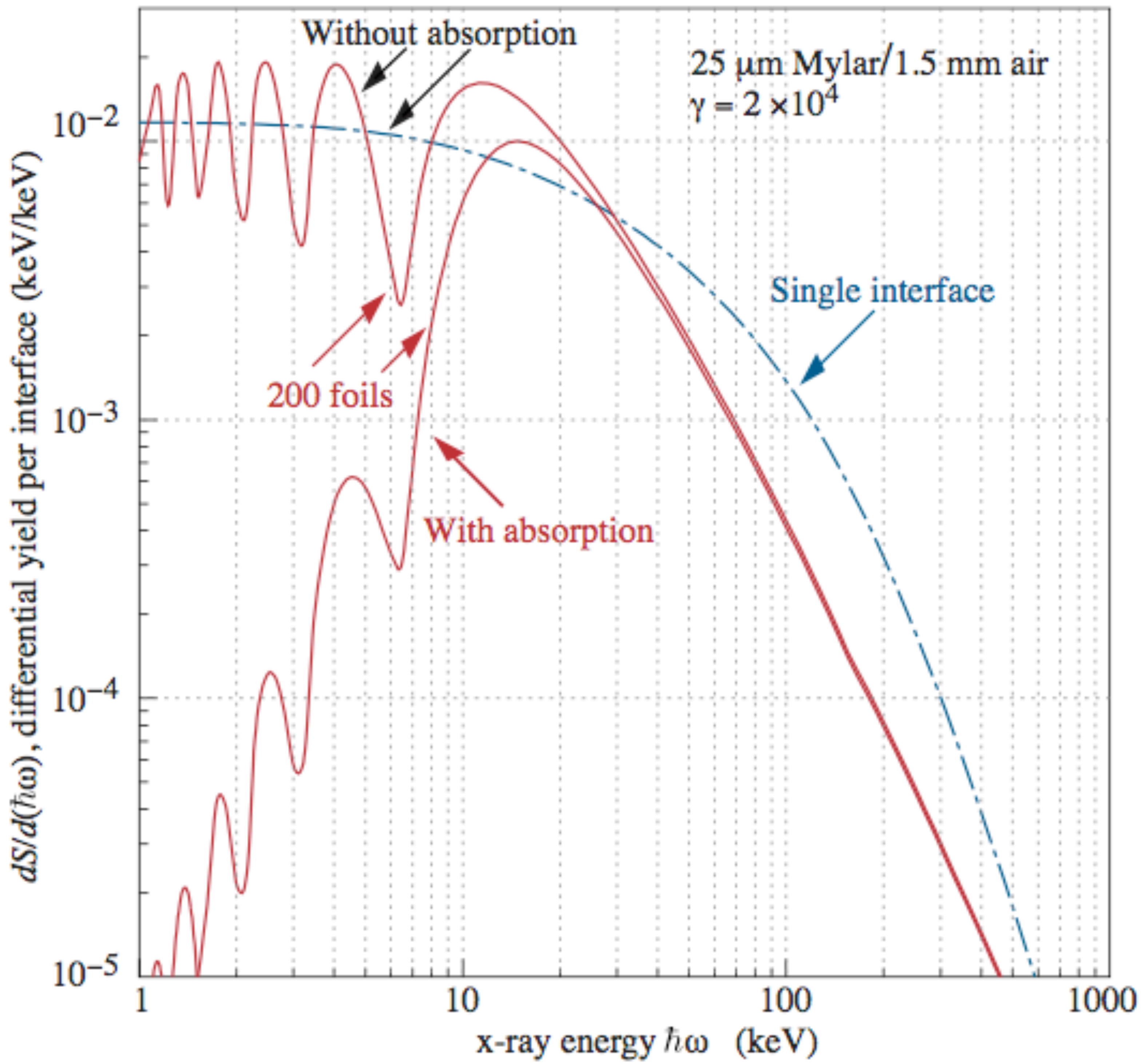
Typical values $\hbar\omega_p \sim 20$ eV (0.7 for air)

Half energy between 0.1 and 1. $\gamma \hbar \omega_p$

Typically $\sim 0.005 \gamma$ with $\hbar\omega > 0.1 \gamma \hbar \omega_p$

Formation length $\sim$ tens of μm

Needs many interfaces to increase photon yield



X-rays detected for instance by photo-electric effect in high Z material like Xenon gas

=> Detector consists of radiator + photon detector

Photon interaction in matter

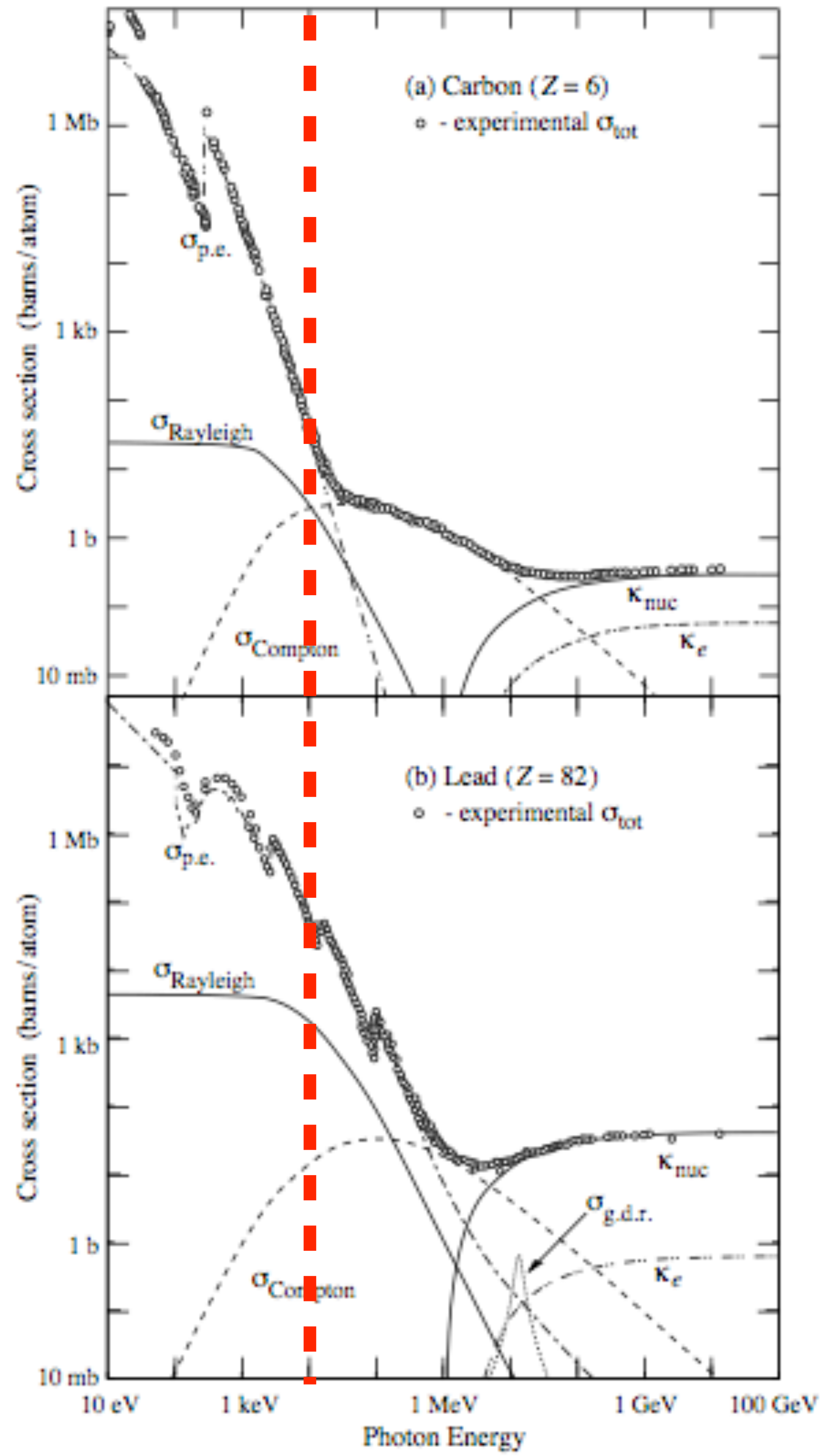
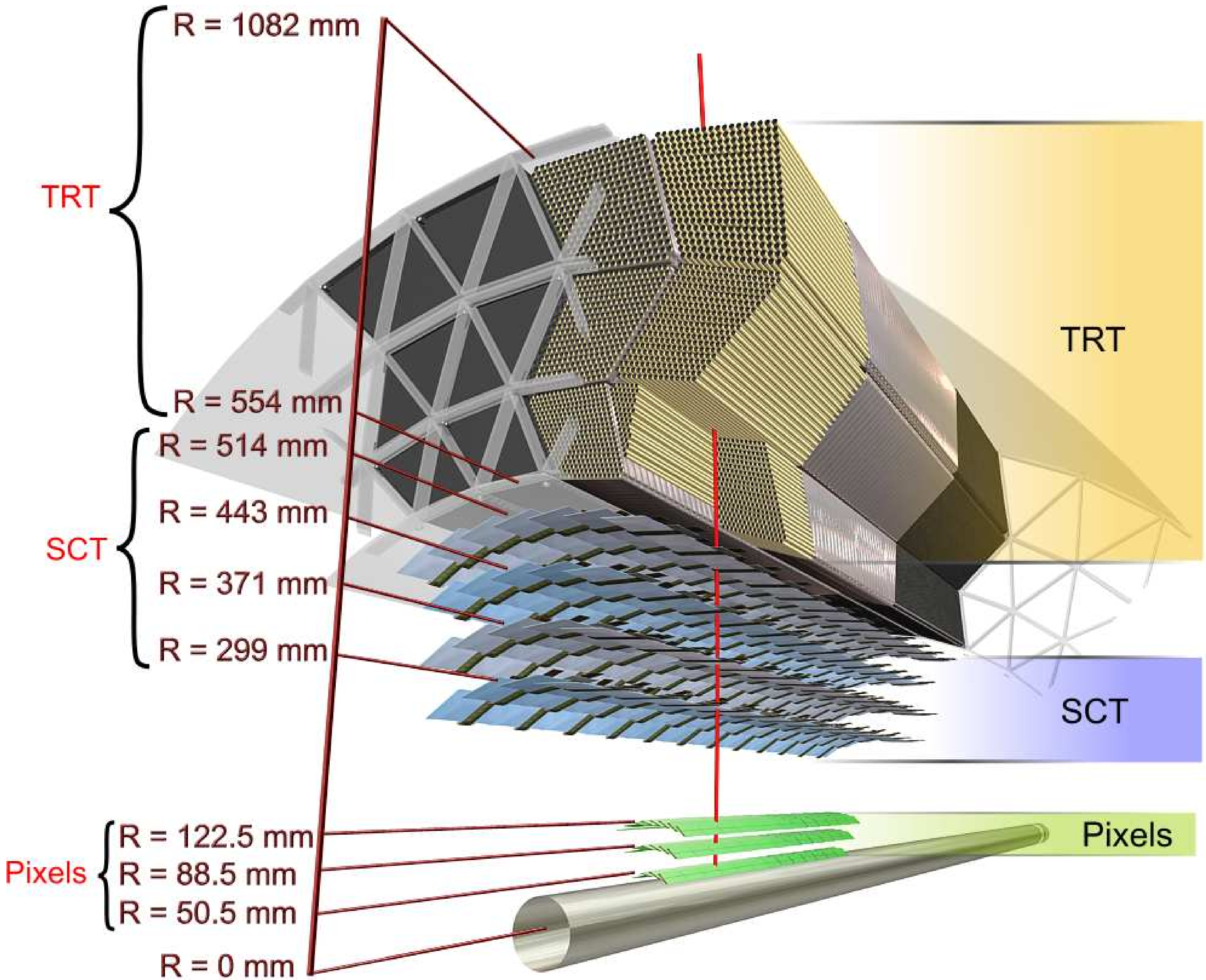
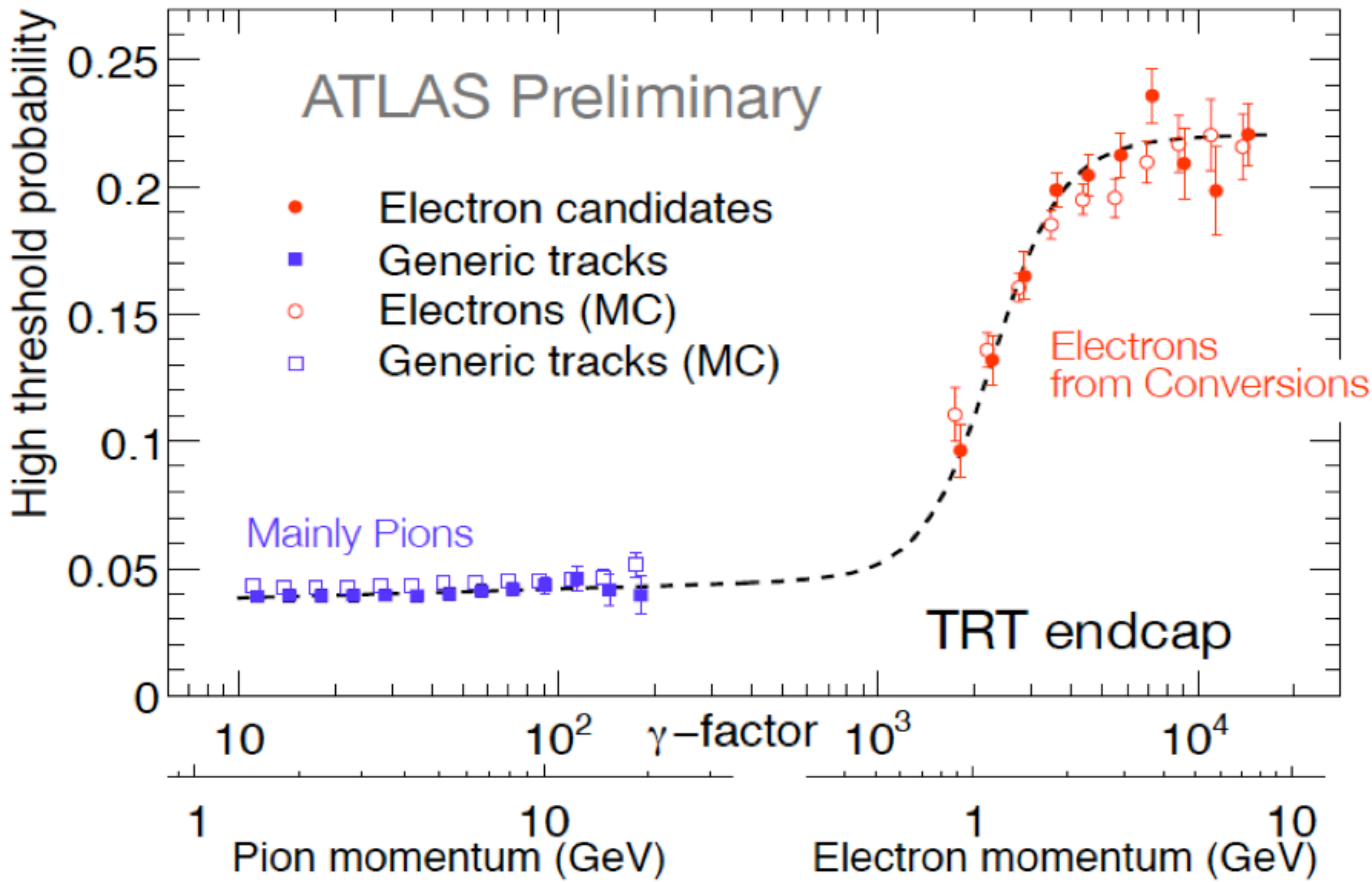


Figure 31.15: Photon total cross sections as a function of energy in carbon and lead, showing the contributions of different processes [51]:

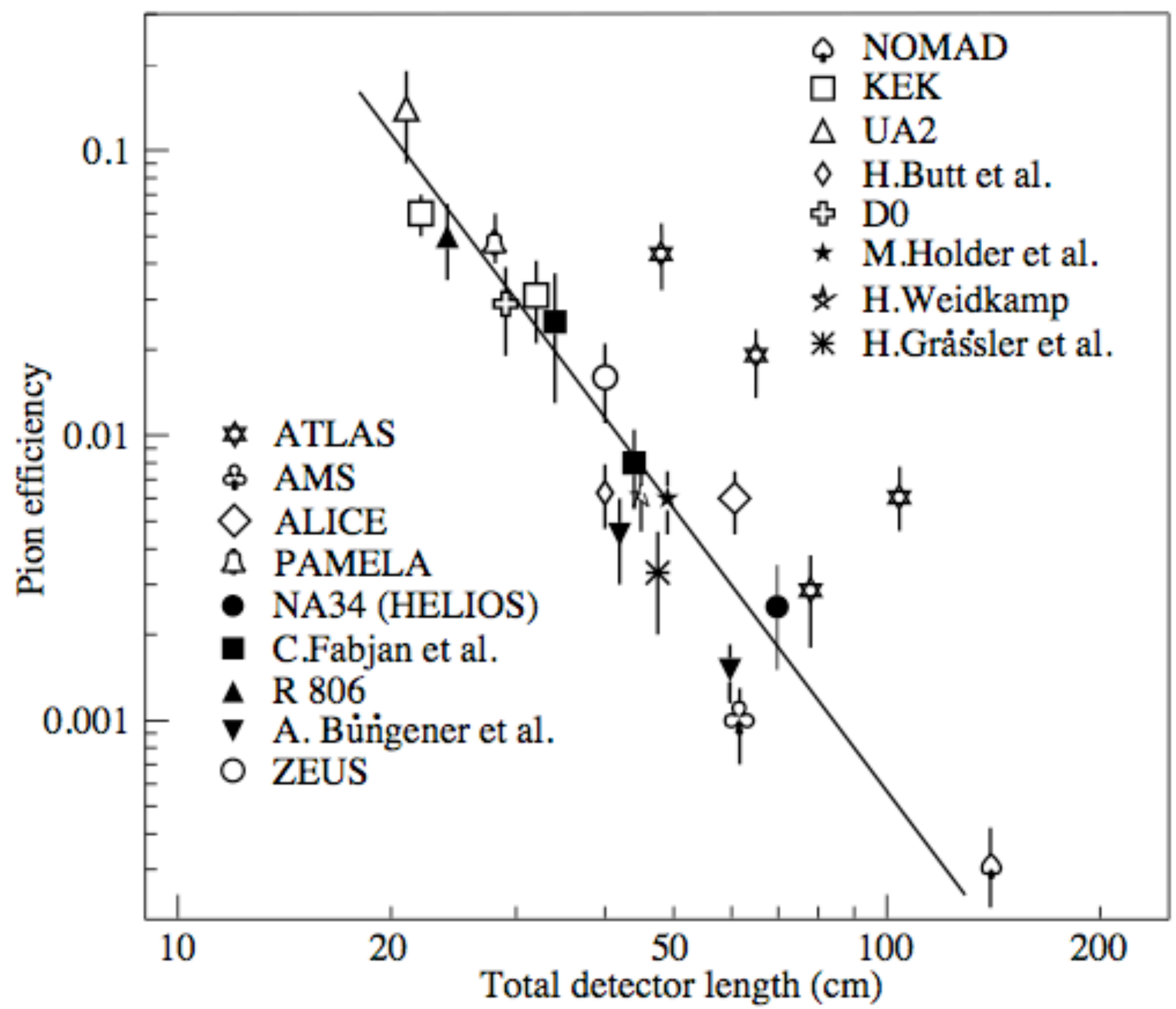
$\sigma_{\text{p.e.}}$ = Atomic photoelectric effect (electron ejection, photon absorption)
σ_{Rayleigh} = Rayleigh (coherent) scattering–atom neither ionized nor excited
σ_{Compton} = Incoherent scattering (Compton scattering off an electron)
κ_{nuc} = Pair production, nuclear field
κ_e = Pair production, electron field
$\sigma_{\text{g.d.r.}}$ = Photonuclear interactions, most notably the Giant Dipole Resonance [52]. In these interactions, the target nucleus is broken up.

Radiator = polypropylene foils
Detector = Straws with wire in the middle containing Xe
(to absorb X-rays)
Edeposited ~2 keV from ionization, ~8-10 KeV from TR photons





ATLAS Preliminary
High threshold probability
0.25
0.2
0.15
0.1
0.05
0
Electron candidates
Generic tracks
Electrons (MC)
Generic tracks (MC)
Electrons
from Conversions
Mainly Pions
TRT endcap
10
10^2
γ−factor
10^3
10^4
1
10
Pion momentum (GeV)
1
10
Electron momentum (GeV)

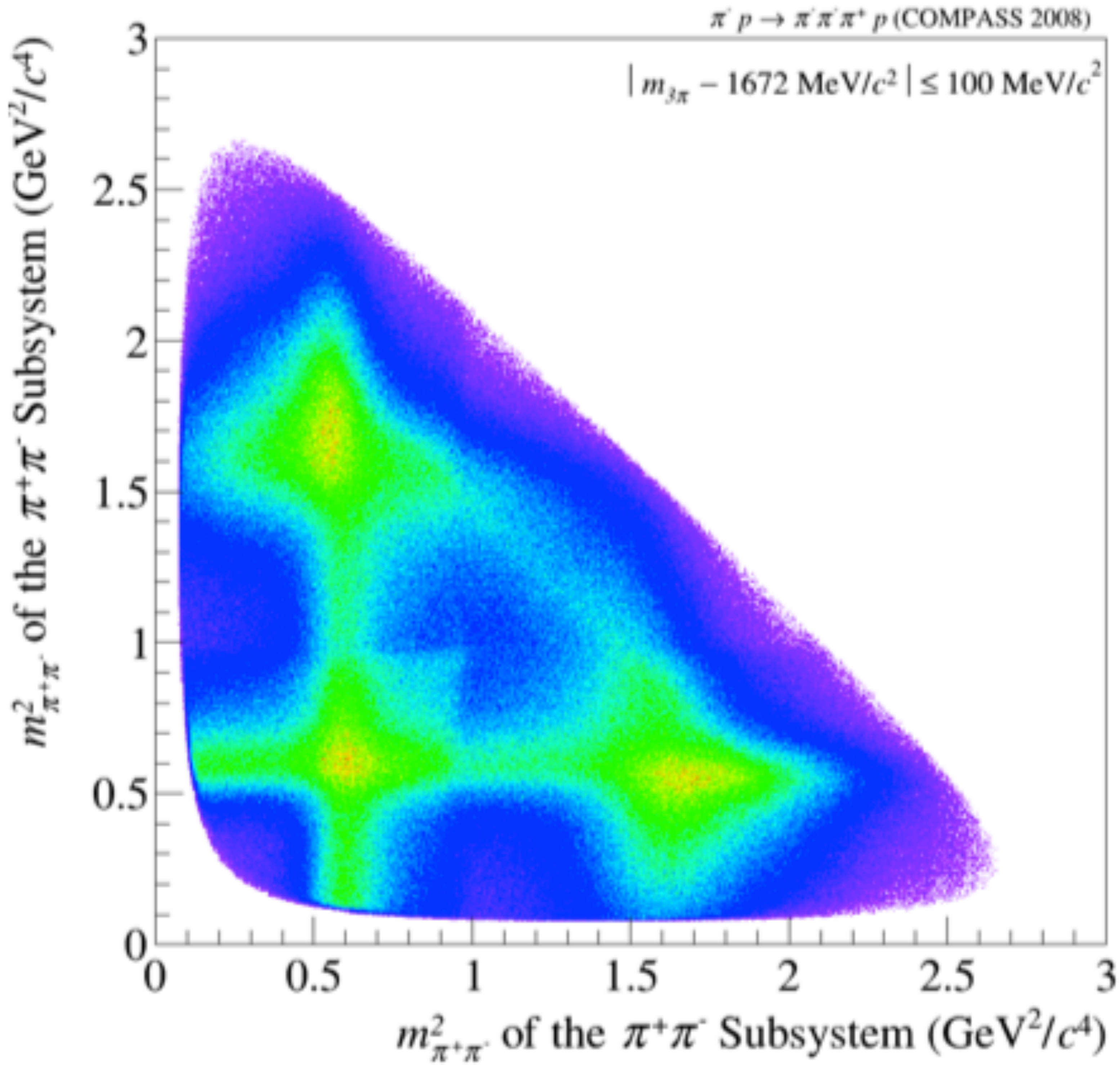


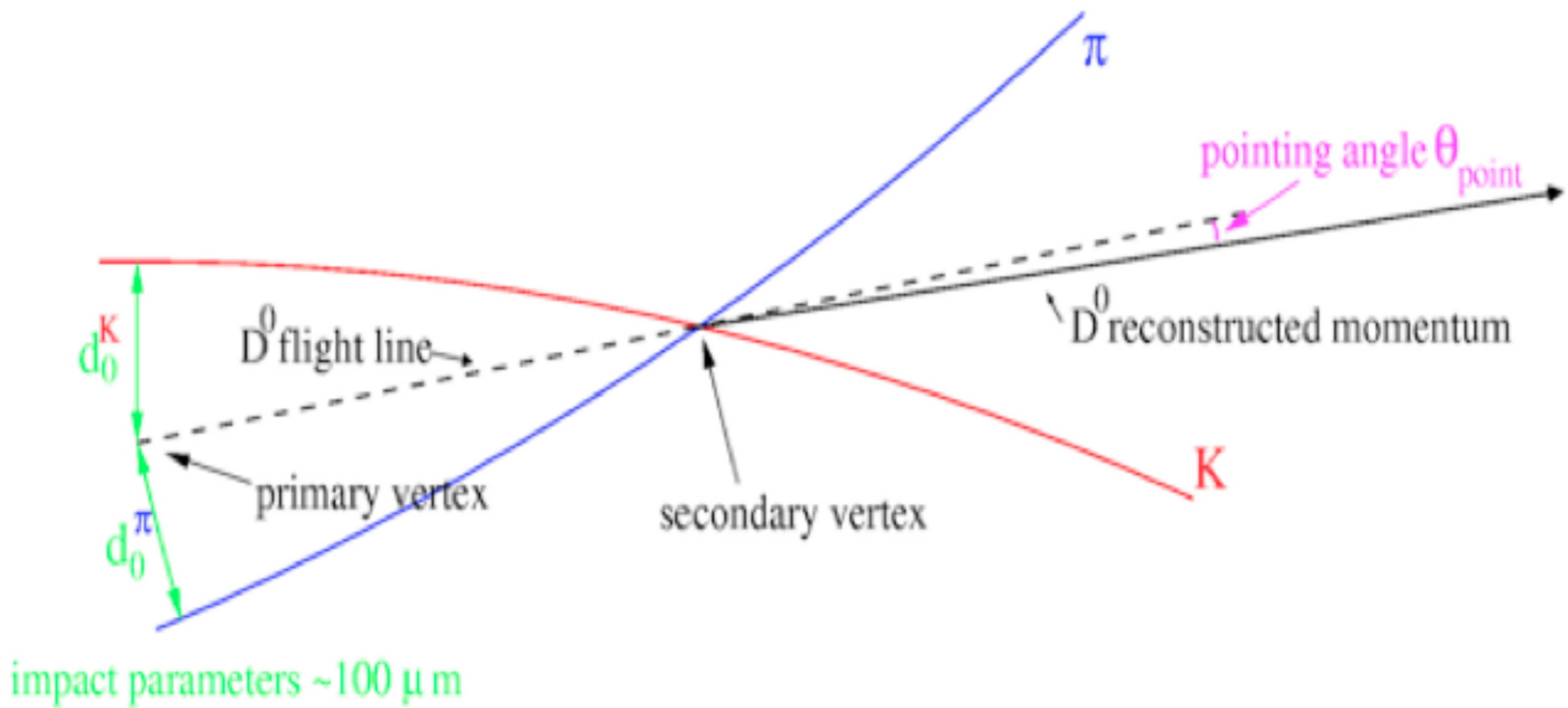
NOMAD
KEK
UA2
H.Butt et al.
D0
M.Holder et al.
H.Weidkamp
H.Gråssler et al.
ATLAS
AMS
ALICE
PAMELA
NA34 (HELIOS)
C.Fabjan et al.
R 806
A. Büngener et al.
ZEUS
Pion efficiency
0.1
0.01
0.001
10
20
50
100
200
Total detector length (cm)

Reconstruction of particle decay

- Useful for short lived particles
  - very short lived => use invariant mass of daughter particles
    - Examples are Ks-> pi+pi-, J/psi-> mu+ mu-, W,Z decays, etc..
  - not so short lived => can measure distance between production and decay positions:
    - tau lepton
    - B-hadron

Exploiting kinematic information from Dalitz plots





lifetimes: D0: 4.10^{-13}s, B0d 1.5 10^{-12}s, tau: 2.9 10^{-13}s

Decay length beta.gamma.c.tau => beta.gamma. 450 microns for B0d

Impact parameter ~ (c.tau)

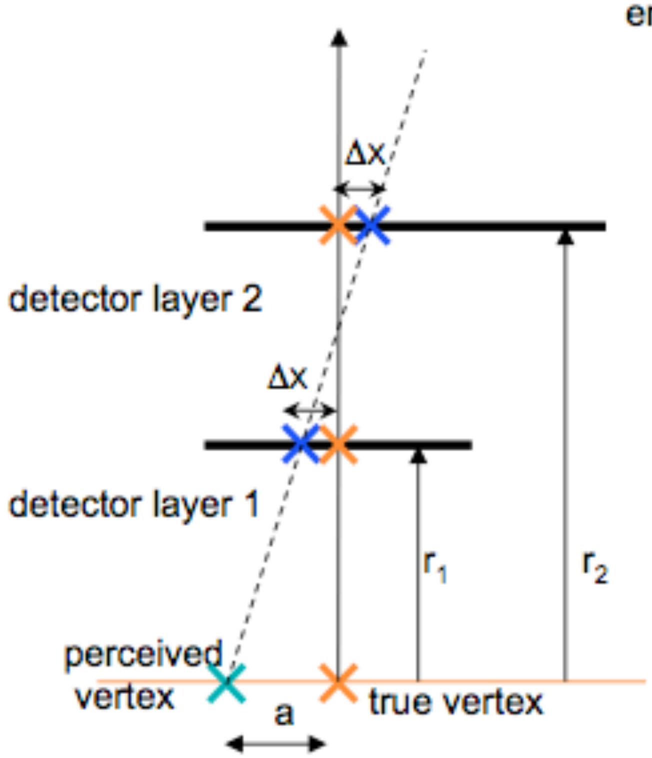
Vertex projection from two points: a simplified approach (telescope equation)

pointing resolution = (a $\oplus$ b) μm

a: from detector position error

b: from coulomb scattering

$$a = \Delta x \bullet \sqrt{\frac{r_2^2 + r_1^2}{(r_2 - r_1)^2}}$$



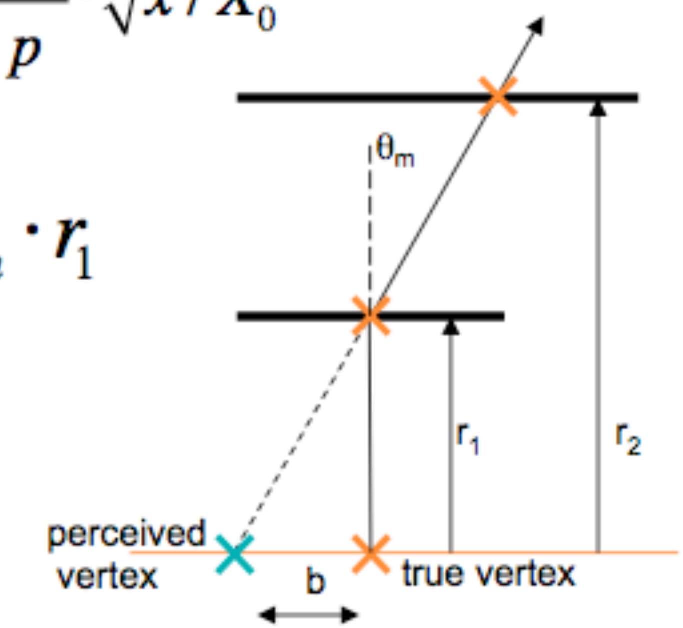
Detector Granularity, minimize Δx:

e.g. 50um pixel and r_2 very large compared to r_1

→ a=Δx=50/$\sqrt{12}$ = 15um

$$\theta_m = \frac{13.6 Mev}{\beta \cdot c \cdot p} \cdot \sqrt{x / X_0}$$

$$b = \theta_m \cdot r_1$$

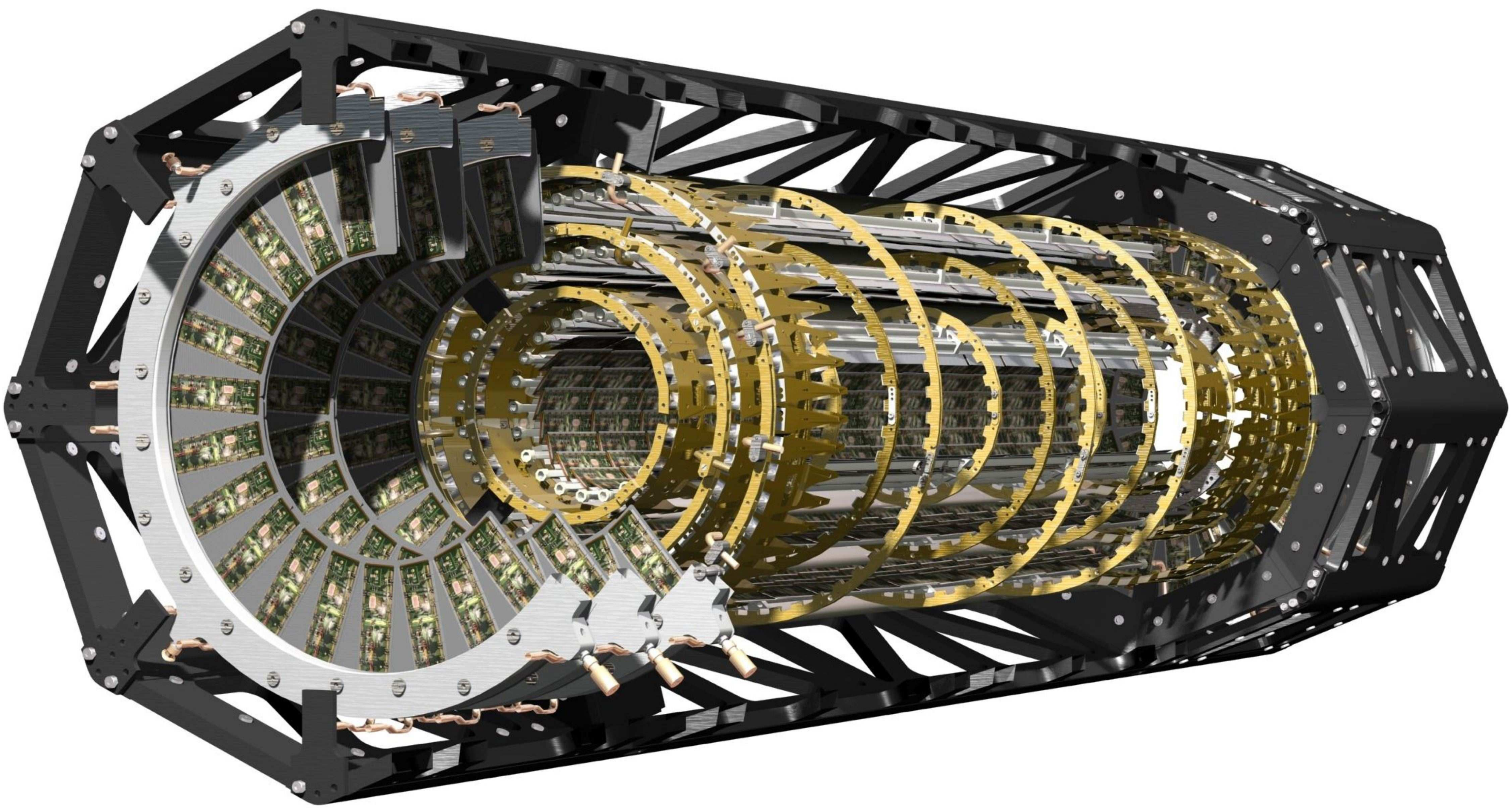


First layer as close as possible to the vertex and First layer with minimal amount of material.

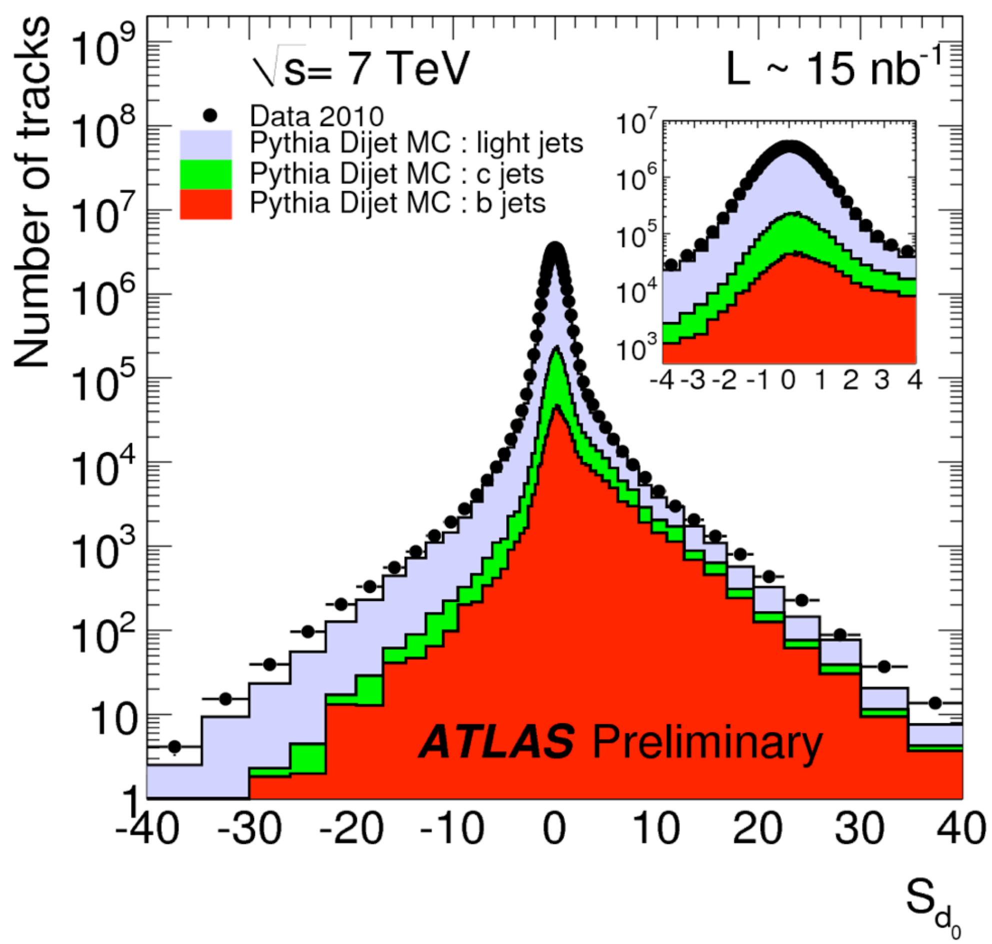
e.g. x/X_0 = 0.0114, r_1= 39mm

→ b= 57um for p=1GeV/c

Example of ATLAS pixel silicon detector

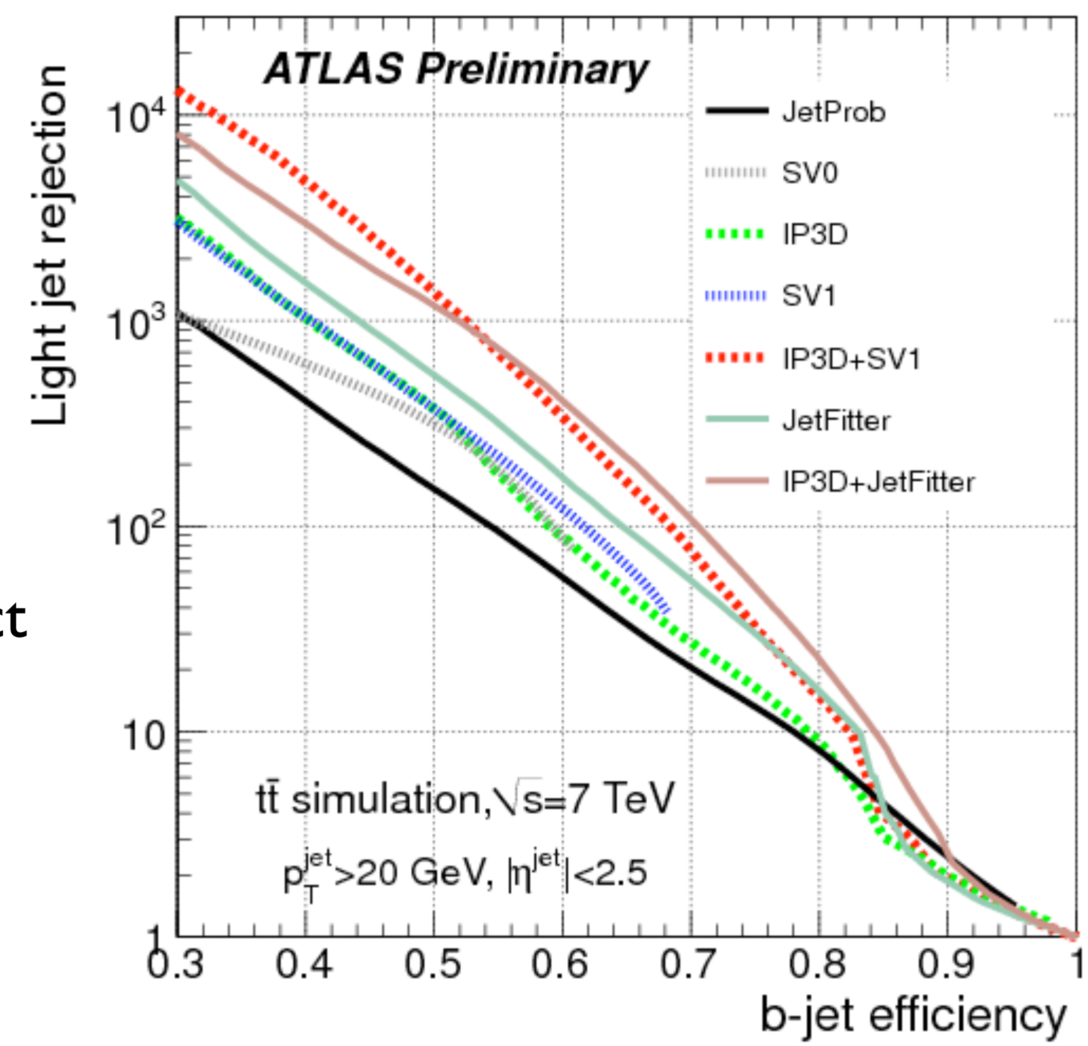


b-tagging performances



Track impact parameter/error

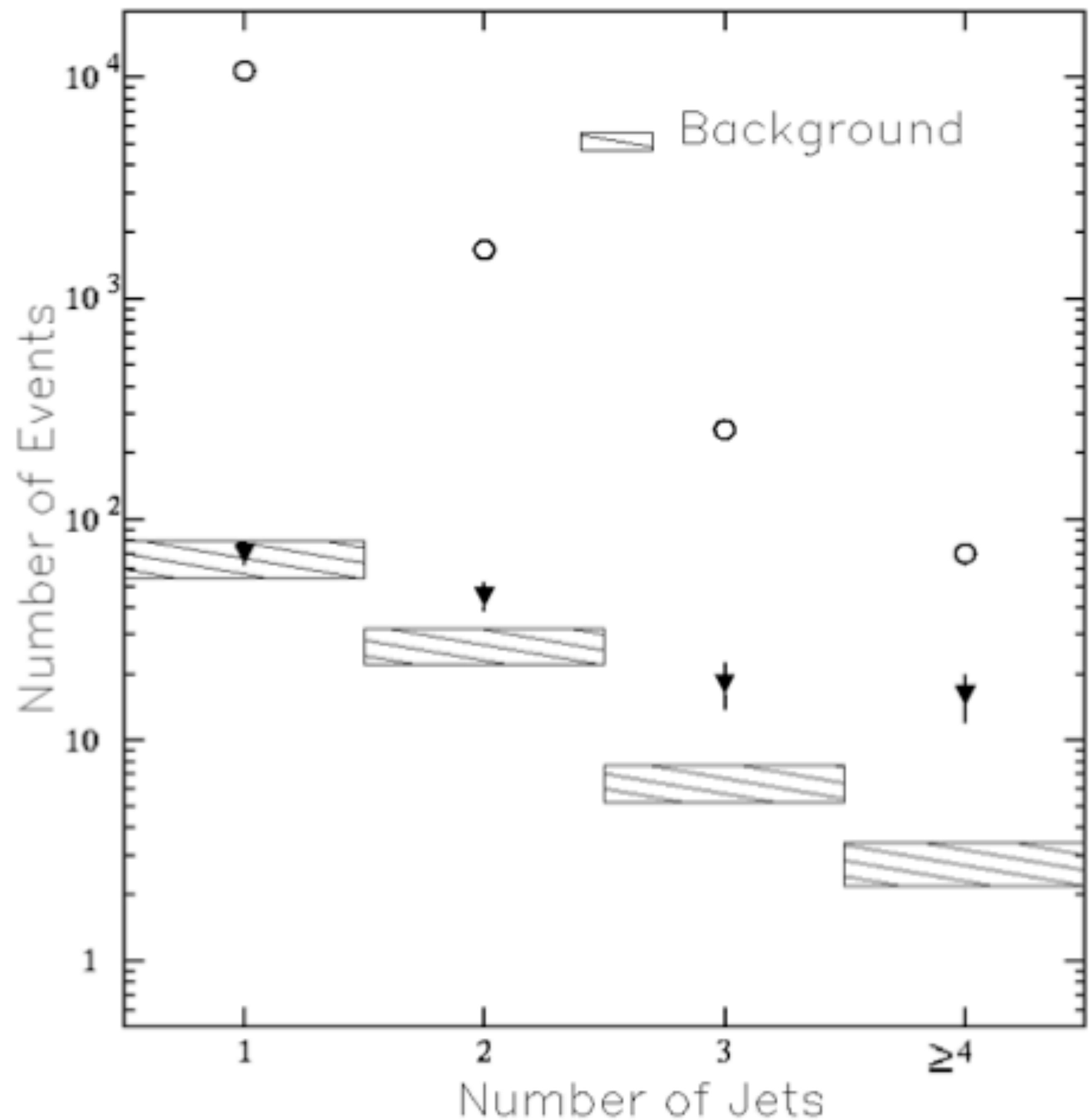
Algorithms combining impact parameter information + secondary vertex reconstruction



Example of b-tagging usage for top quark discovery

Signal t tbar -> W W b bbar, one W->lepton, one W->jets

Background: W(->lepton)+ jet
Only a small amount of these jets have b quarks.



e + 4 jet event

40758_44414

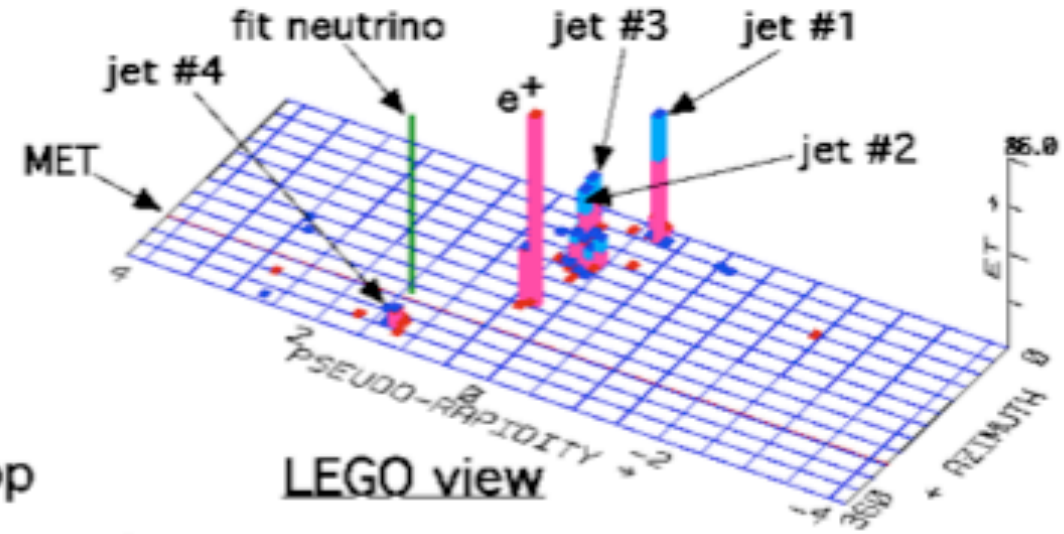
24-September, 1992

TWO jets tagged by SVX

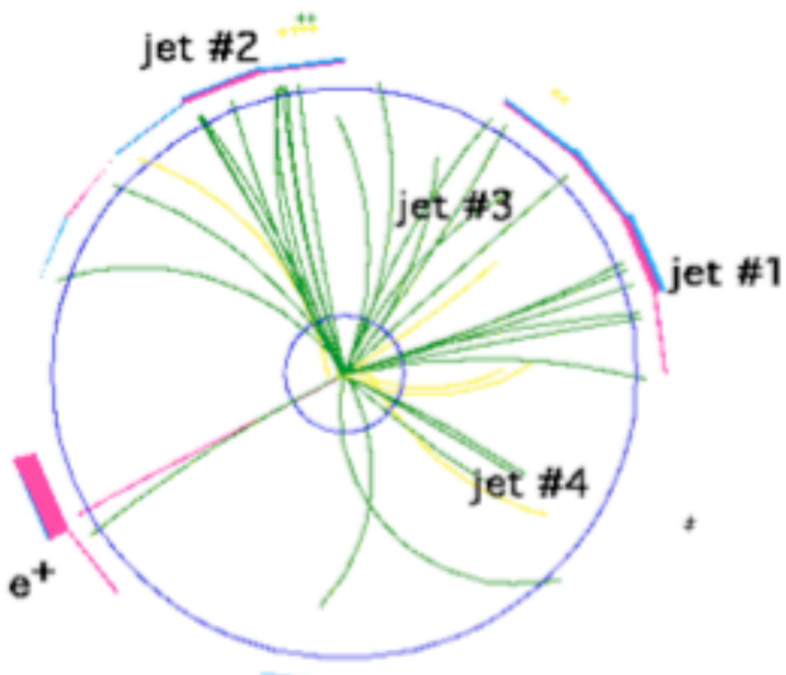
fit top mass is 170 +- 10 GeV

e^+, Missing E_t, jet #4 from top

jets 1,2,3 from top (2&3 from W)

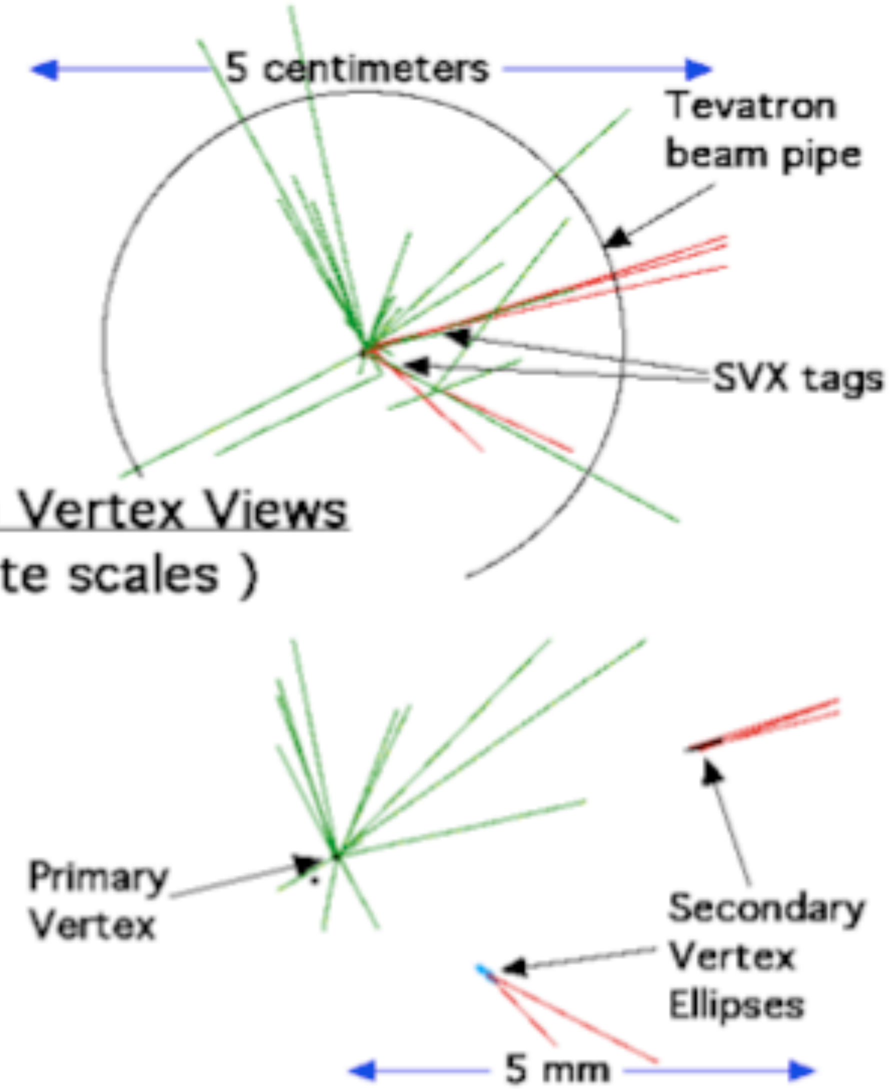


LEGO view



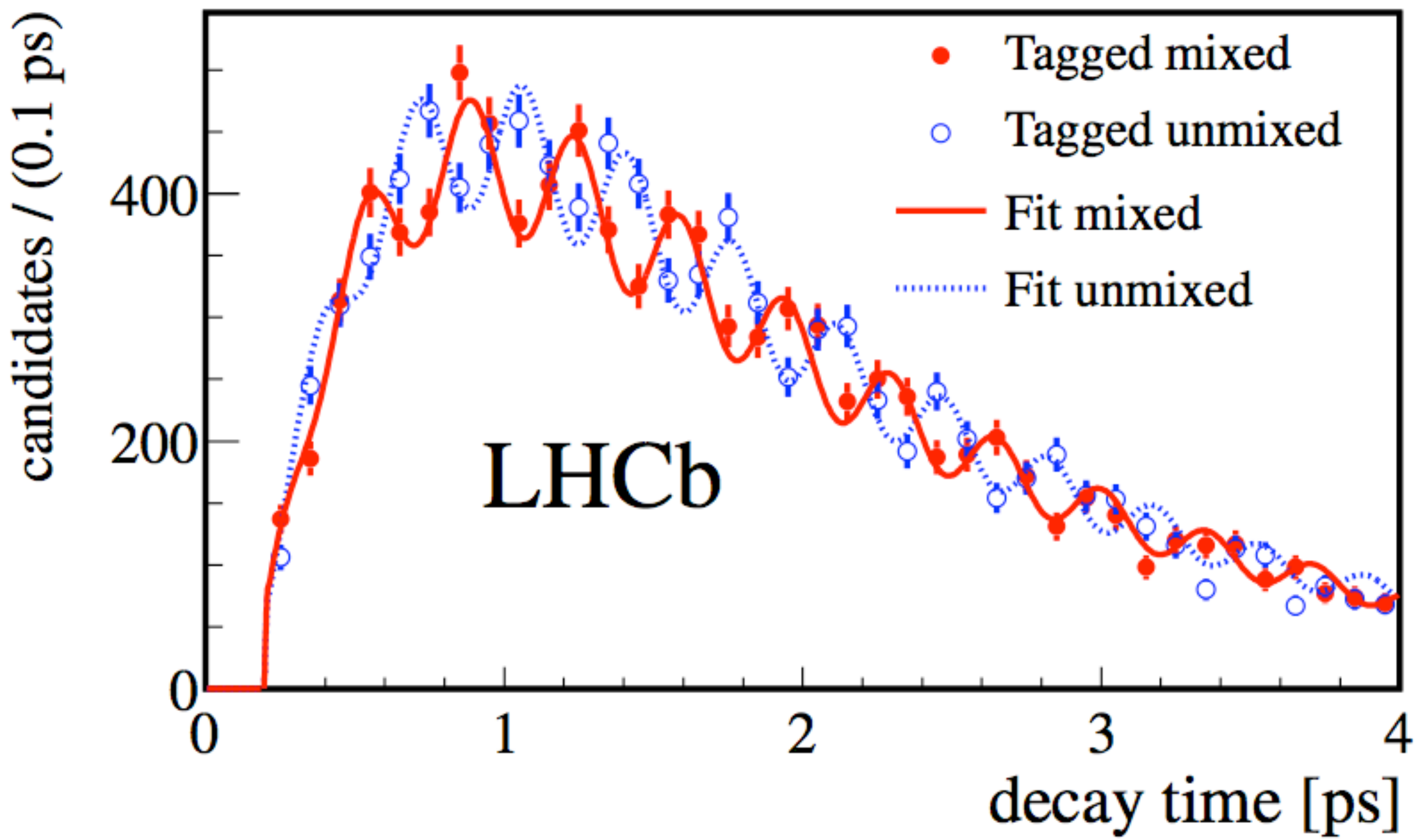
3 meters

Tracking View



Two Vertex Views
(note scales)

Example of application of particle ID and secondary vertex : Bs mixing measurement



Particle Flow techniques in collider experiments

- Different particles species are measured more accurately with different techniques
  - What is the most precise technique for E=100 GeV electron energy measurement in a LHC experiment ?
  - What is the most precise technique to measure a few GeV charged pion ?
  - What is the most precise technique to measure a 5 GeV K0L ?
  - How can one separate particles from different interactions in the same bunch crossing at the LHC ?

- Charged particle momentum measurement

Detector resolution: $\sigma\left(\frac{1}{p_t}\right) = \text{Cte} \Rightarrow \frac{\sigma_{p_t}}{p_t} = a \cdot p_t \qquad a \propto \frac{1}{BL^2}$

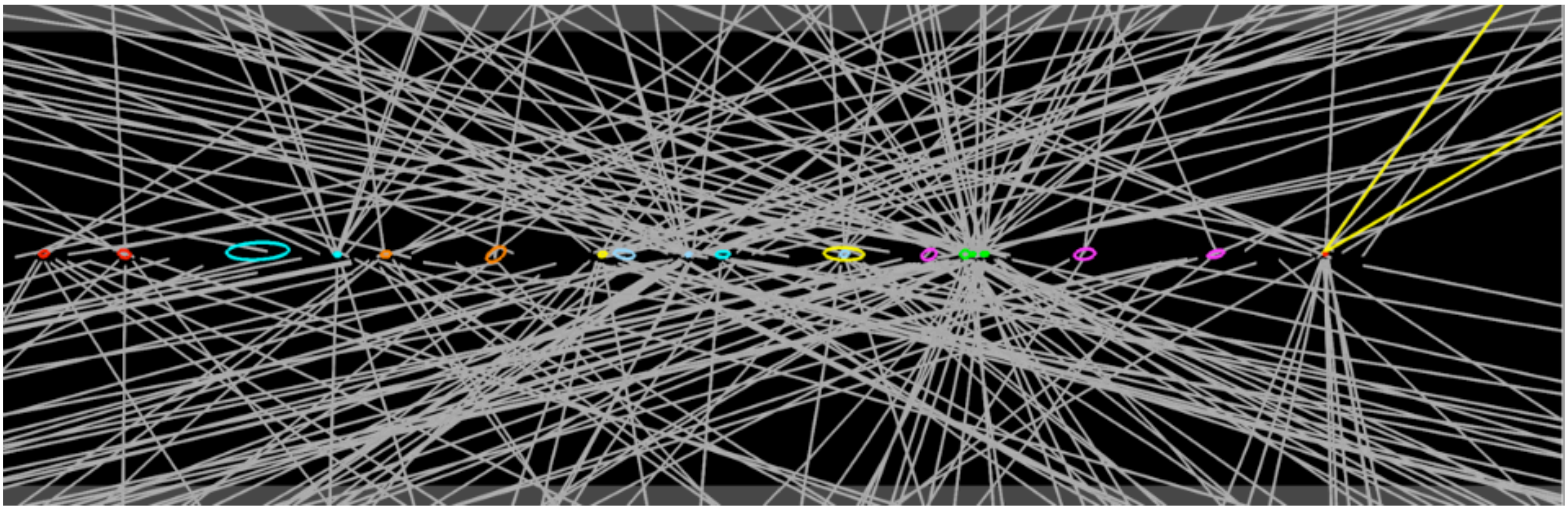
Multiple scattering: $\sigma_\theta \propto \frac{13.6\ \text{MeV}}{p} \sqrt{\frac{x}{X_0}}$

$\sigma\left(\frac{1}{p_t}\right) \propto \sigma_\theta \Rightarrow \frac{\sigma_{p_t}}{p_t} = b \qquad \left(b \propto \frac{1}{B}\right)$

- Calorimetric energy measurement

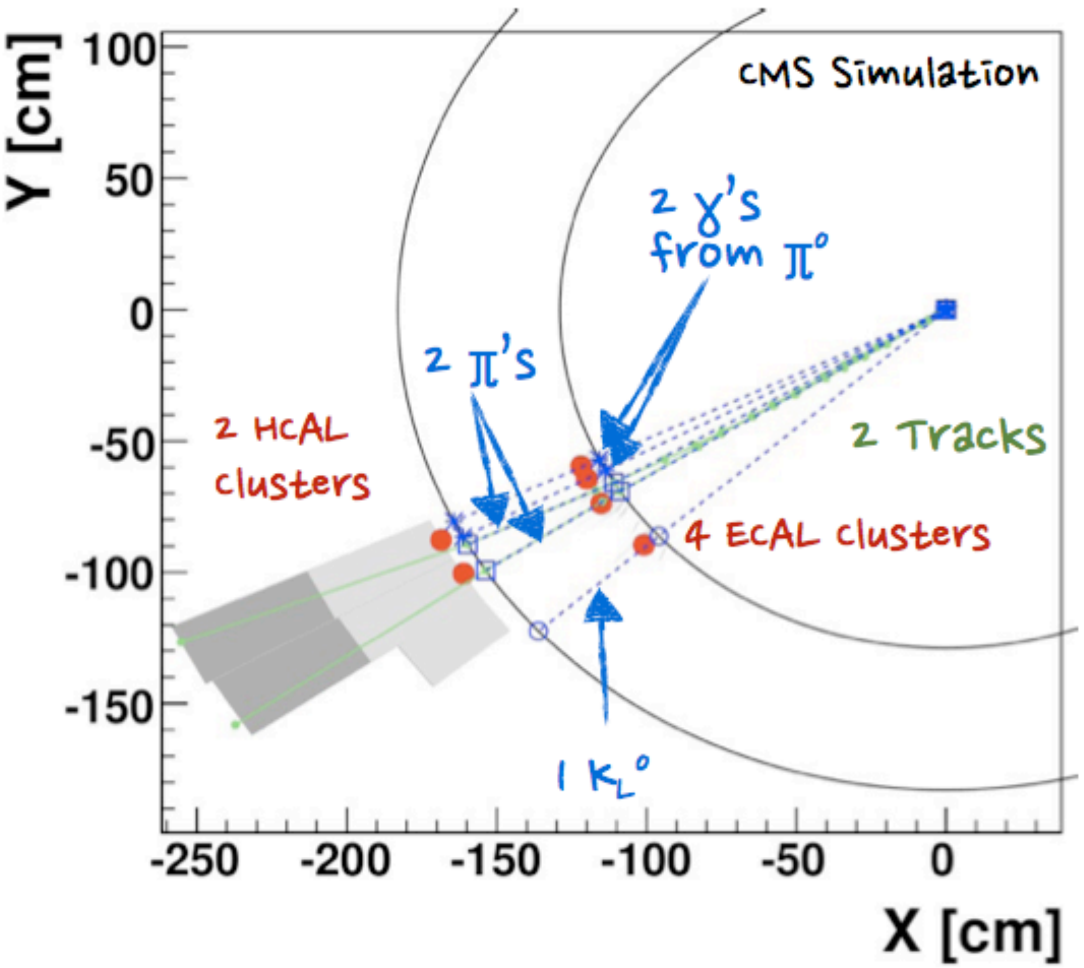
$$\frac{\sigma_E}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

Also have to deal with pileup interactions



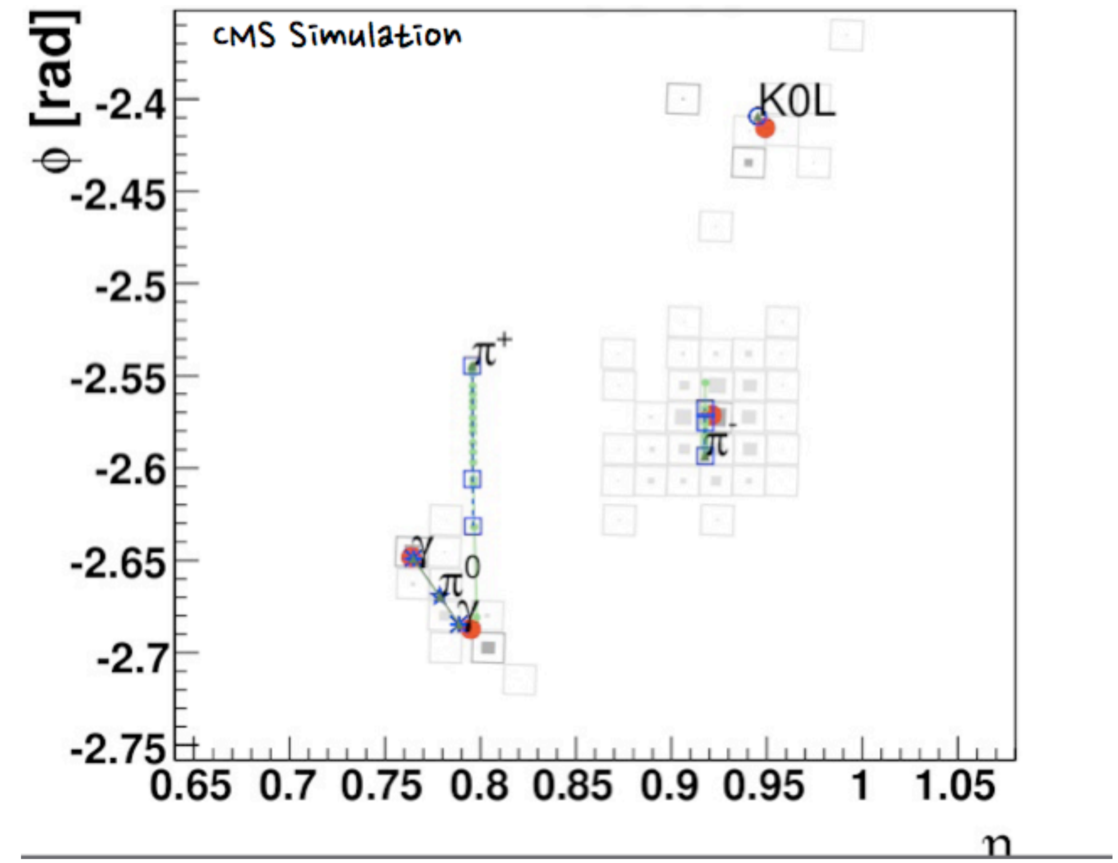
Can be distinguished for charged tracks
but not easily for calorimeter energy deposits

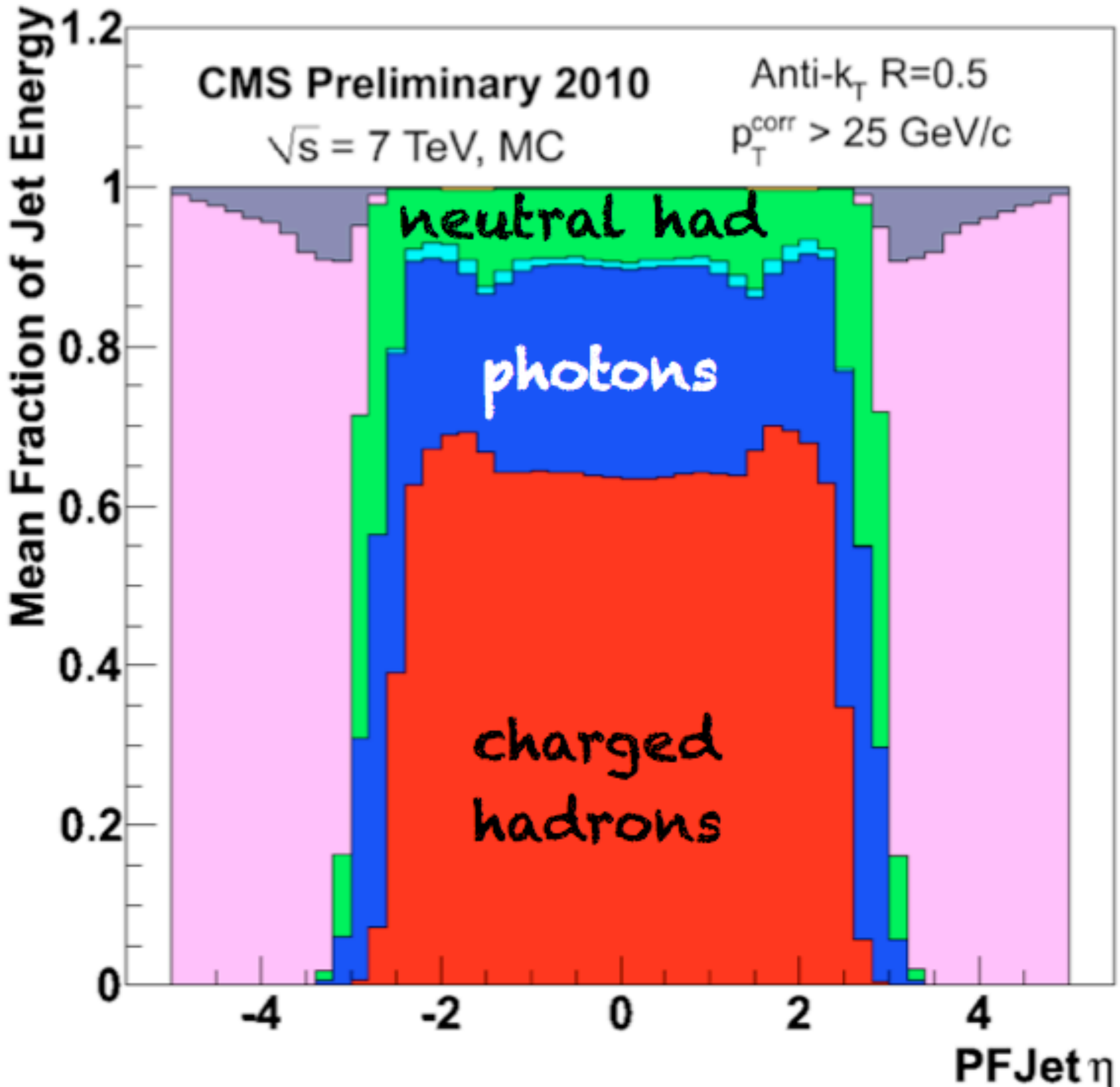
Particle flow principle

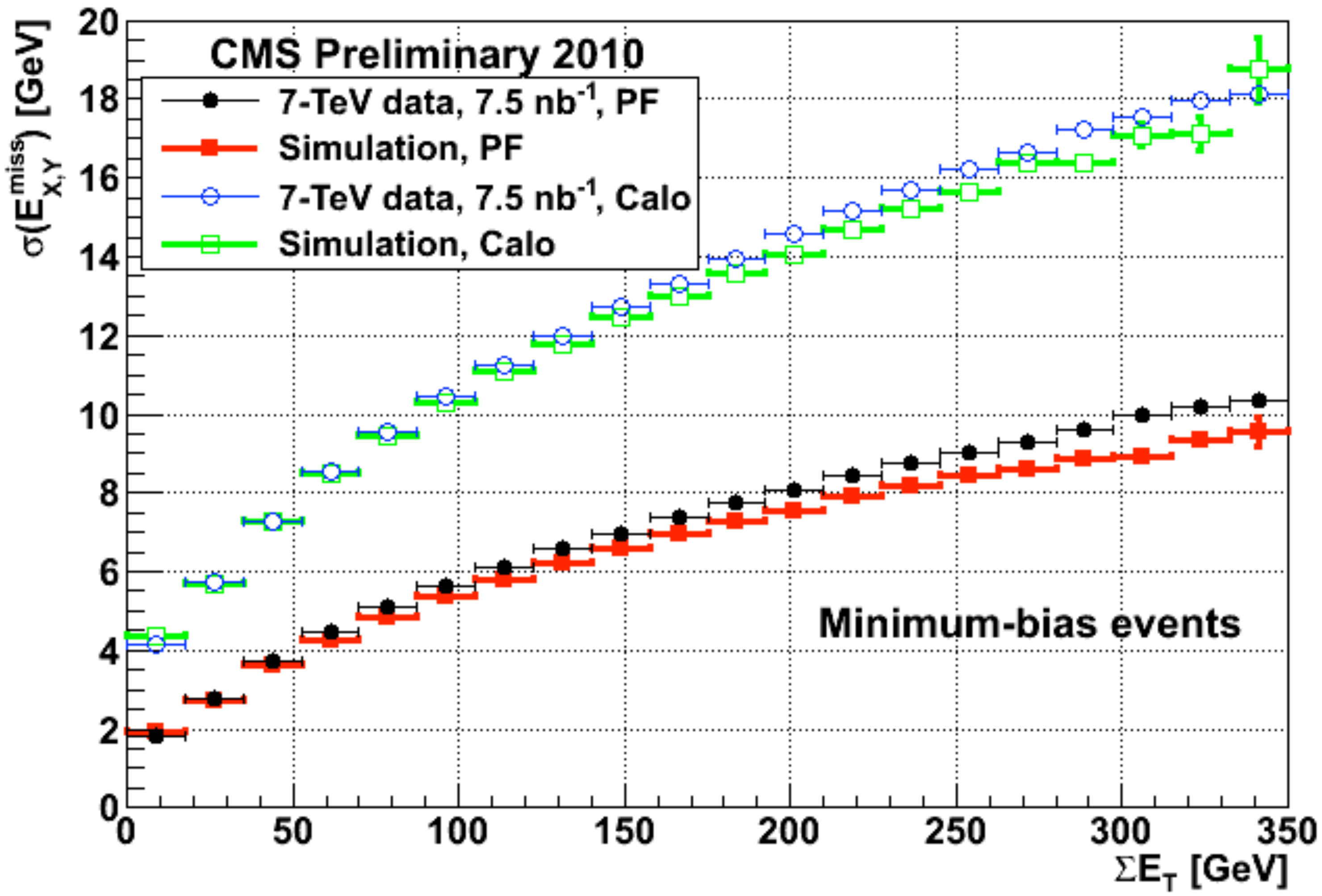


Jet PT = 65 GeV/c

Four true particles:
π^+, π^-, π^0, K_L^0




CMS Preliminary 2010
$\sqrt{s}$ = 7 TeV, MC
Anti-k_T R=0.5
p_T^{corr} > 25 GeV/c
neutral had
photons
charged
hadrons
Mean Fraction of Jet Energy
PFJet η
1.2
1
0.8
0.6
0.4
0.2
0
-4
-2
0
2
4



CMS Preliminary 2010
7-TeV data, 7.5 nb-1, PF
Simulation, PF
7-TeV data, 7.5 nb-1, Calo
Simulation, Calo
Minimum-bias events
σ(E_{X,Y}^{miss}) [GeV]
ΣE_T [GeV]
0
2
4
6
8
10
12
14
16
18
20
50
100
150
200
250
300
350

some references/links

- PDG reviews on particle interactions and particle detectors http://pdg.lbl.gov/
- C.Lippmann, hep-ex arXiv:1101.3276
- ATLAS, CMS, LHCB, ALICE performance papers
- R.Cavanaugh's lectures at HCP school 2012
- D.Bortoletto's lectures for CERN summer student
- W.Riegler's CERN academic training lectures, February 2014