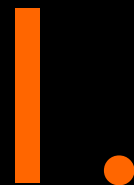
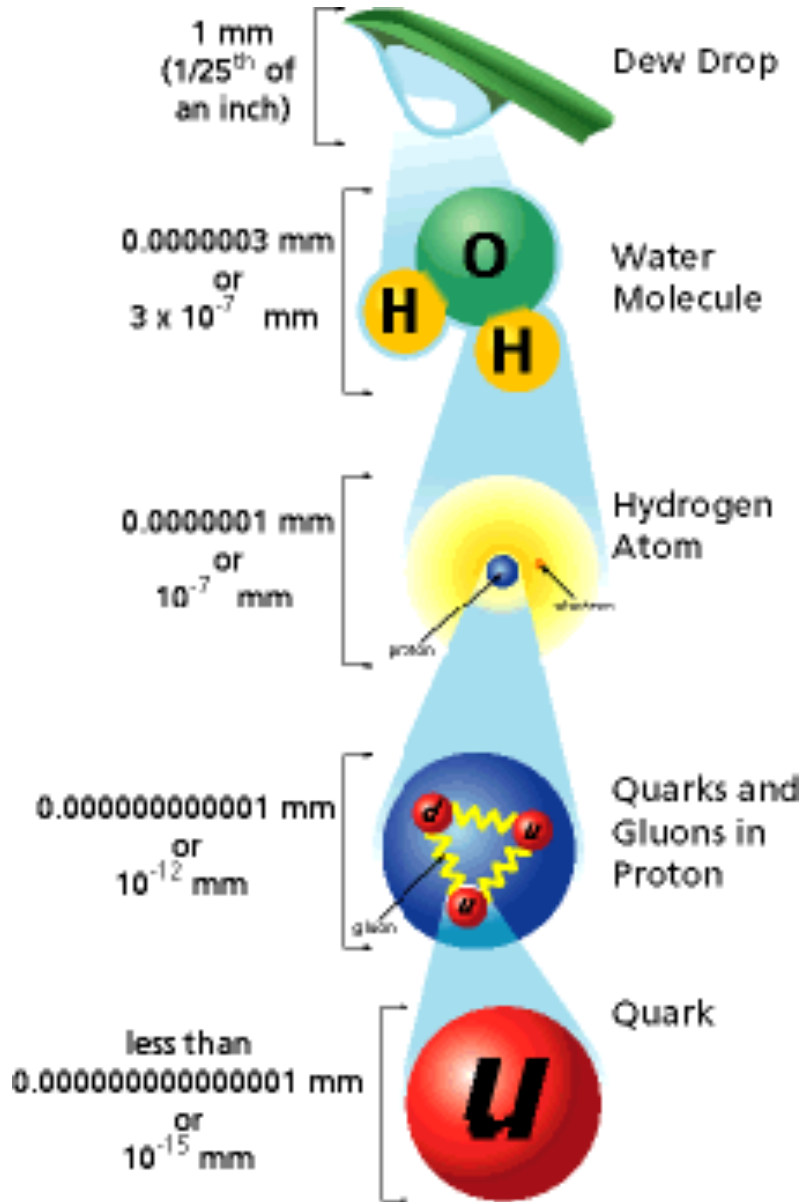


Experimental particle. physics



units, quantities,
kinematics, measurements

Order of magnitudes



Optical microscope resolution

$$\Delta r \sim \frac{1}{\sin \theta}$$

with θ = angular aperture of the light beam

De Broglie wave length

$$\lambda = \frac{h}{p} \qquad \Delta r \sim \frac{h}{p}$$

with p = transferred momentum

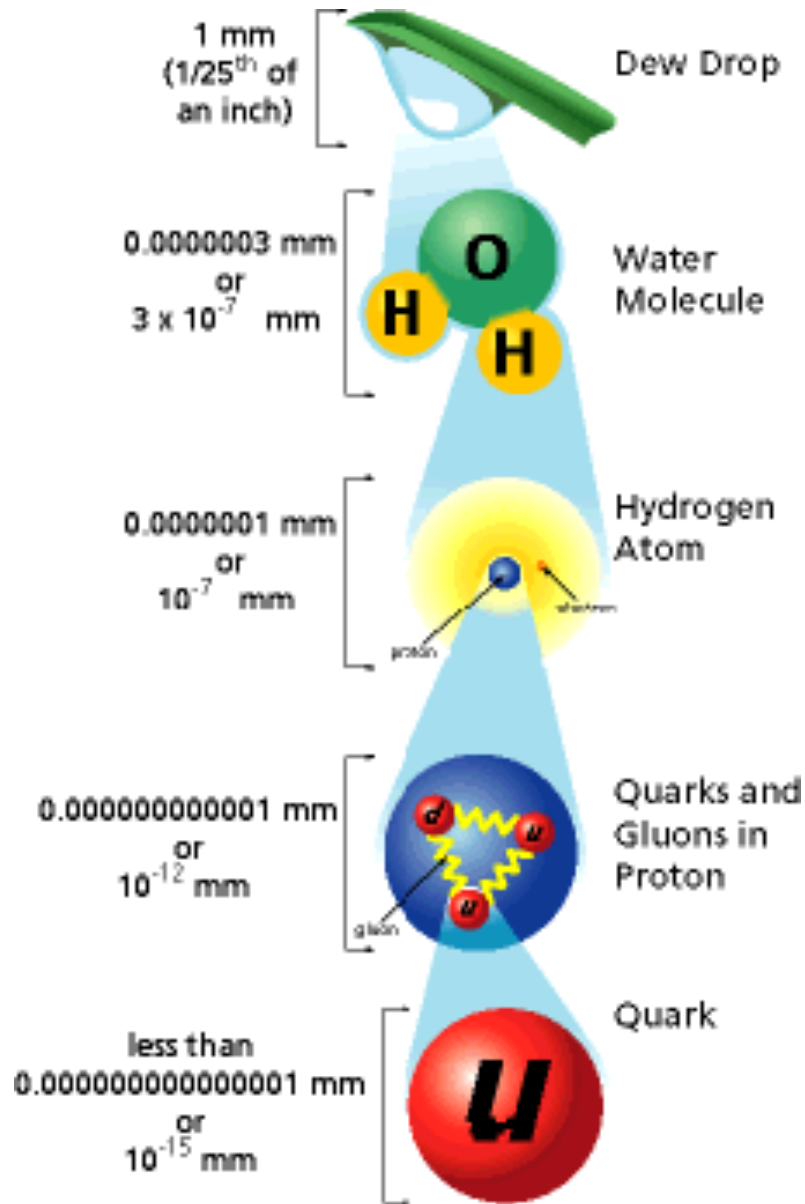
HEP, SI and “natural” units

Quantity	HEP units	SI units
length	1 fm	10^{-15} m
charge	e	$1.602 \cdot 10^{-19}$ C
energy	1 GeV	1.602×10^{-10} J
mass	1 GeV/c ²	1.78×10^{-27} kg
$\hbar = h/2\pi$	6.588×10^{-25} GeV s	1.055×10^{-34} Js
c	2.988×10^{23} fm/s	2.988×10^8 m/s
$\hbar c$	197 MeV fm	...

“natural” units ($\hbar = c = 1$)

mass	1 GeV
length	1 GeV ⁻¹ = 0.1973 fm
time	1 GeV ⁻¹ = 6.59×10^{-25} s

How much energy to probe these distances?



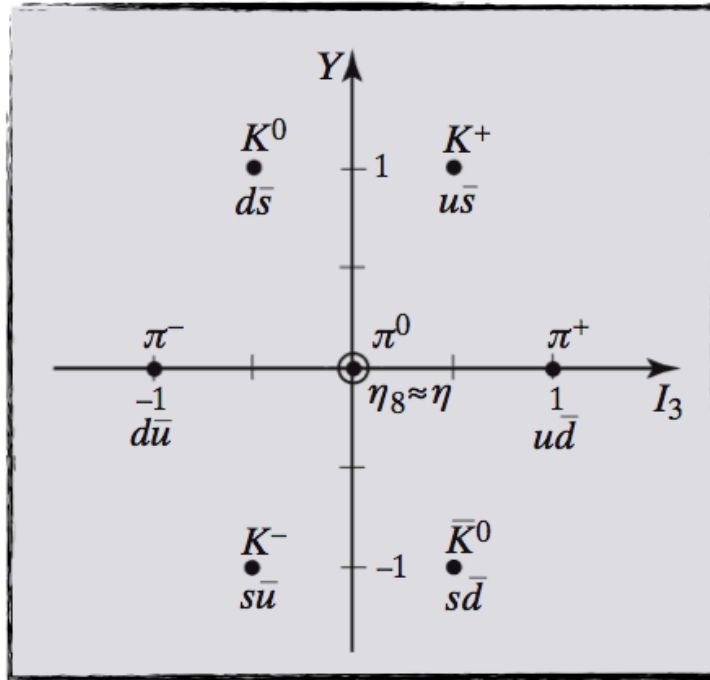
$$\lambda = \frac{h}{p} = \frac{2\pi\hbar c}{pc} = \frac{2\pi \times 197 \text{ MeV fm}}{pc}$$

What?	Dimension [m]	p [GeV/c]
Atom	10^{-10}	
Nucleus	10^{-14}	
Nucleon	10^{-15}	
Quark	10^{-18}	

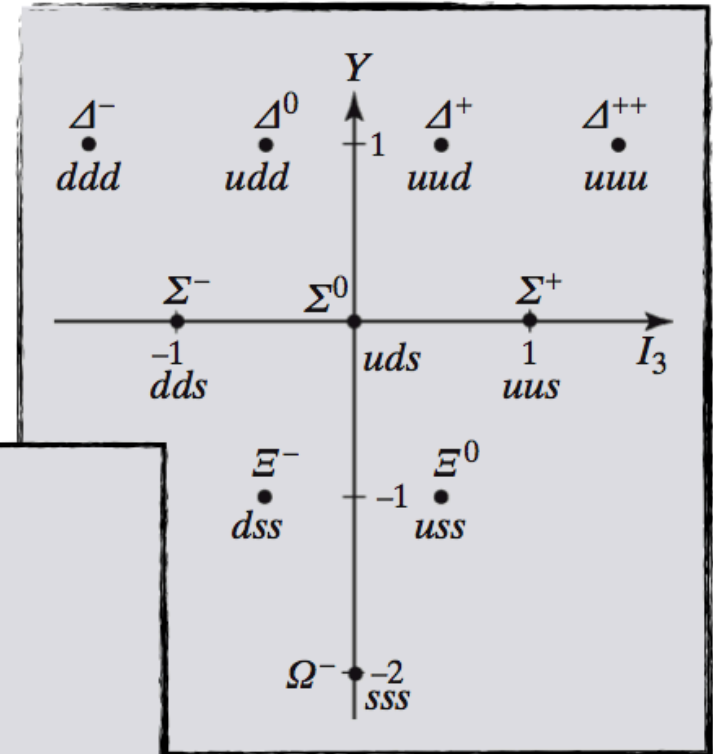
What do we want to measure?

1968: SLAC u up quark	1974: Brookhaven & SLAC c charm quark	1995: Fermilab t top quark	1979: DESY g gluon
1968: SLAC d down quark	1947: Manchester University s strange quark	1977: Fermilab b bottom quark	1923: Washington University* γ photon
1956: Savannah River Plant ν_e electron neutrino	1962: Brookhaven ν_μ muon neutrino	2000: Fermilab ν_τ tau neutrino	1983: CERN W W boson
1897: Cavendish Laboratory e electron	1937: Caltech and Harvard μ muon	1976: SLAC τ tau	1983: CERN Z Z boson
			2012: CERN H Higgs boson

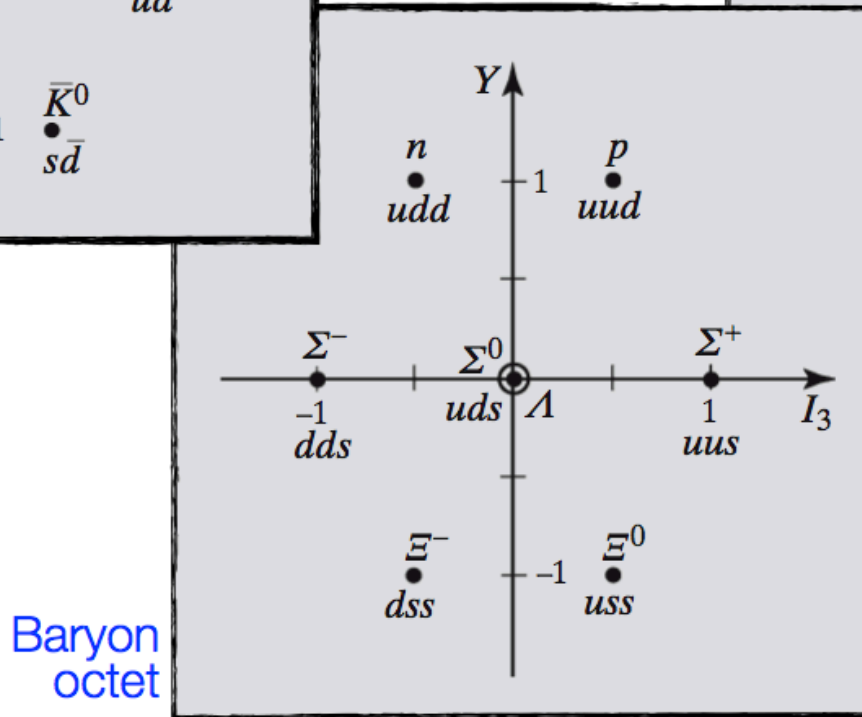
Baryons and Mesons



Meson octet



Baryon decuplet



Baryon octet

Measuring particles

- Particles are characterized by
 - ✓ **Mass** [Unit: eV/c² or eV]
 - ✓ **Charge** [Unit: e]
 - ✓ **Energy** [Unit: eV]
 - ✓ **Momentum** [Unit: eV/c or eV]
 - ✓ (+ spin, lifetime, ...)

Particle identification via measurement of:

e.g. (E, p, Q) or (p, β, Q)
(p, m, Q) ...

- ... and move at **relativistic speed**

$$\beta = \frac{v}{c} \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

$$\ell = \frac{\ell_0}{\gamma} \quad \text{length contraction}$$

$$t = t_0 \gamma \quad \text{time dilatation}$$

$$E^2 = \vec{p}^2 c^2 + m^2 c^4$$

$$E = m\gamma c^2 = mc^2 + E_{\text{kin}}$$

$$\vec{\beta} = \frac{\vec{p}c}{E} \quad \vec{p} = m\gamma \vec{\beta} c$$

Relativistic kinematics in a nutshell

$$E^2 = \vec{p}^2 + m^2$$

$$\ell = \frac{\ell_0}{\gamma}$$

$$E = m\gamma$$

$$t = t_0\gamma$$

$$\vec{p} = m\gamma\vec{\beta}$$

$$\vec{\beta} = \frac{\vec{p}}{E}$$

Center of mass energy

- In the **center of mass frame** the total momentum is 0
- In **laboratory frame** center of mass energy can be computed as:

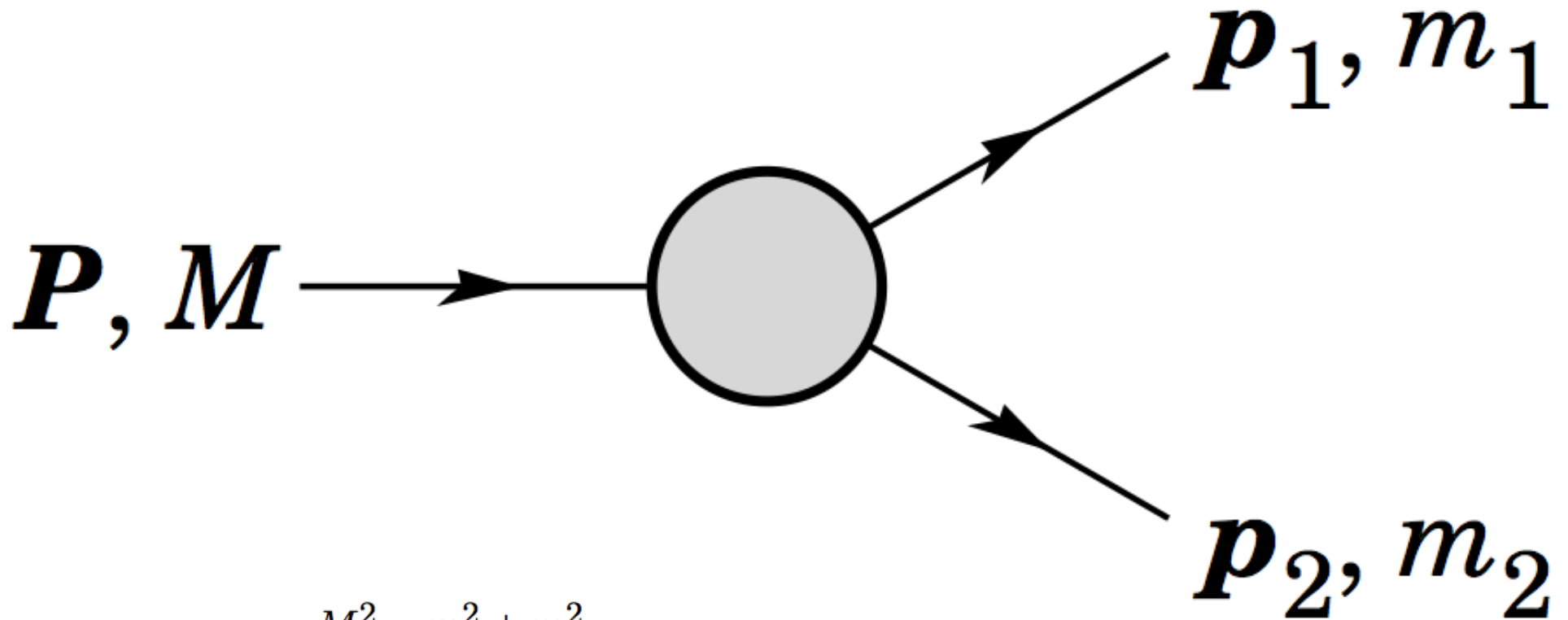
$$E_{\text{cm}} = \sqrt{s} = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$

Hint: it can be computed as the “length” of the total four-momentum, that is invariant:

$$p = (E, \vec{p}) \quad \sqrt{p \cdot p}$$

What is the “length” of a the four-momentum of a particle?

2-bodies decay



$$E_1 = \frac{M^2 - m_2^2 + m_1^2}{2M}, \text{ we'll compute this now...}$$

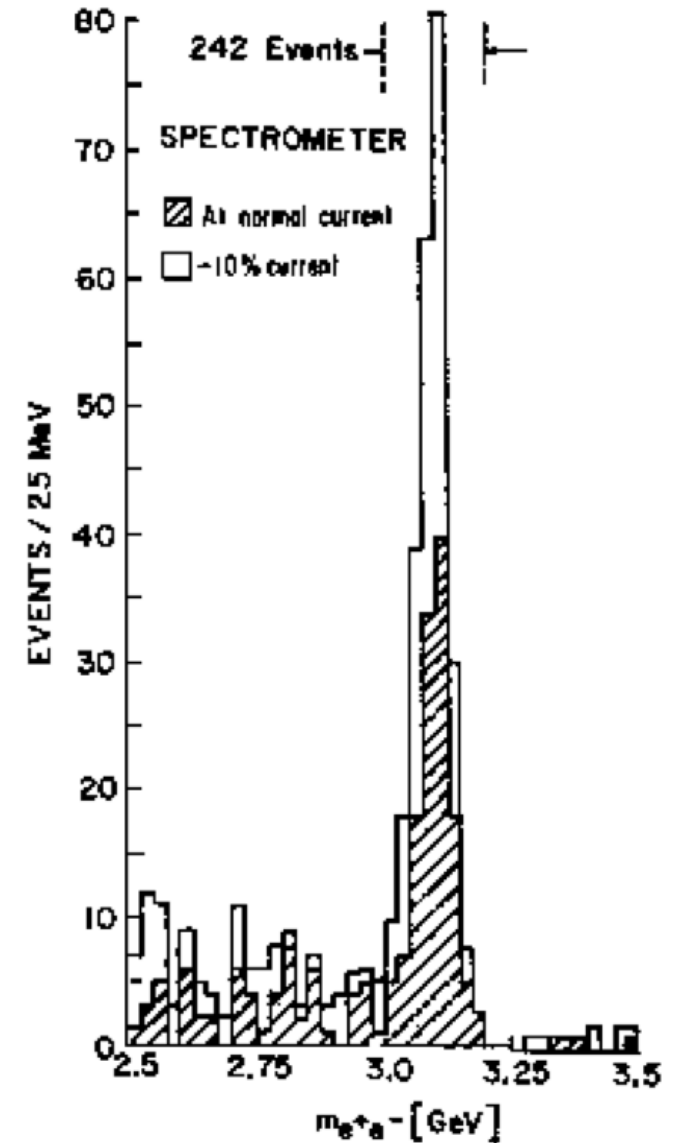
$$|\mathbf{p}_1| = |\mathbf{p}_2|$$

$$= \frac{[(M^2 - (m_1 + m_2)^2)(M^2 - (m_1 - m_2)^2)]^{1/2}}{2M}$$

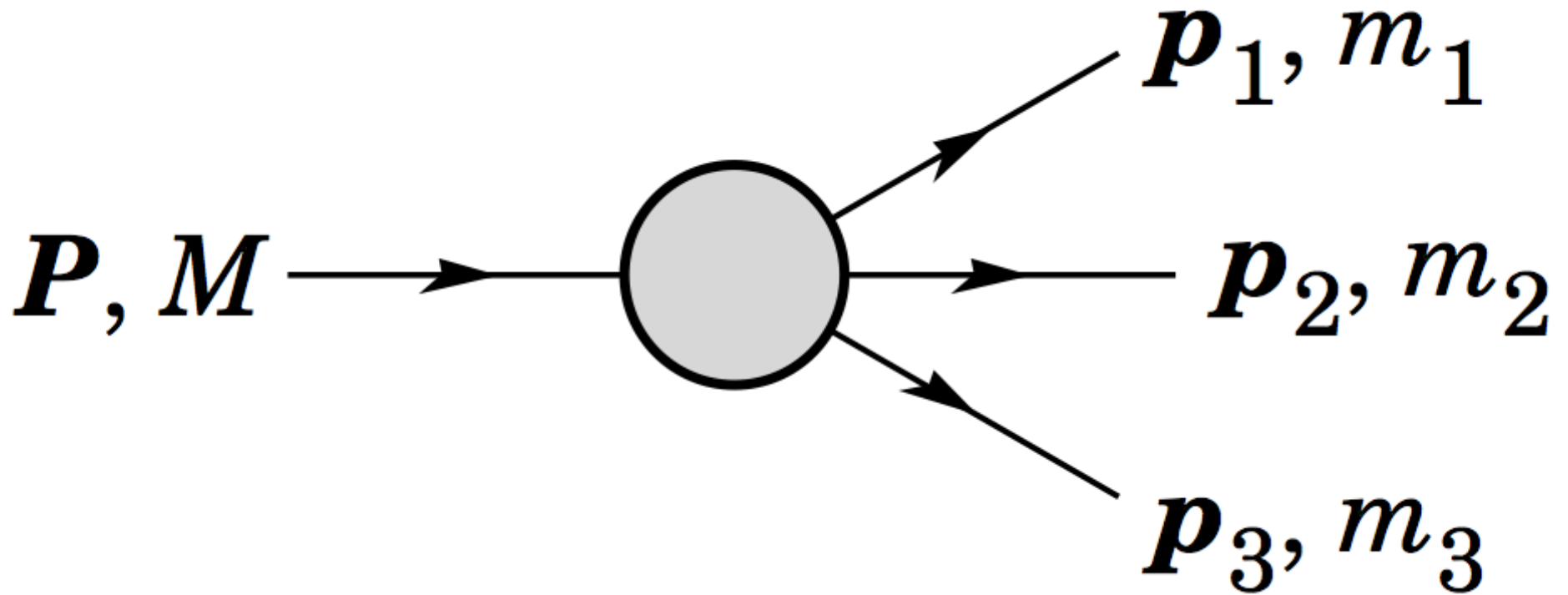
... you'll compute this as homework!

Invariant mass

$$M = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$



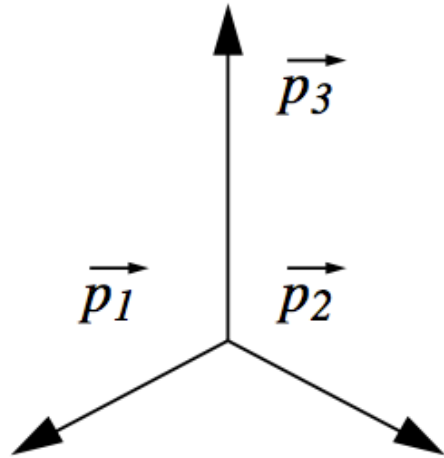
3-bodies decay



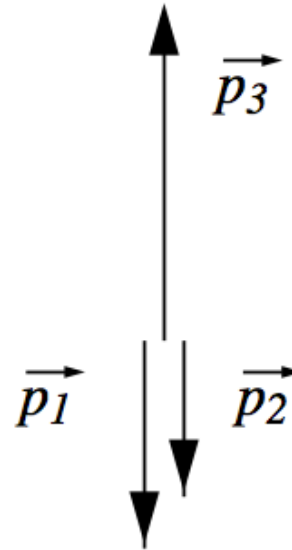
$$|\mathbf{p}_3| = \frac{[(M^2 - (m_{12} + m_3)^2)(M^2 - (m_{12} - m_3)^2)]^{1/2}}{2M}$$

3-bodies decay

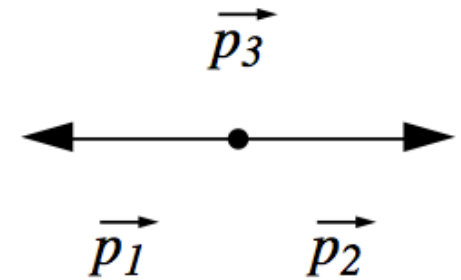
$$|\vec{p}_3| = \frac{[(M^2 - (m_{12} + m_3)^2)(M^2 - (m_{12} - m_3)^2)]^{1/2}}{2M}$$



(a)



(b)



(c)

$$\max(|\vec{p}_3|)$$

$$\min(|\vec{p}_3|)$$

$$(m_{12})_{min} = m_1 + m_2$$

$$(m_{12})_{max} = M - m_3$$

A real example: pion decay(s)

pion decays at rest

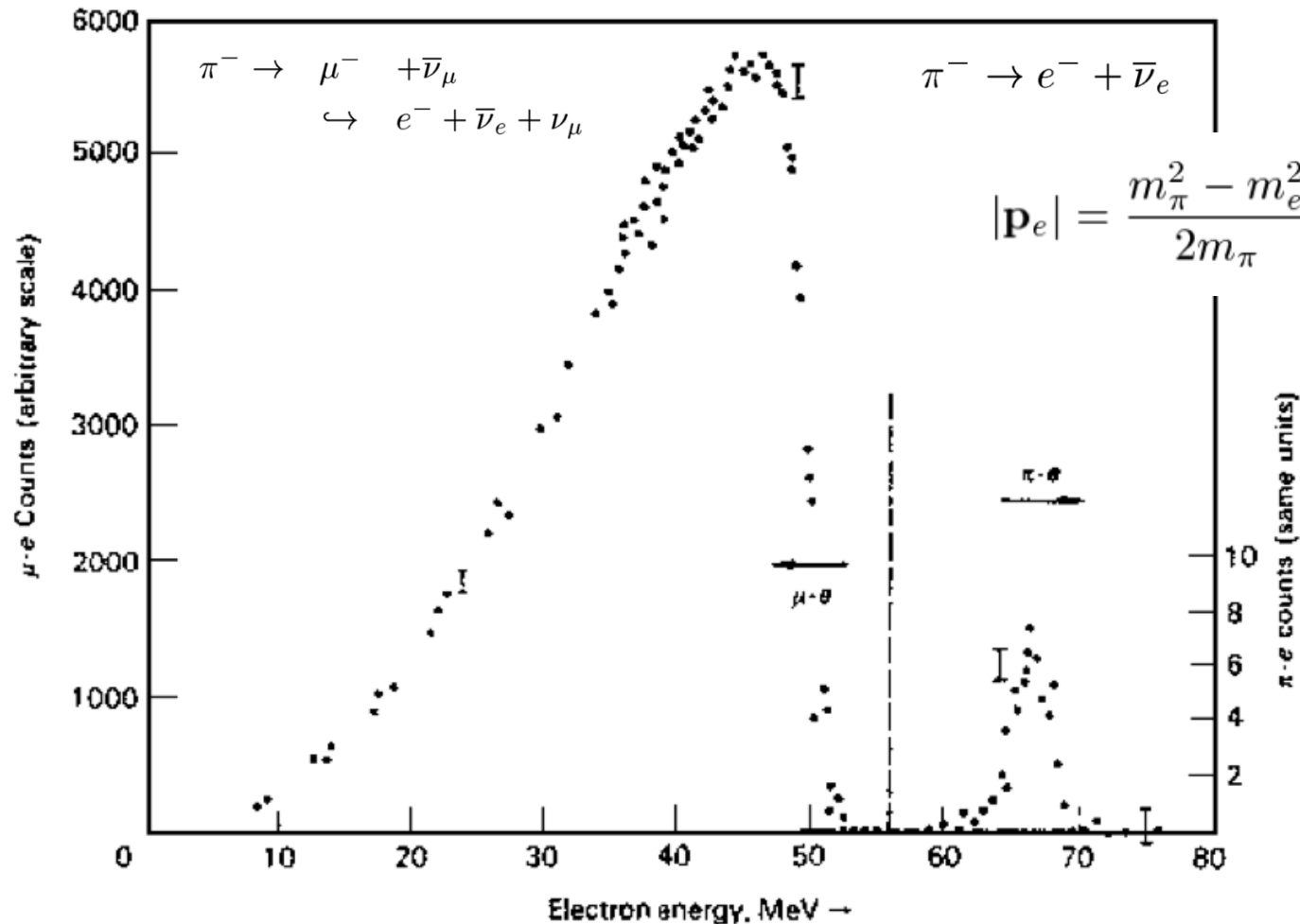
$$|\mathbf{p}_\mu| = \frac{m_\pi^2 - m_\mu^2}{2m_\pi} c \simeq 30 \text{ MeV}/c$$

$$m_\nu = 0$$

in most cases
muon decays
at rest

$$|\mathbf{p}_e|_{max} = \frac{m_\mu^2 - m_e^2}{2m_\mu} c \simeq 52 \text{ MeV}/c$$

$$|\mathbf{p}_e|_{min} = 0$$



3-bodies decay: Dalitz plot

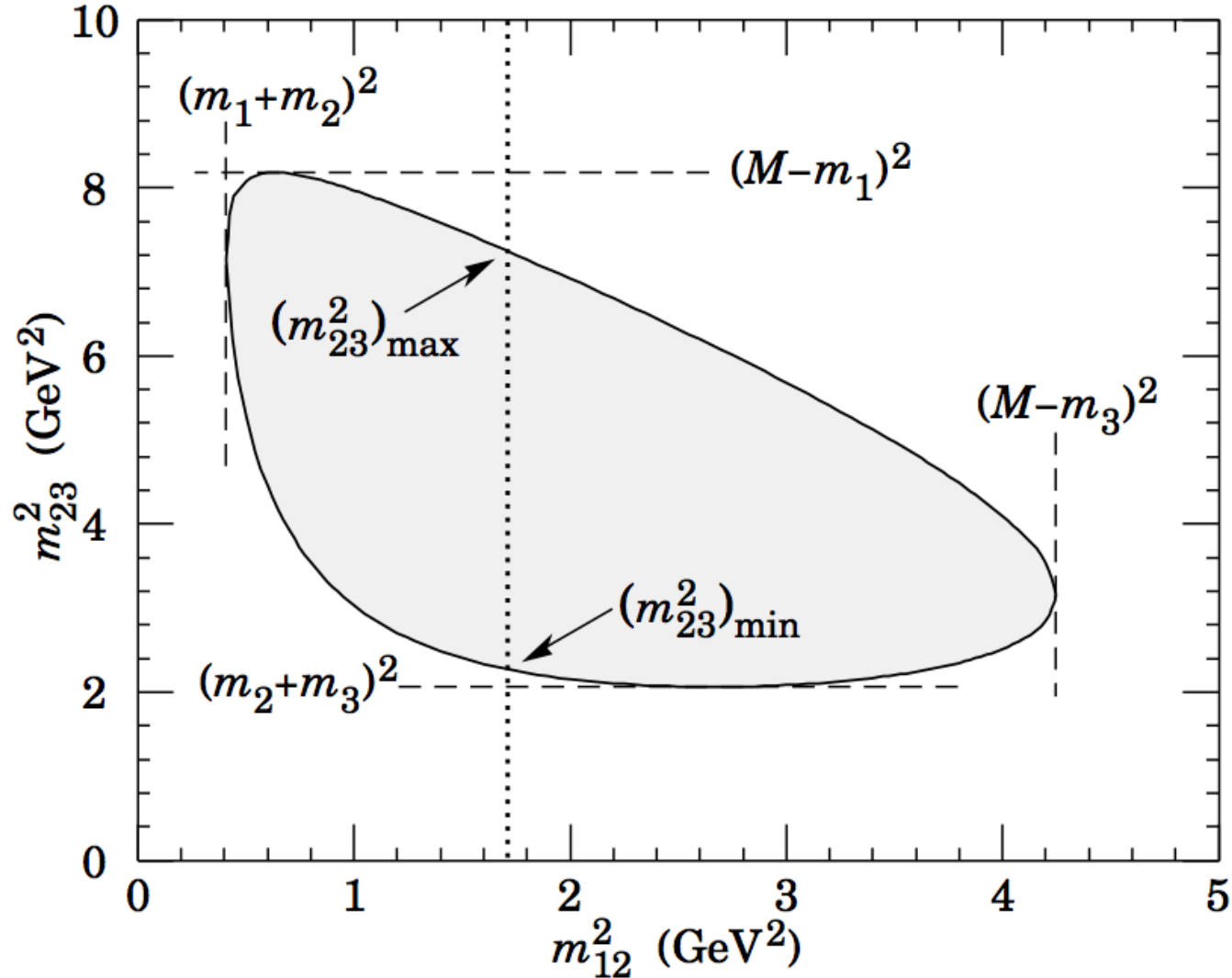
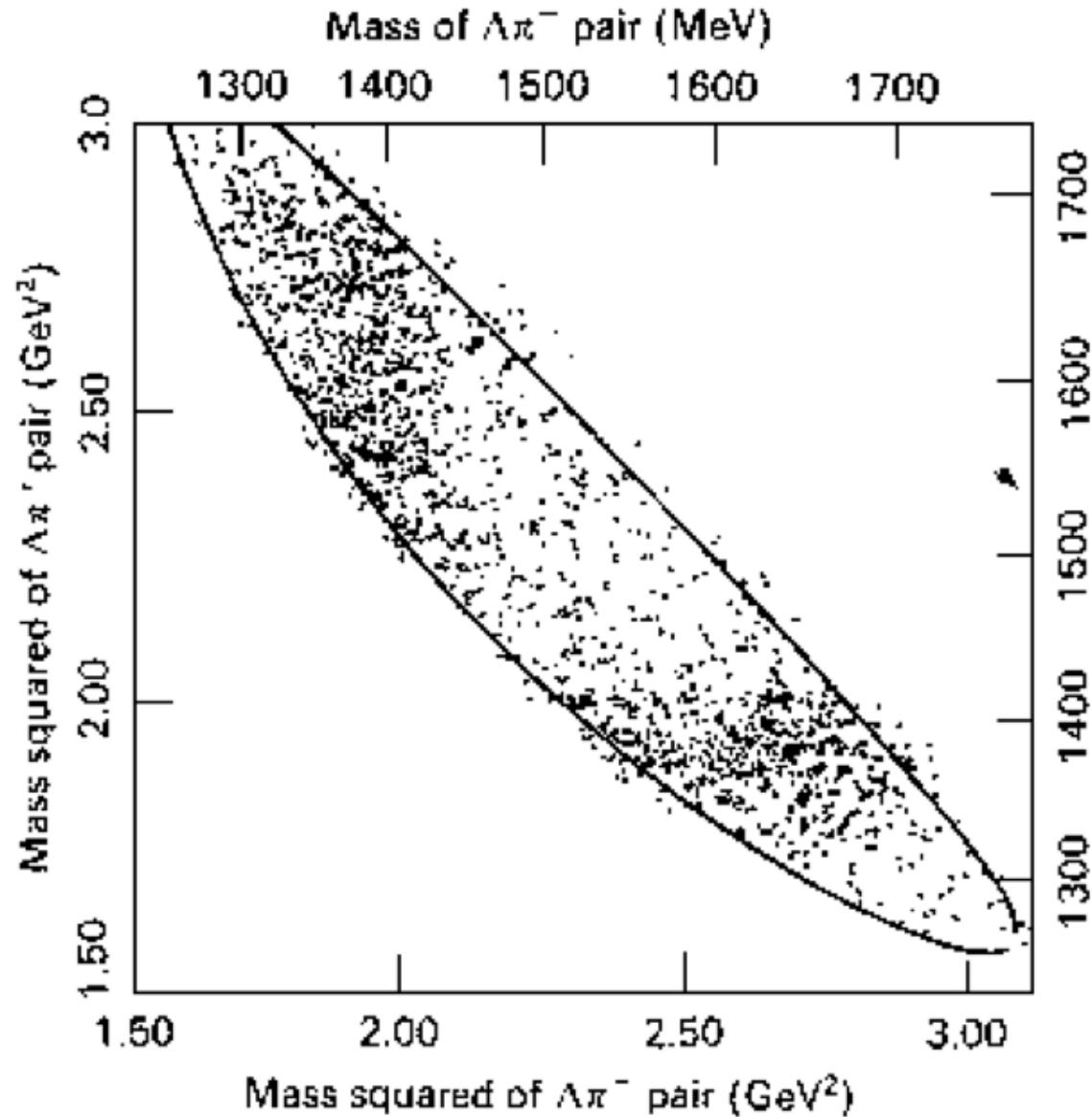


Figure 45.3: Dalitz plot for a three-body final state. In this example, the state is $\pi^+ \bar{K}^0 p$ at 3 GeV. Four-momentum conservation restricts events to the shaded region.

Multi-bodies decay

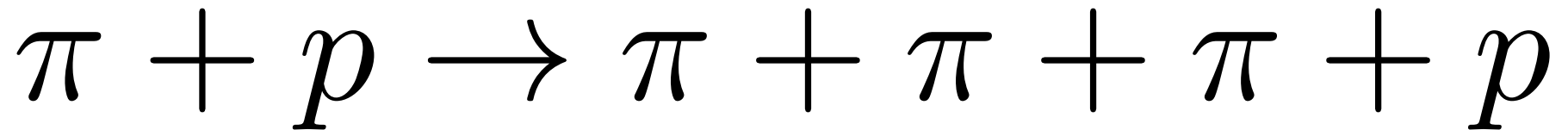
$$K^- + p \rightarrow \pi^+ + \pi^- + \Lambda$$



Reaction threshold

$$\sqrt{s} \geq \sum_i m_i c^2$$

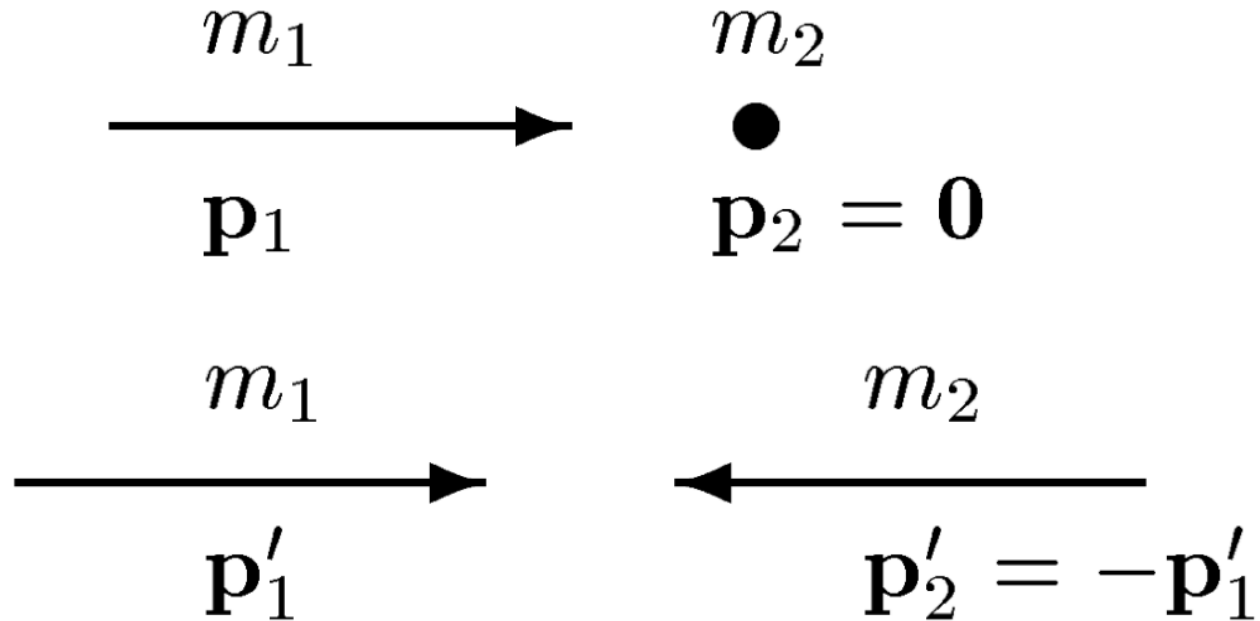
What energy should the pion have for this reaction to happen?



$$\begin{aligned} s &= (p_\pi + p_p)^2 c^2 = (E_\pi + m_p c^2)^2 - |\mathbf{p}_\pi|^2 \\ &= (m_\pi c^2)^2 + (m_p c^2)^2 + 2E_\pi(m_p c^2) \end{aligned}$$

$$E_\pi \geq \frac{(\sum_i m_i c^2)^2 - (m_\pi c^2)^2 - (m_p c^2)^2}{2m_p c^2} \simeq 500 \text{ MeV}$$

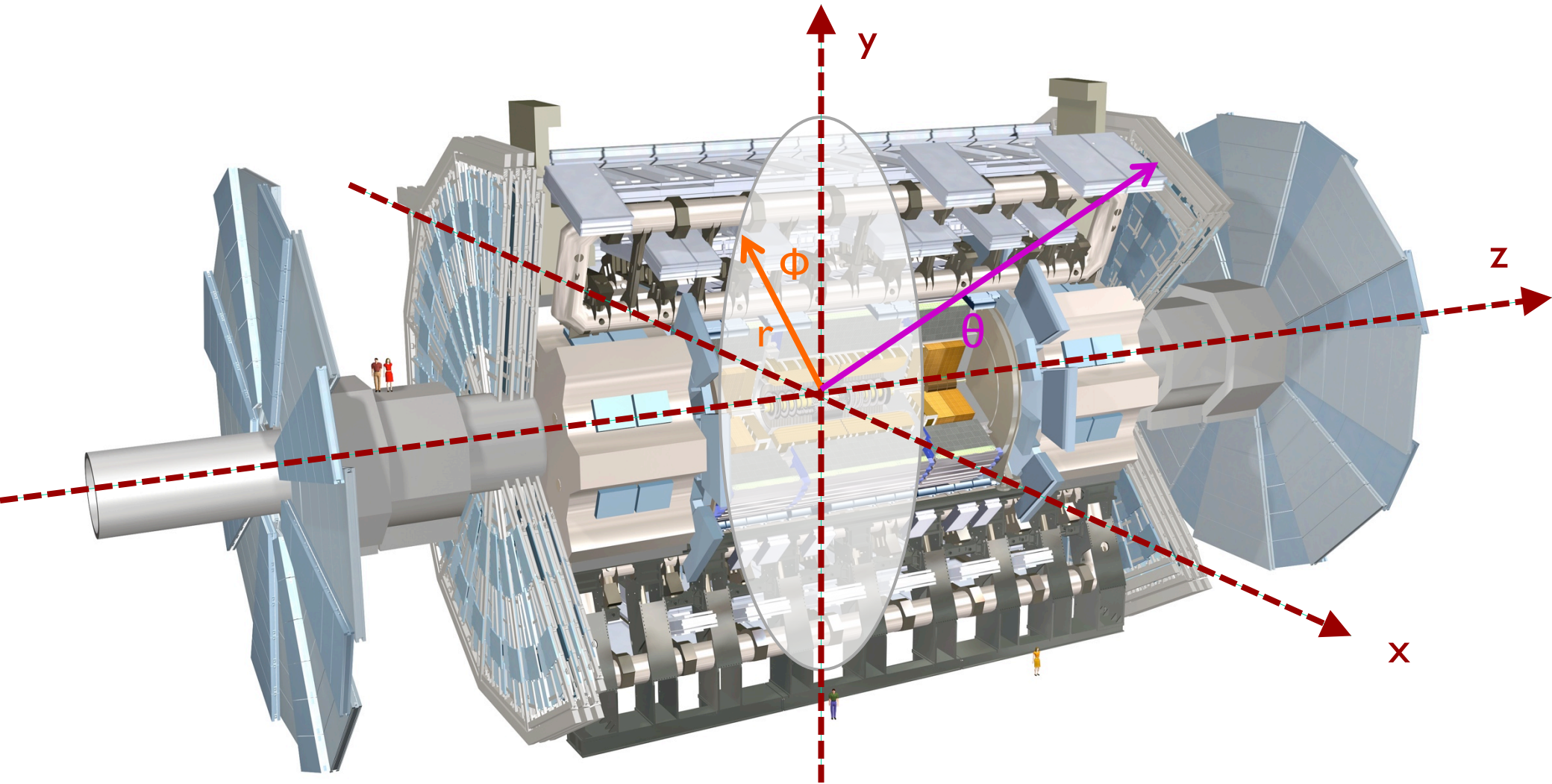
Fixed target vs. collider



How much energy should a fixed target experiment have to equal the center of mass energy of two colliding beam?

$$E_{\text{fix}} = 2 \frac{E_{\text{col}}^2}{m} - m$$

Collider experiment coordinates



Rapidity

Lorentz factor $\gamma = \frac{1}{\sqrt{1 - \beta^2}} = \cosh \varphi$ Hyperbolic cosine of “rapidity”

$$\begin{aligned} E &= m \cosh \varphi \\ |\vec{p}| &= m \sinh \varphi \end{aligned} \quad \varphi = \tanh^{-1} \frac{E}{|\vec{p}|} = \frac{1}{2} \ln \frac{E + |\vec{p}|}{E - |\vec{p}|}$$

- Particle physicists prefer to use modified rapidity along beam axis

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

Pseudorapidity

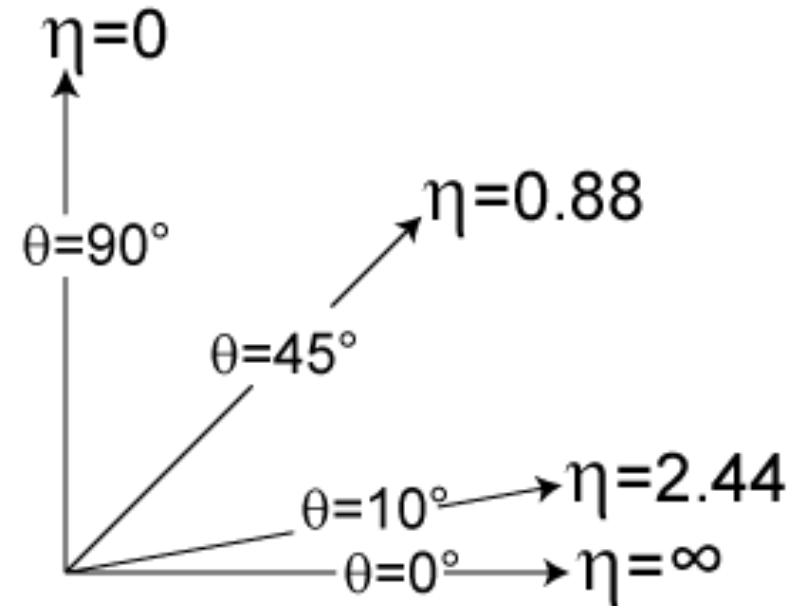
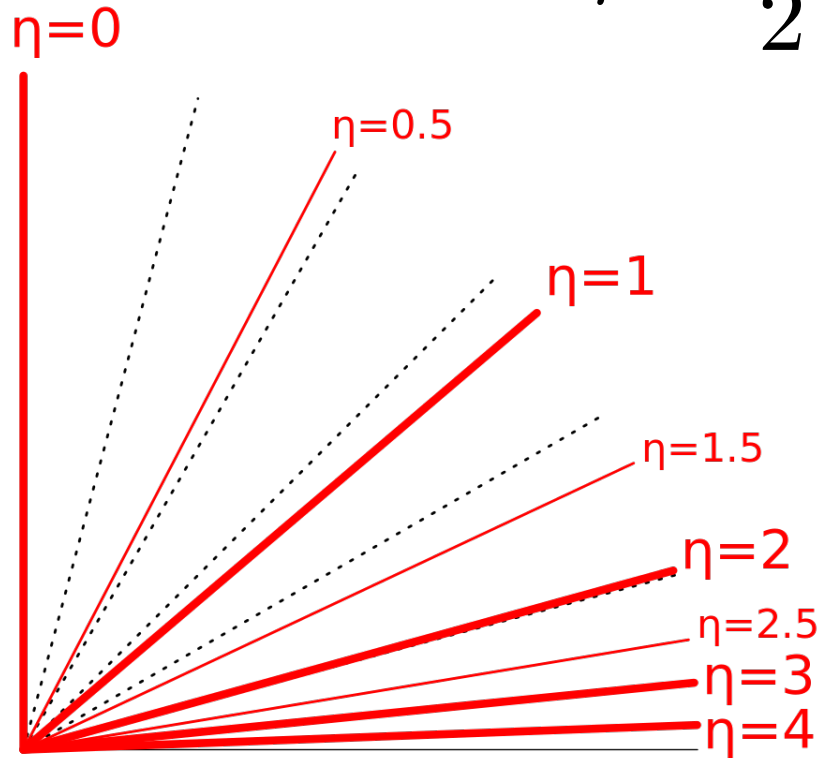
$$\eta = \frac{1}{2} \ln \frac{|\vec{p}| + p_z}{|\vec{p}| - p_z}$$

$$\eta = \frac{1}{2} \ln \left(\tan \frac{\theta}{2} \right)$$

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

$$\eta \simeq y$$

if $E \gg m$



Transverse variables

- At hadron colliders, a significant and unknown fraction of the beam energy in each event escapes down the beam pipe.
- Net momentum can only be constrained in the plane transverse to the beam z-axis!

$$\sum p_T(i) = 0$$

$$p_T = \sqrt{p_x^2 + p_y^2}$$
$$p_x = p_T \cos \phi$$
$$p_y = p_T \sin \phi$$
$$p_z = p_T \sinh \eta$$

$$|p| = p_T \cosh \eta$$

$$E_T = \frac{E}{\cosh \eta}$$

Missing transverse energy and transverse mass

- If invisible particles are created, only their transverse momentum can be constrained: **missing transverse energy**

$$E_T^{\text{miss}} = \sum p_T(i)$$

- If a heavy particle is produced and decays into two particles one of which is invisible, the mass of the parent particle can be constrained with the **transverse mass quantity**

$$\begin{aligned} M_T^2 &\equiv [E_T(1) + E_T(2)]^2 - [\mathbf{p}_T(1) + \mathbf{p}_T(2)]^2 \\ &= m_1^2 + m_2^2 + 2[E_T(1)E_T(2) - \mathbf{p}_T(1) \cdot \mathbf{p}_T(2)] \end{aligned}$$

$$\text{if } m_1 = m_2 = 0 \quad M_T^2 = 2|\mathbf{p}_T(1)||\mathbf{p}_T(2)|(1 - \cos \phi_{12})$$

$W \rightarrow e \nu$ discovery

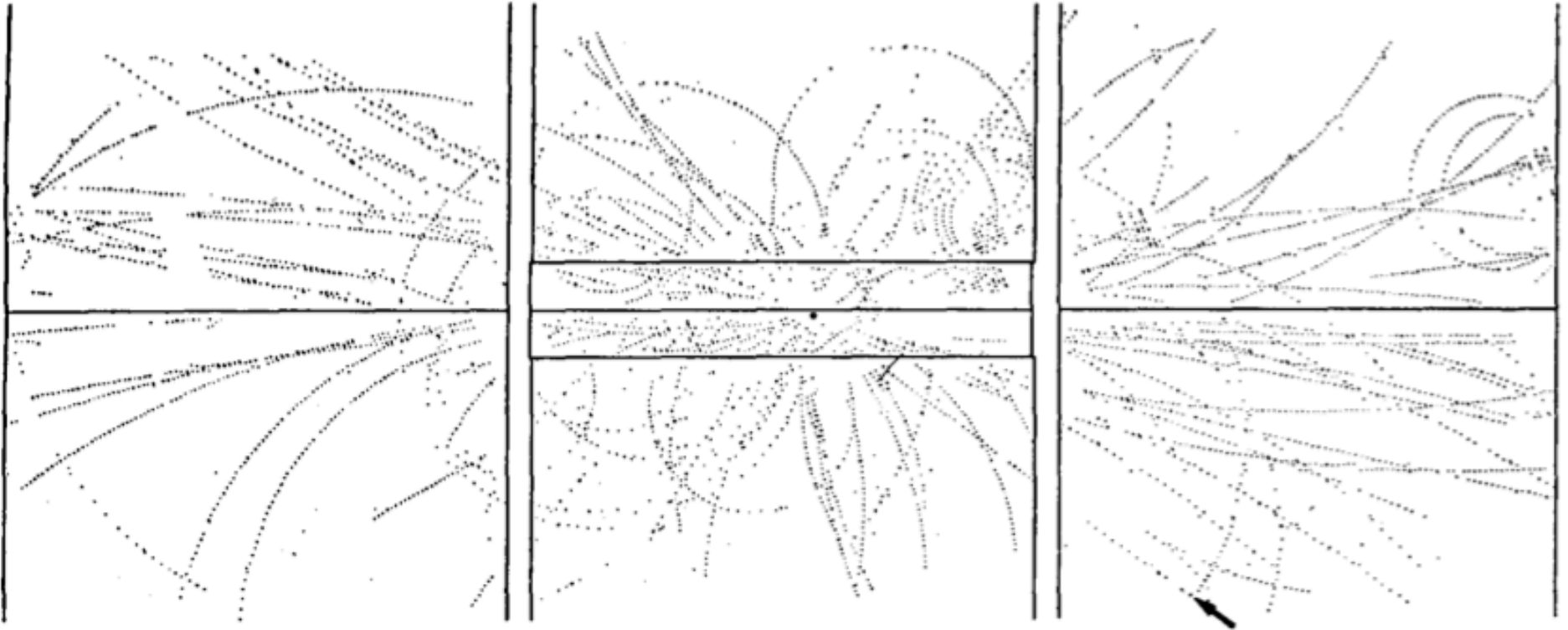
Volume 122B, number 1

PHYSICS LETTERS

24 February 1983

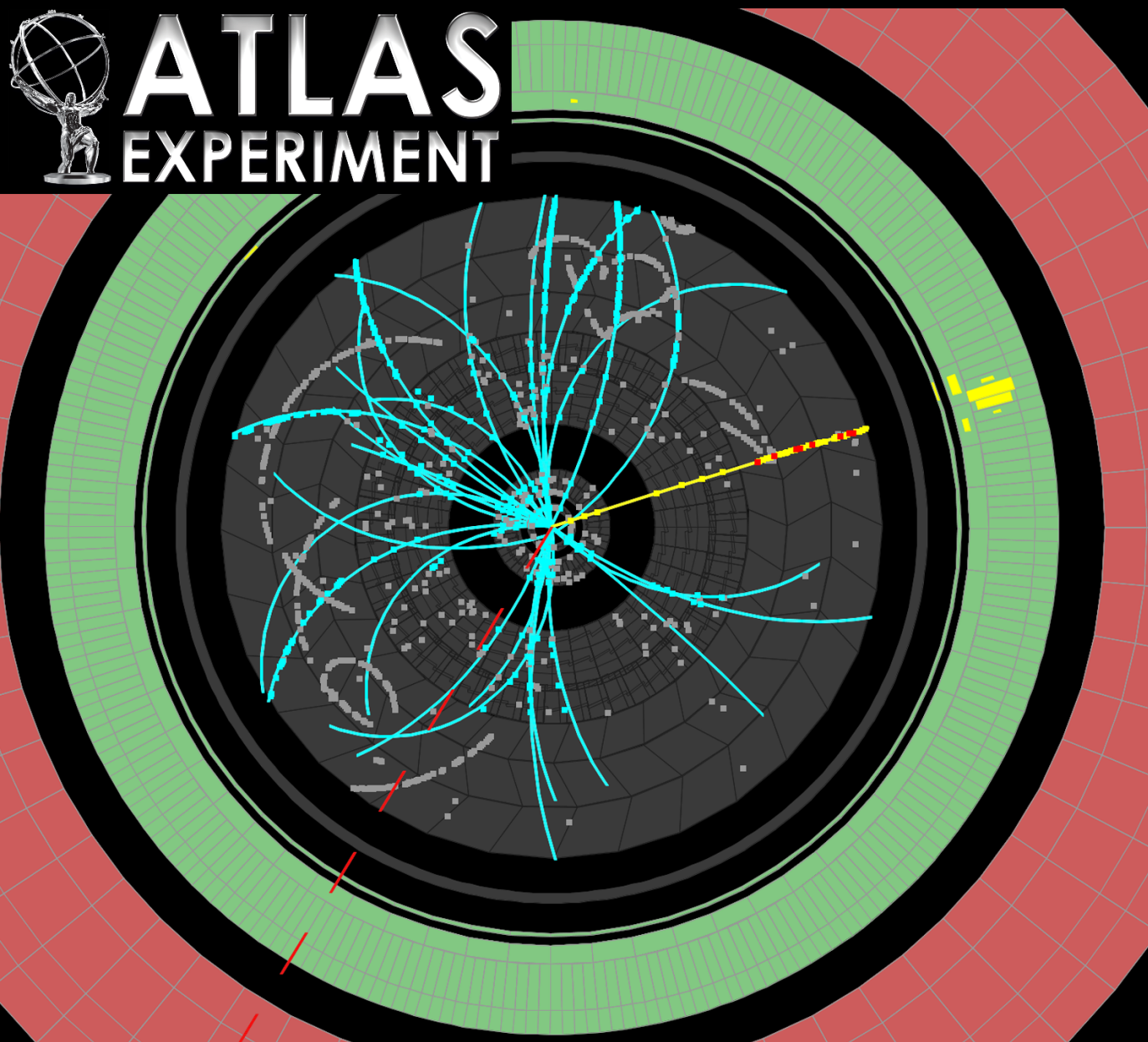
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EVENT 2958. 1279.



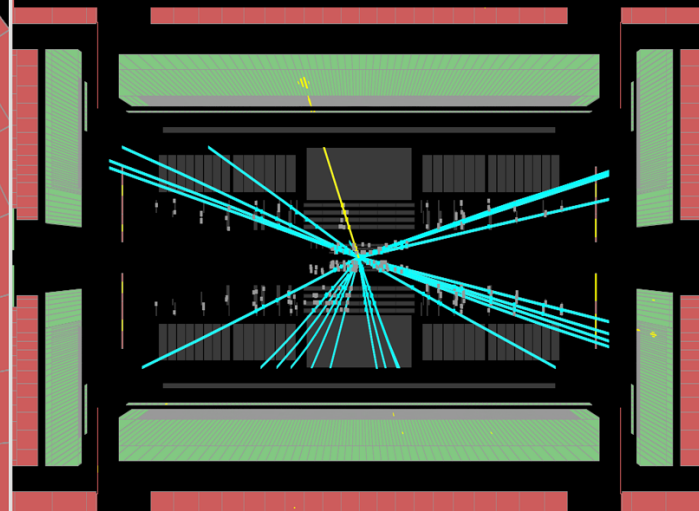


ATLAS EXPERIMENT



Run Number: 152409, Event Number: 5966801

Date: 2010-04-05 06:54:50 CEST



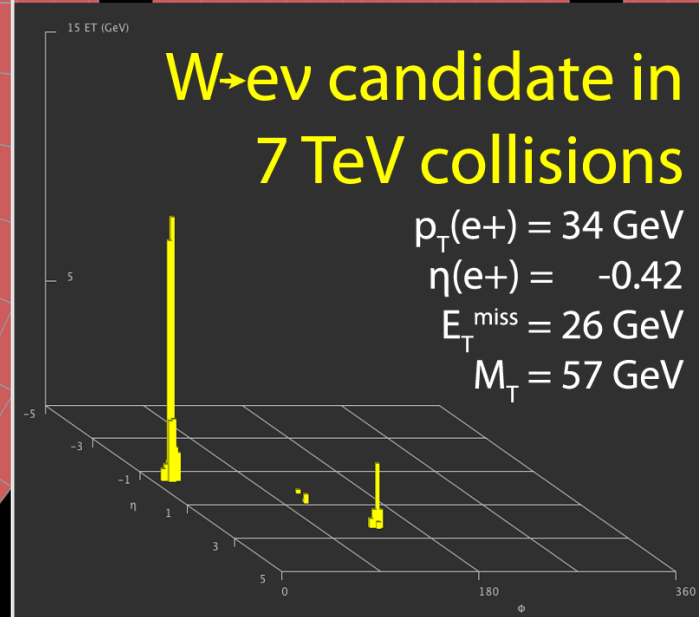
W→ev candidate in 7 TeV collisions

$p_T(e^+) = 34 \text{ GeV}$

$\eta(e^+) = -0.42$

$E_T^{\text{miss}} = 26 \text{ GeV}$

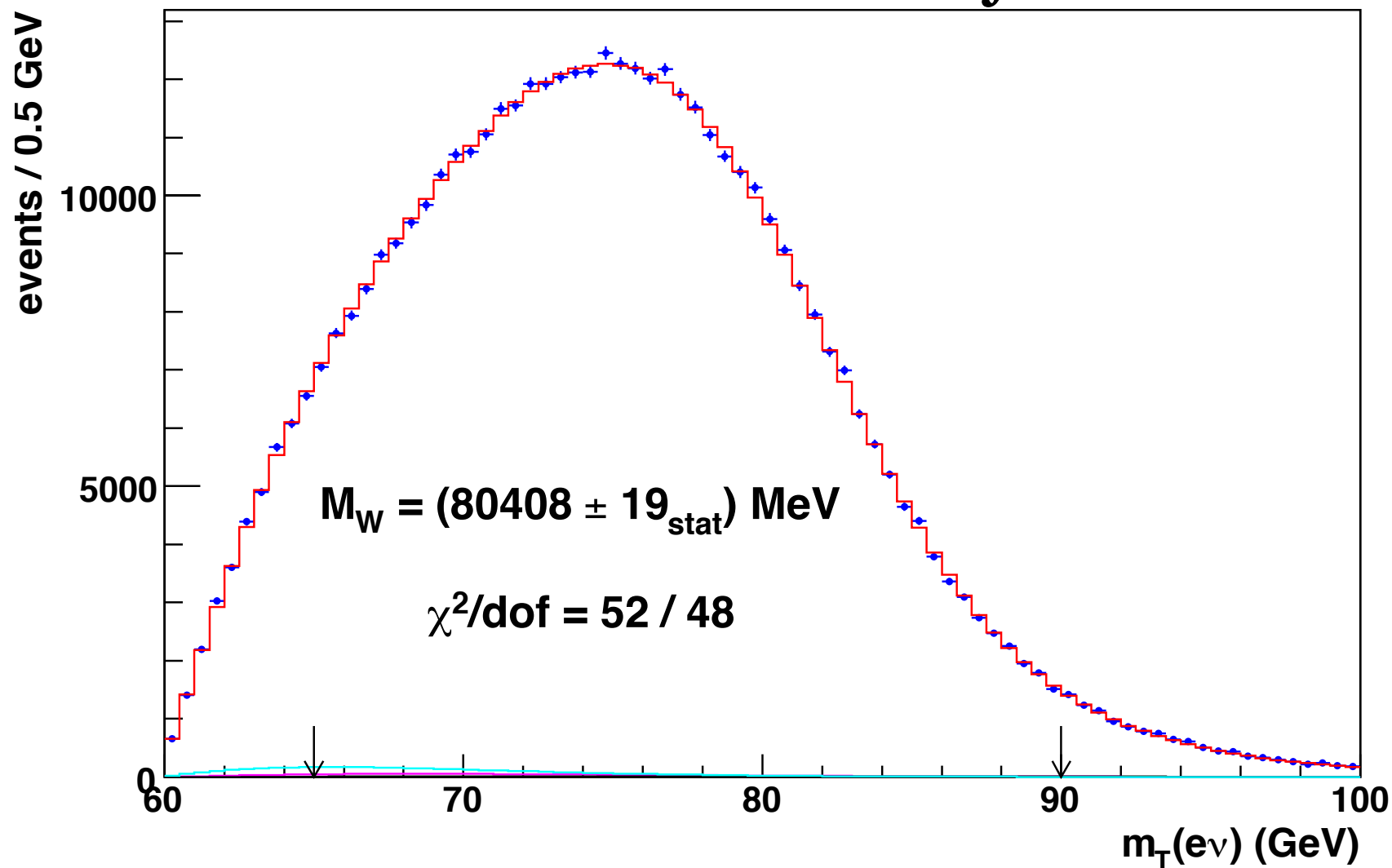
$M_T = 57 \text{ GeV}$



$W \rightarrow e \nu$

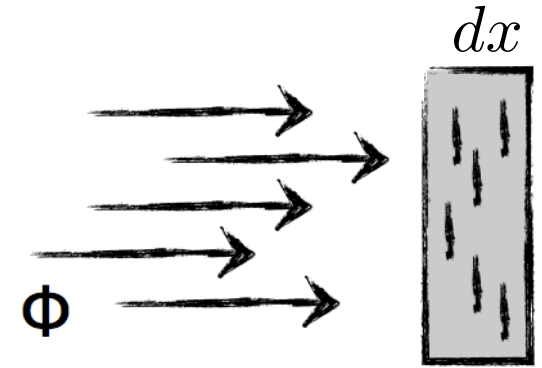
CDF II preliminary

$$\int L dt \approx 2.2 \text{ fb}^{-1}$$



Interaction cross section

Flux $\Phi = \frac{1}{S} \frac{dN_i}{dt}$ $[L^{-2} t^{-1}]$



Reactions per unit of time $\frac{dN_{\text{reac}}}{dt} = \Phi \overbrace{\sigma N_{\text{target}} dx}^{\text{area obscured by target particle}}$ $[t^{-1}]$

$[L^{-2} t^{-1}]$ $[?]$ $[L^{-1}]$ $[L]$

Reaction rate per target particle $W_{if} = \Phi \sigma$ $[t^{-1}]$

Cross section per target particle $\sigma = \frac{W_{if}}{\Phi}$ $[L^2]$ = reaction rate per unit of flux

$1b = 10^{-28} \text{ m}^2$ (roughly the area of a nucleus with $A = 100$)

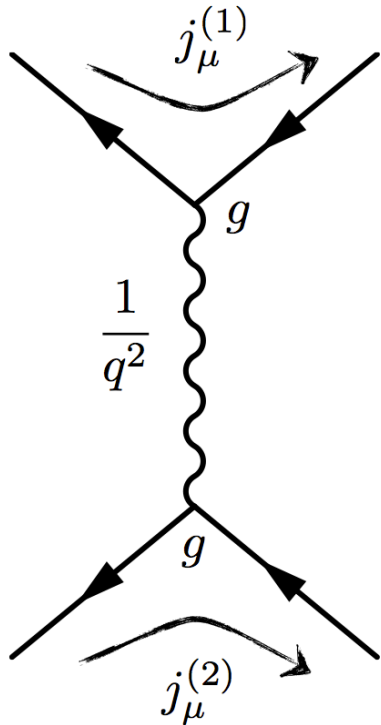
Fermi Golden Rule

From non-relativistic perturbation theory...

transition probability matrix element energy density of final states

$$W_{if} = \frac{2\pi}{\hbar} |M_{if}|^2 \frac{dN}{dE_f}$$

[τ^{-1}] [E] [E^{-1}]



$$M_{if} = -i \int j_\mu^{(1)} \left(\frac{1}{q^2} \right) j_\mu^{(2)} d^4x$$

$$\sigma \sim |M_{if}|^2 \sim g^4 \left(\frac{1}{q^4} \right)$$

Cross section: magnitude and units

Standard

cross section unit:

$$[\sigma] = \text{mb}$$

with $1 \text{ mb} = 10^{-27} \text{ cm}^2$

or in

natural units:

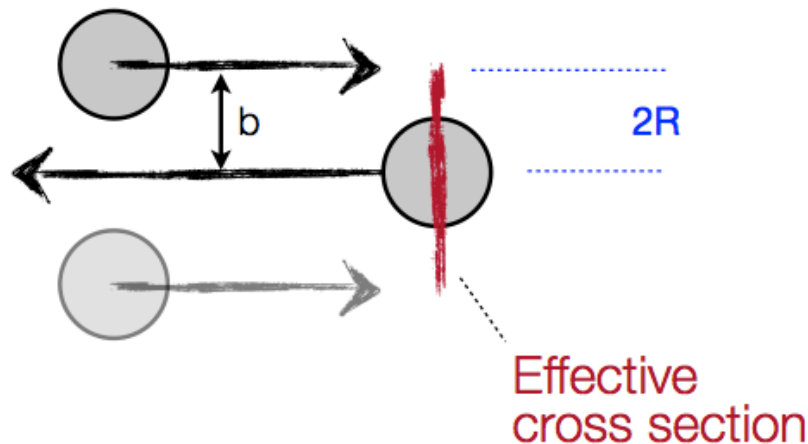
$$[\sigma] = \text{GeV}^{-2}$$

with $1 \text{ GeV}^{-2} = 0.389 \text{ mb}$

$$1 \text{ mb} = 2.57 \text{ GeV}^{-2}$$

Estimating the
proton-proton cross section:

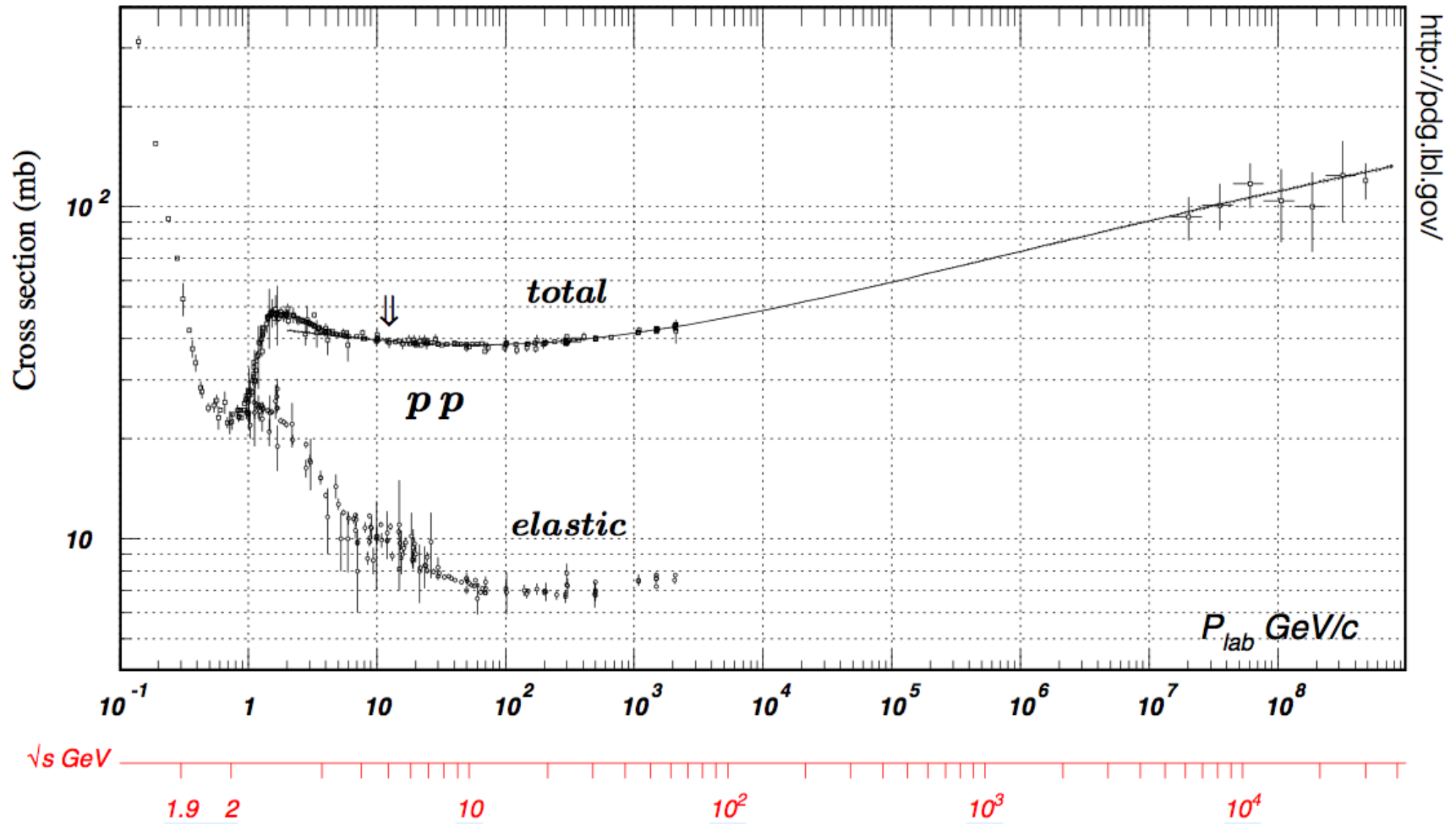
using: $\hbar c = 0.1973 \text{ GeV fm}$
 $(\hbar c)^2 = 0.389 \text{ GeV}^2 \text{ mb}$



Proton radius: $R = 0.8 \text{ fm}$

Strong interactions happens up to $b = 2R$

Proton-proton scattering cross-section



Cross-sections at LHC

