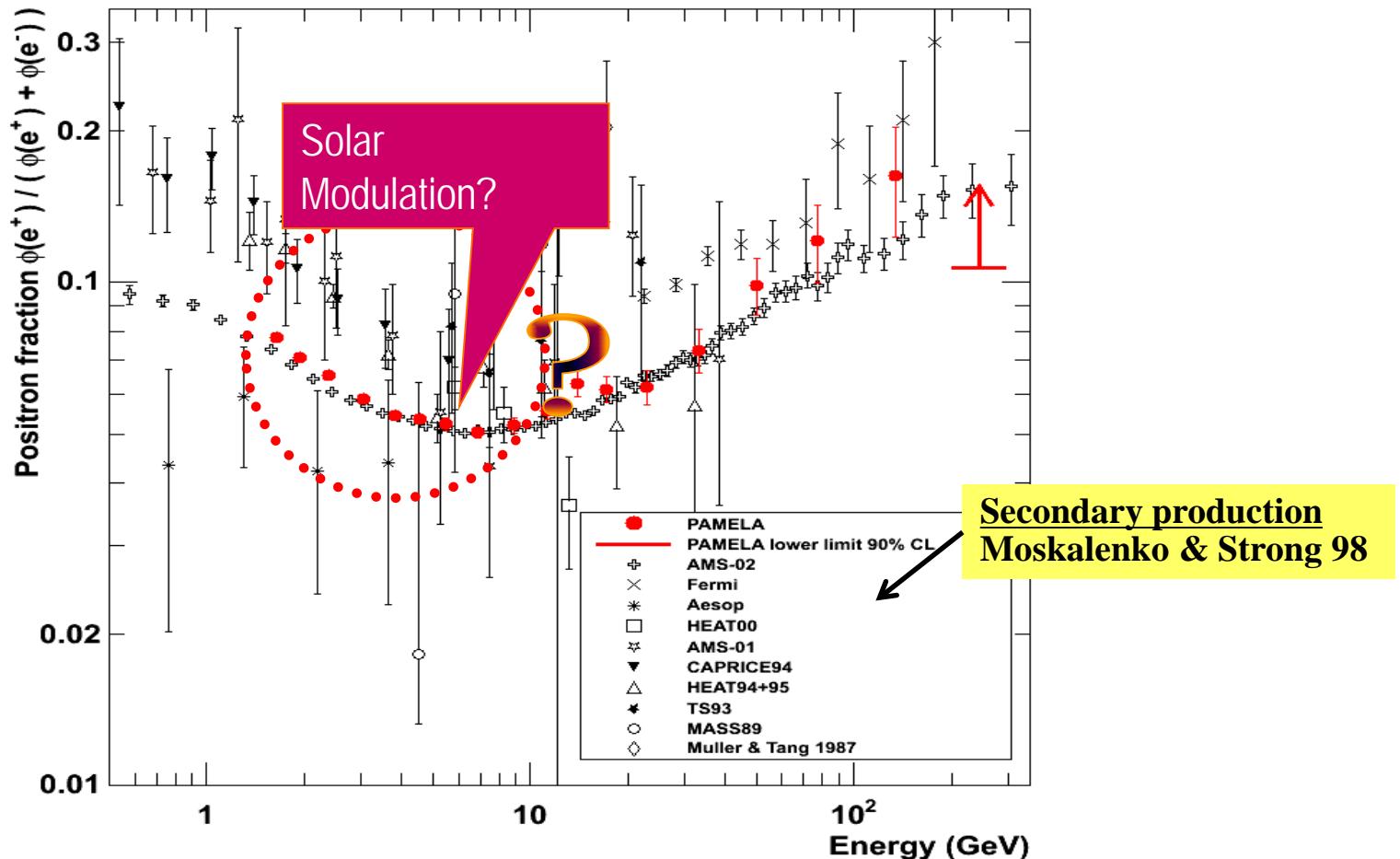
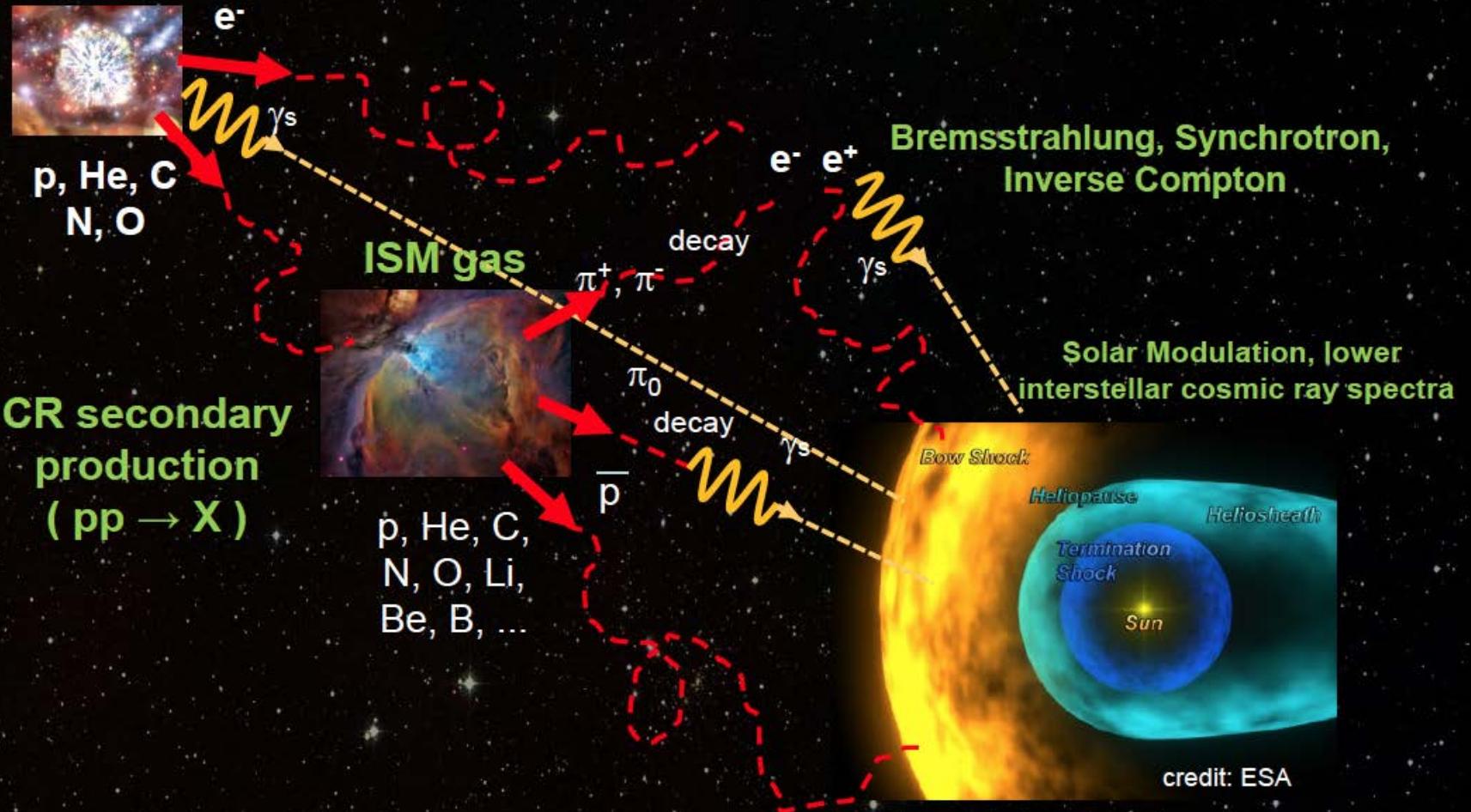


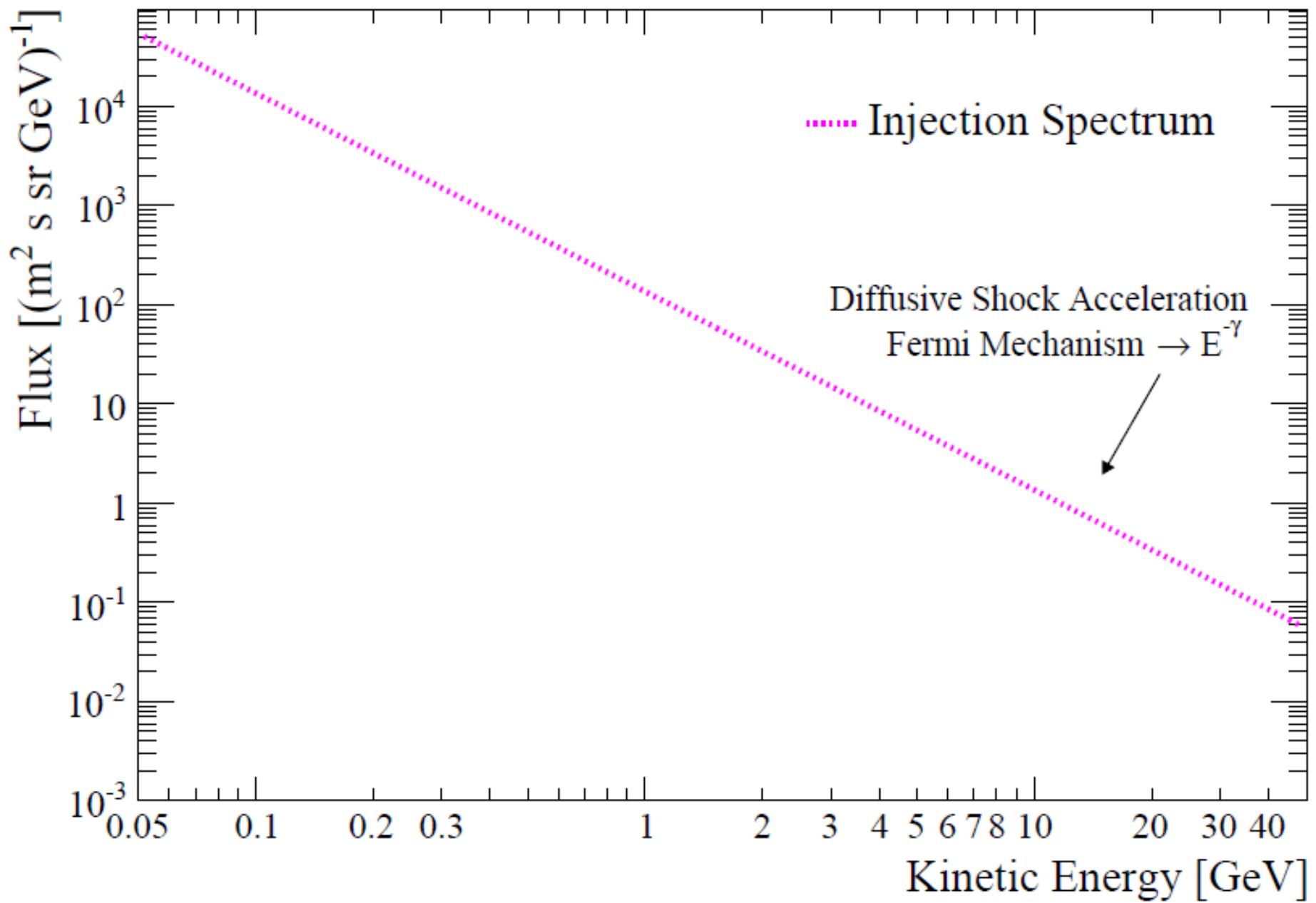
Positron to Electron Fraction

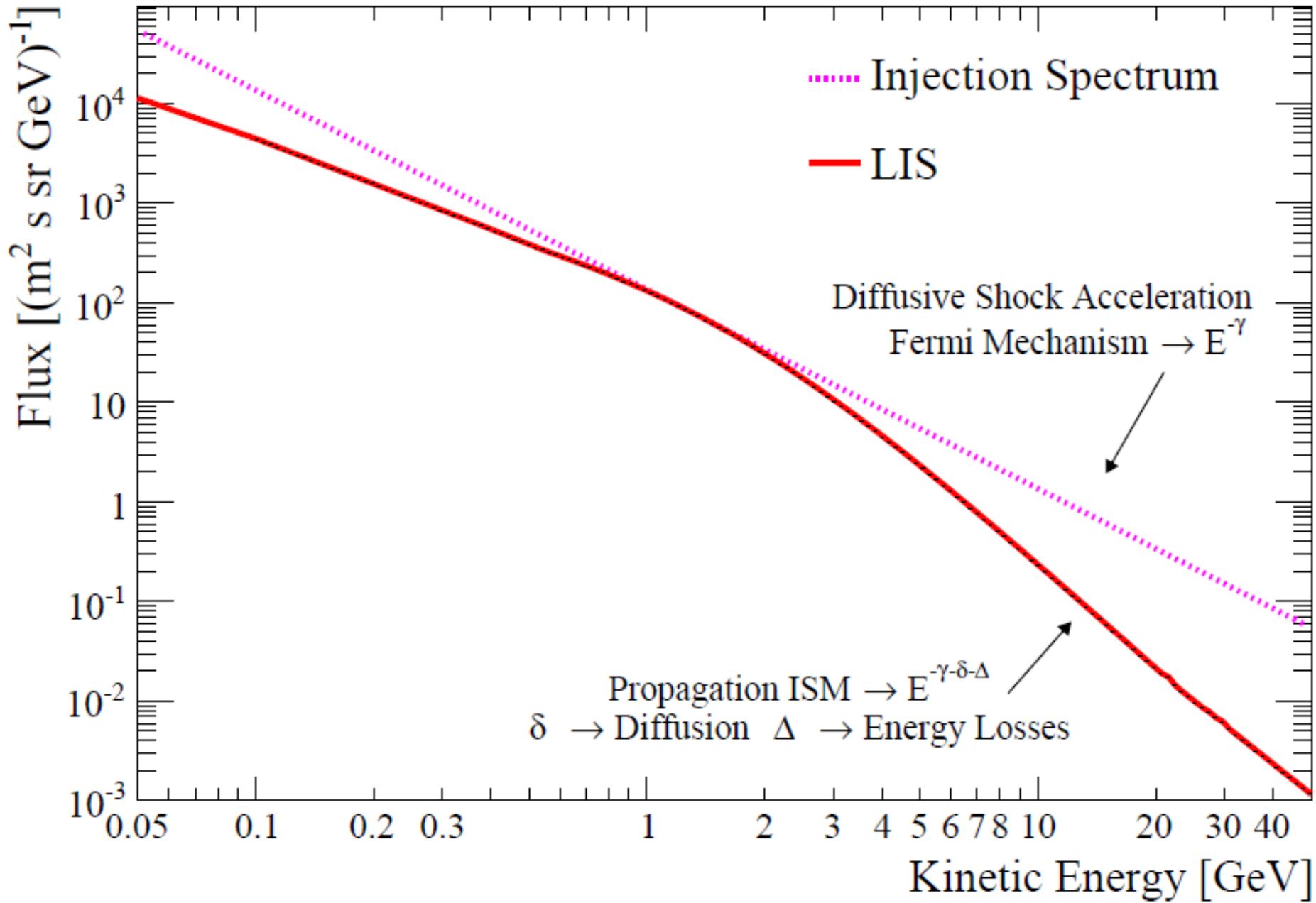


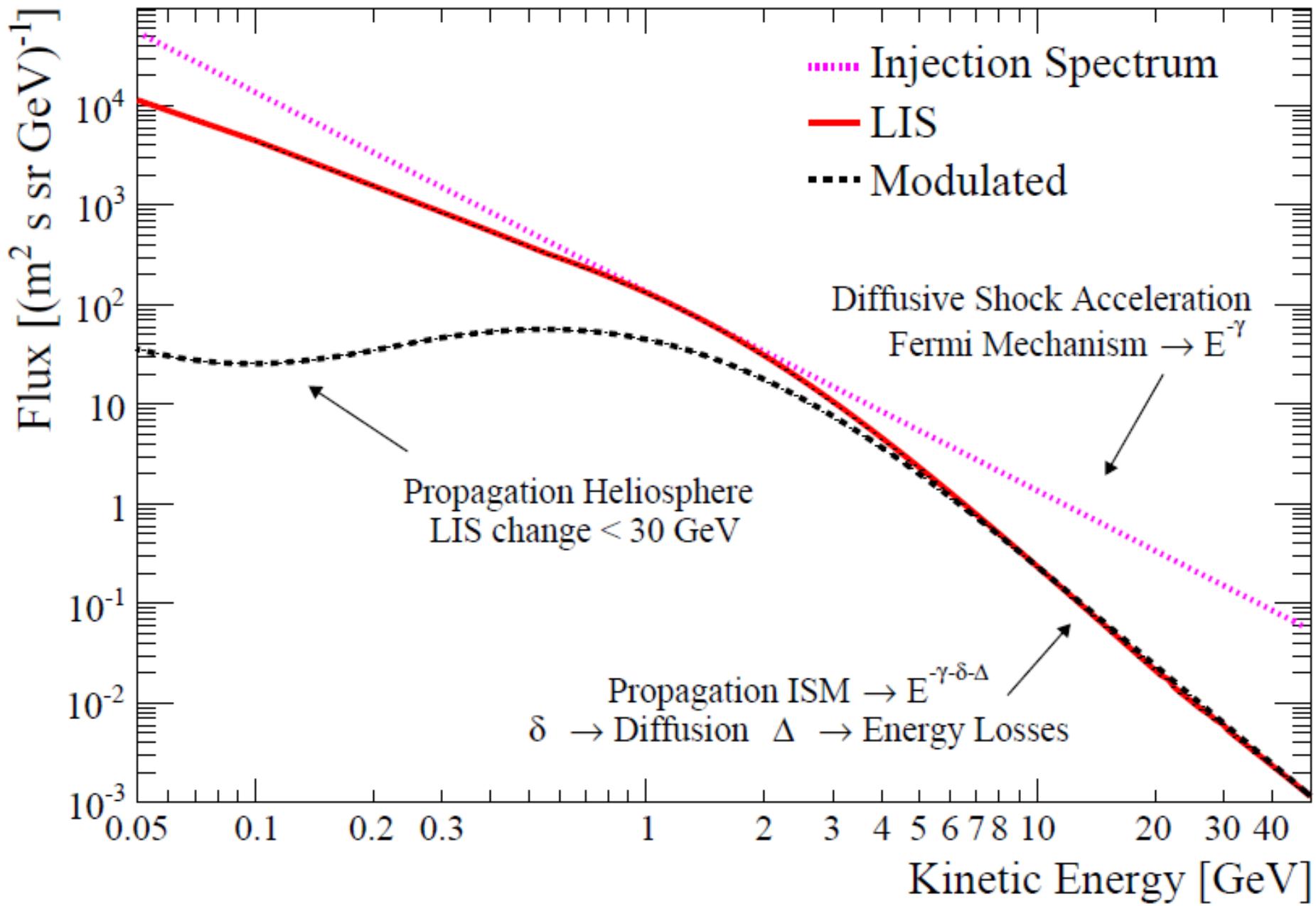
Adriani et al, Astropart. Phys. 34 (2010) 1;
arXiv:1001.3522 [astro-ph.HE]



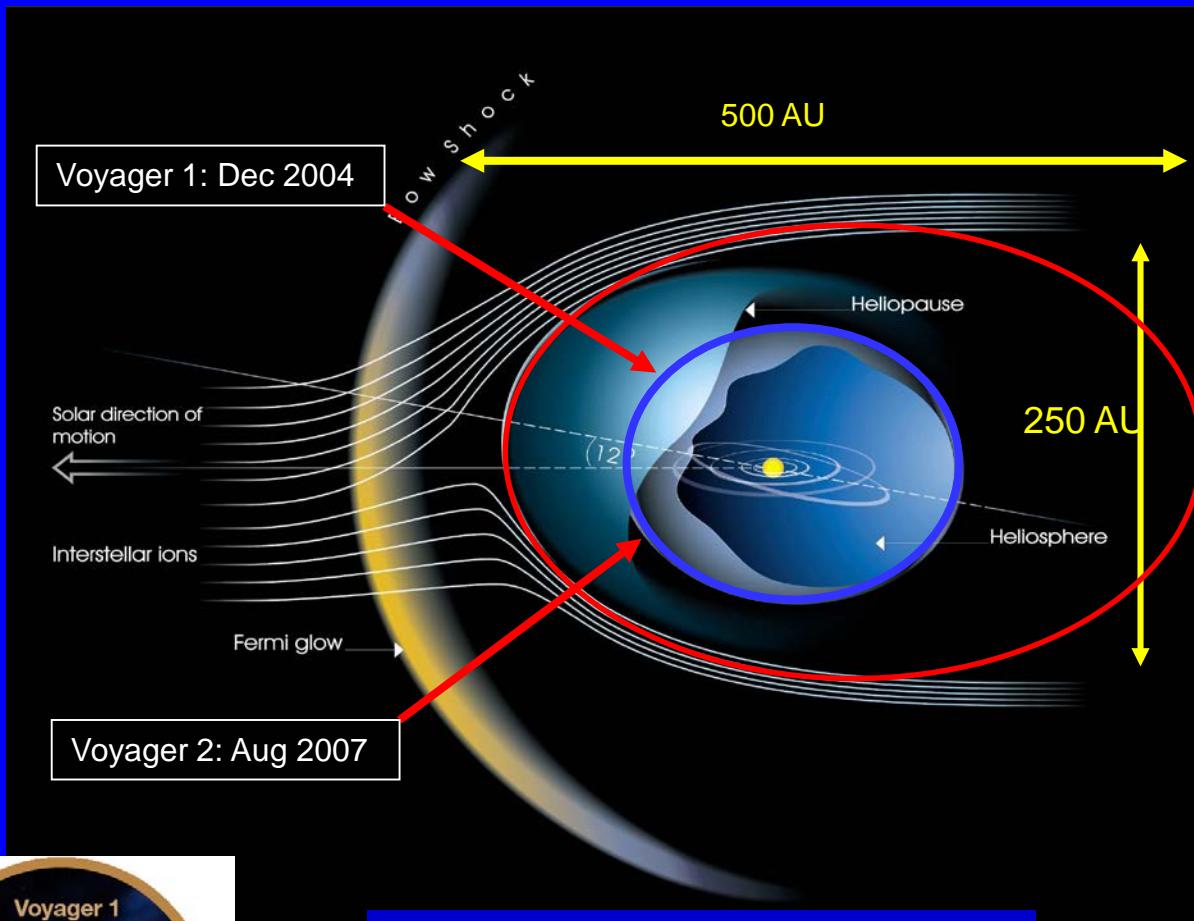
Cosmic rays in the heliosphere



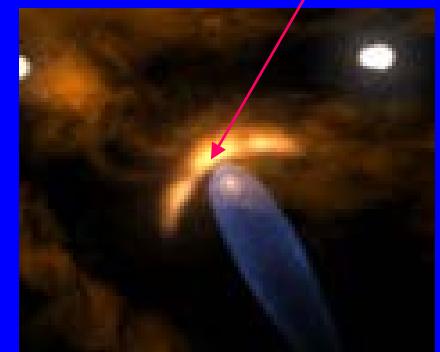




Heliospace: The Heliosphere



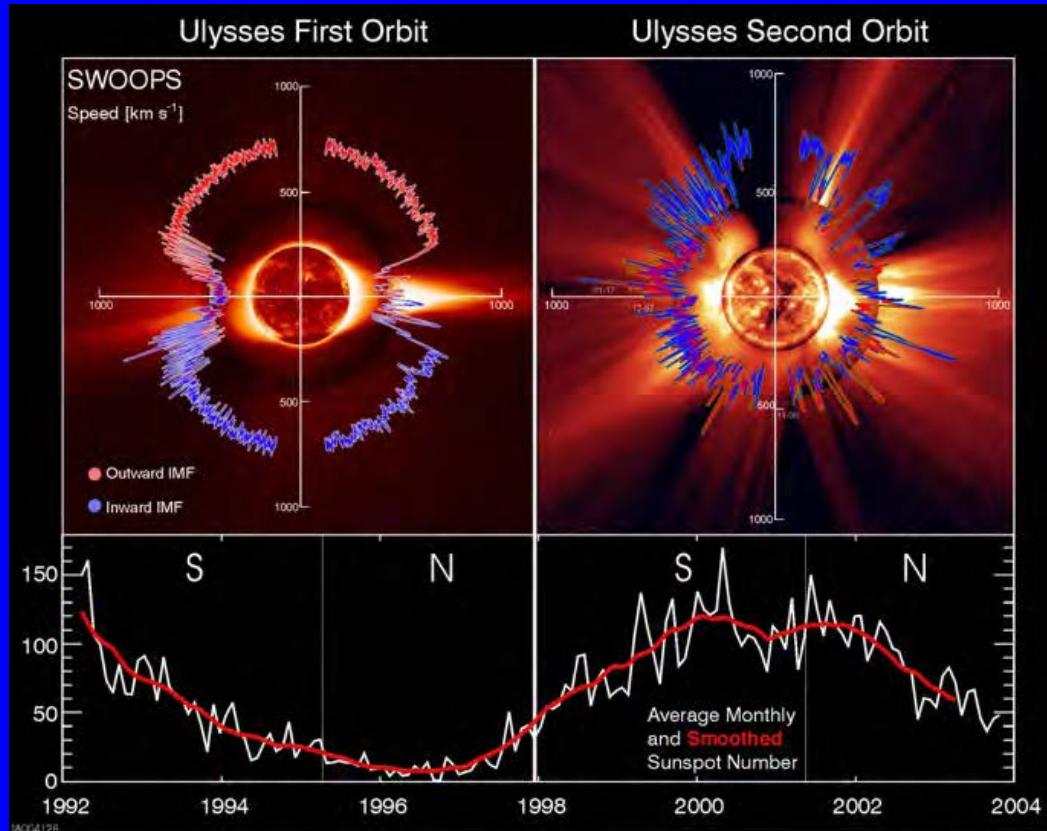
Bow Shocks



- Launched in 1977
- Voyager 1: ~ 121 AU
- Voyager 2: ~ 99 AU

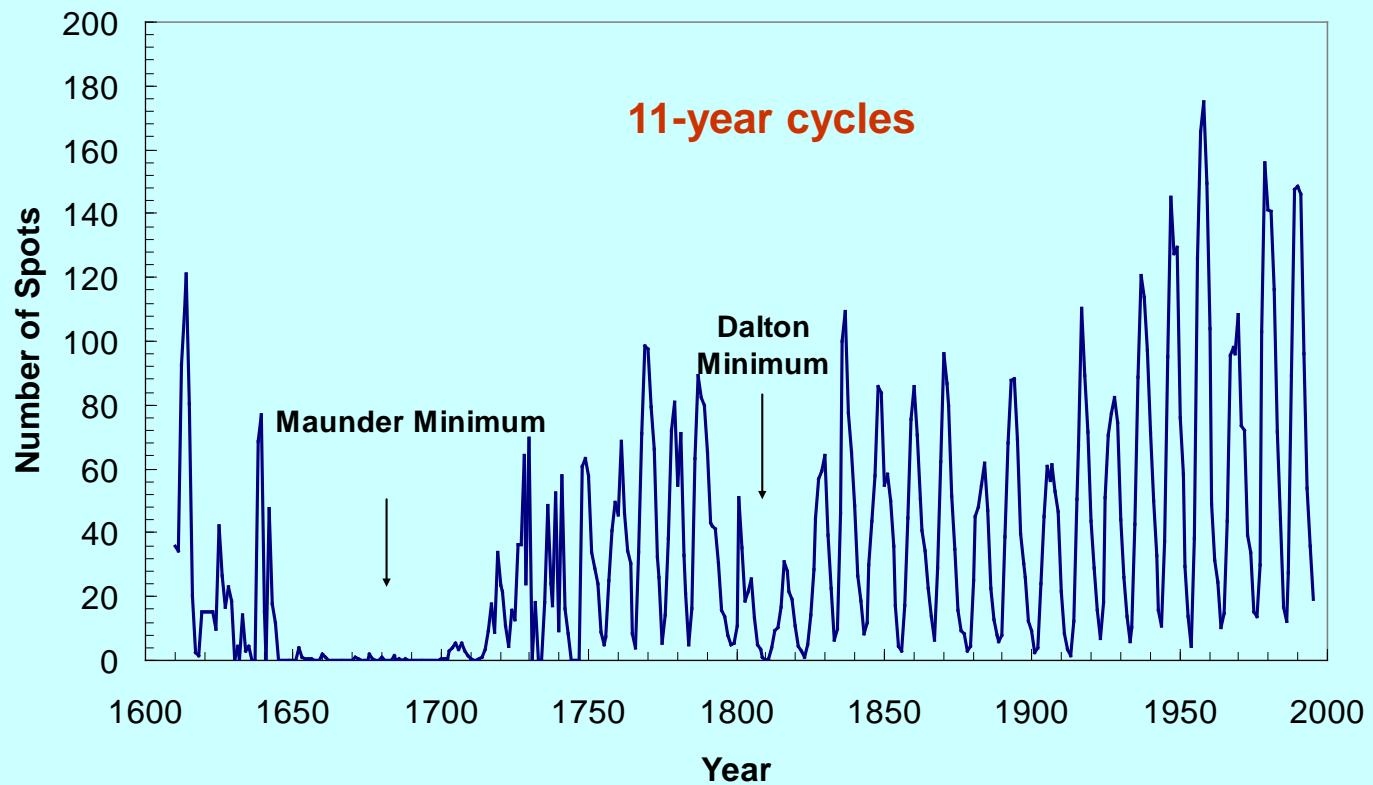
Courtesy of M. Potgieter

The Solar Wind

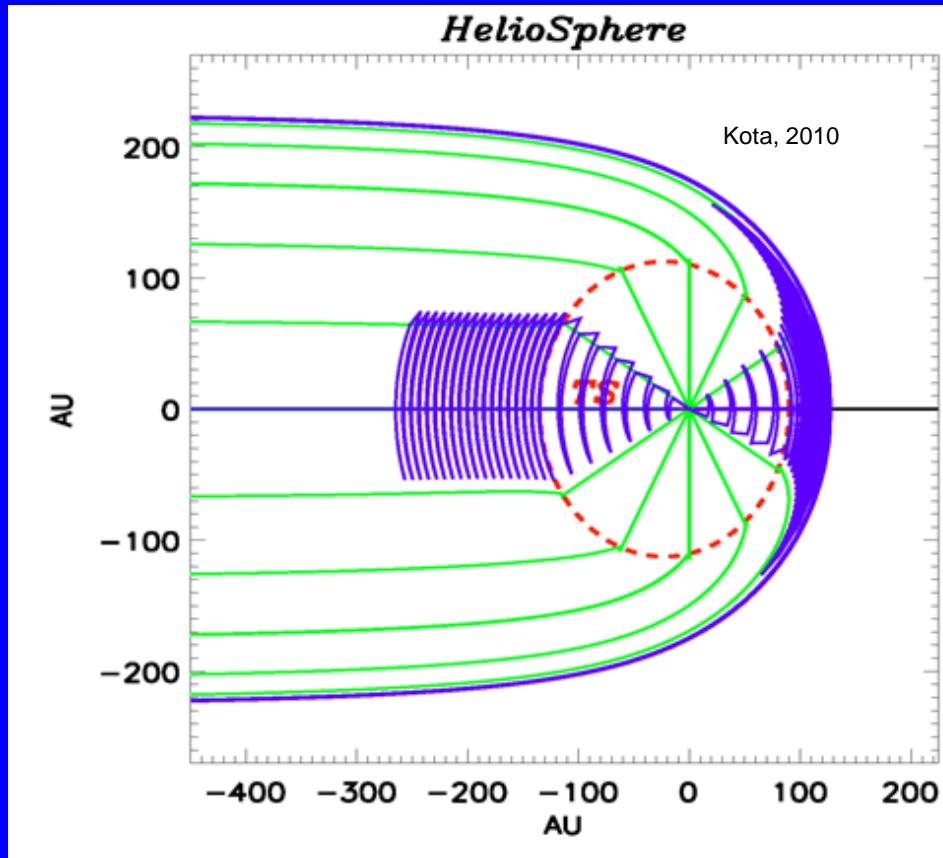


Strong latitude dependence at solar minimum

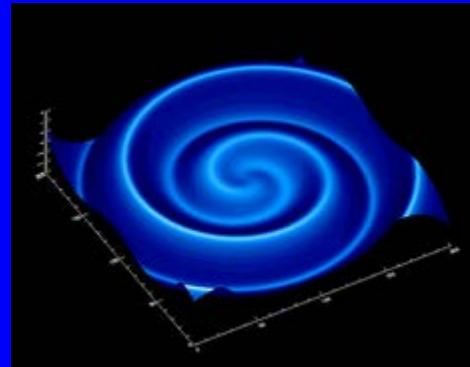
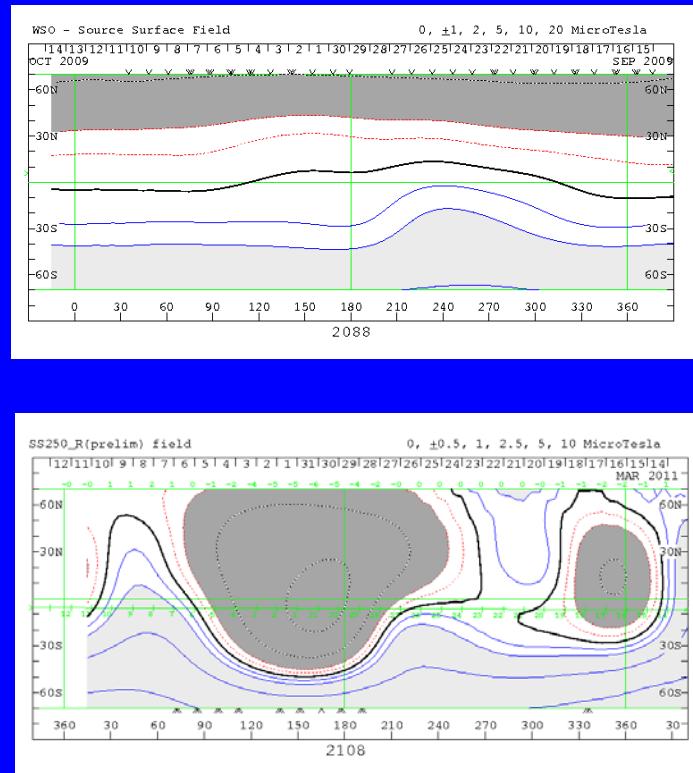
Solar Activity and Sunspot Numbers



Wavy Heliospheric Current Sheet



Conceptual wavy HCS



Magnetic flux freezing

- *The equation of continuity*

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 ,$$

where ρ is the mass density and \mathbf{v} is the velocity at a point in the fluid.

- *Force equation*

$$\rho \frac{d\mathbf{v}}{dt} = -\nabla p + \mathbf{J} \times \mathbf{B} + \mathbf{F}_v + \rho \mathbf{g} ,$$

where p is the pressure, \mathbf{J} is the current density, \mathbf{B} is the magnetic flux density, \mathbf{F}_v represents viscous forces and \mathbf{g} is the gravitational acceleration.

Maxwell's equations:

Slow varying phenomena -> no displacement term $\partial \mathbf{D} / \partial t$, hence no space charge effects

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} ,$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} ,$$

$$\nabla \cdot \mathbf{B} = 0 ,$$

$$\nabla \cdot \mathbf{E} = \frac{\rho_e}{\epsilon_0} .$$

Ohm's law

$$\mathbf{J} = \sigma(\mathbf{E} + \mathbf{v} \times \mathbf{B}),$$

where σ is the electrical conductivity of the plasma. Now substituting for \mathbf{E}

$$\nabla \times (\mathbf{J}/\sigma - \mathbf{v} \times \mathbf{B}) = -\frac{\partial \mathbf{B}}{\partial t}.$$

Now, eliminating \mathbf{J}

$$\nabla \times \left(\frac{\nabla \times \mathbf{B}}{\sigma \mu_0} - \mathbf{v} \times \mathbf{B} \right) = -\frac{\partial \mathbf{B}}{\partial t}.$$

Therefore,

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) - \frac{\nabla \times (\nabla \times \mathbf{B})}{\sigma \mu_0}.$$

We now use the identity $\nabla \times (\nabla \times \mathbf{B}) = \nabla(\nabla \cdot \mathbf{B}) - \nabla^2 \mathbf{B}$. Since $\nabla \cdot \mathbf{B}$ is always zero, we find

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \frac{1}{\sigma \mu_0} \nabla^2 \mathbf{B}.$$

Infinite conductivity $\sigma \rightarrow \infty$:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) . \quad (11.47)$$

consider a current loop S in the plasma and the two contributions to changes in the magnetic flux density ϕ through it with time. First, there may be changes in the magnetic flux density due to external causes, and second, there is an induced component of the flux density due to motion of the loop. The first contribution is

$$\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} . \quad (11.48)$$

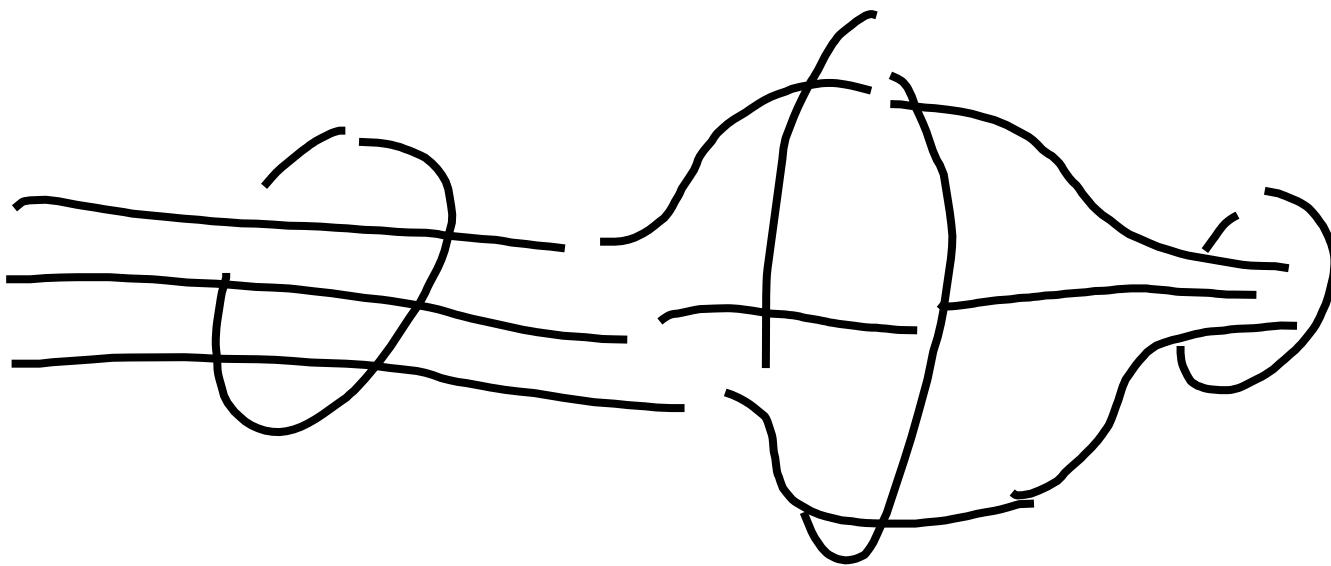
The second contribution results from the fact that, because of the motion of the loop, there is an induced electric field $E = \mathbf{v} \times \mathbf{B}$. Because $\nabla \times \mathbf{E} = -\partial \phi / \partial t$, there is an additional contribution to the total magnetic flux through the loop,

$$- \int_S \nabla \times (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{S} . \quad (11.49)$$

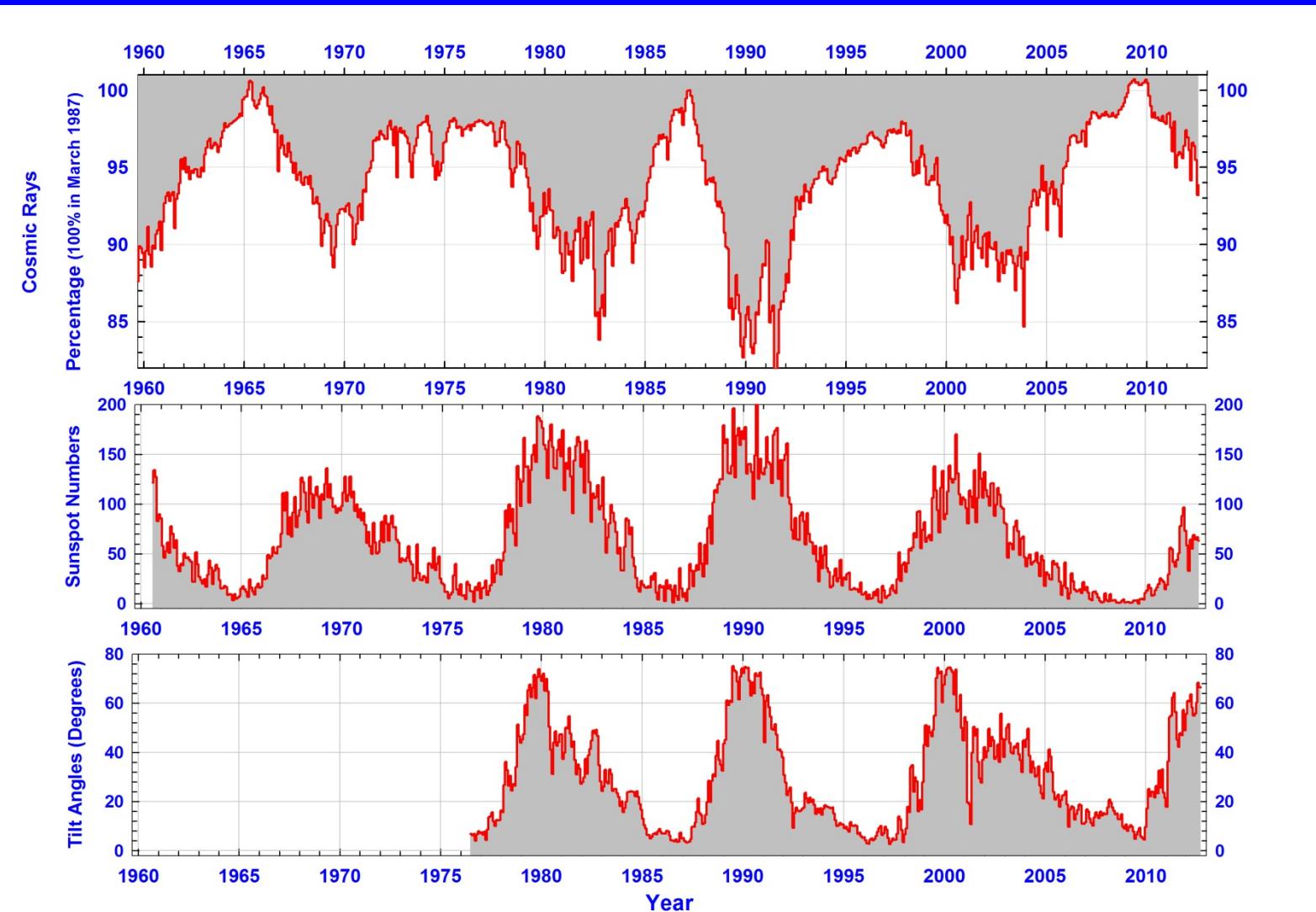
Adding together both contributions, we obtain

$$\begin{aligned} \frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S} &= \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} - \int_S \nabla \times (\mathbf{v} \times \mathbf{B}) \cdot d\mathbf{S} \\ &= \int_S \left(\frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{v} \times \mathbf{B}) \right) \cdot d\mathbf{S} = 0 , \end{aligned}$$

$$\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S} = 0$$

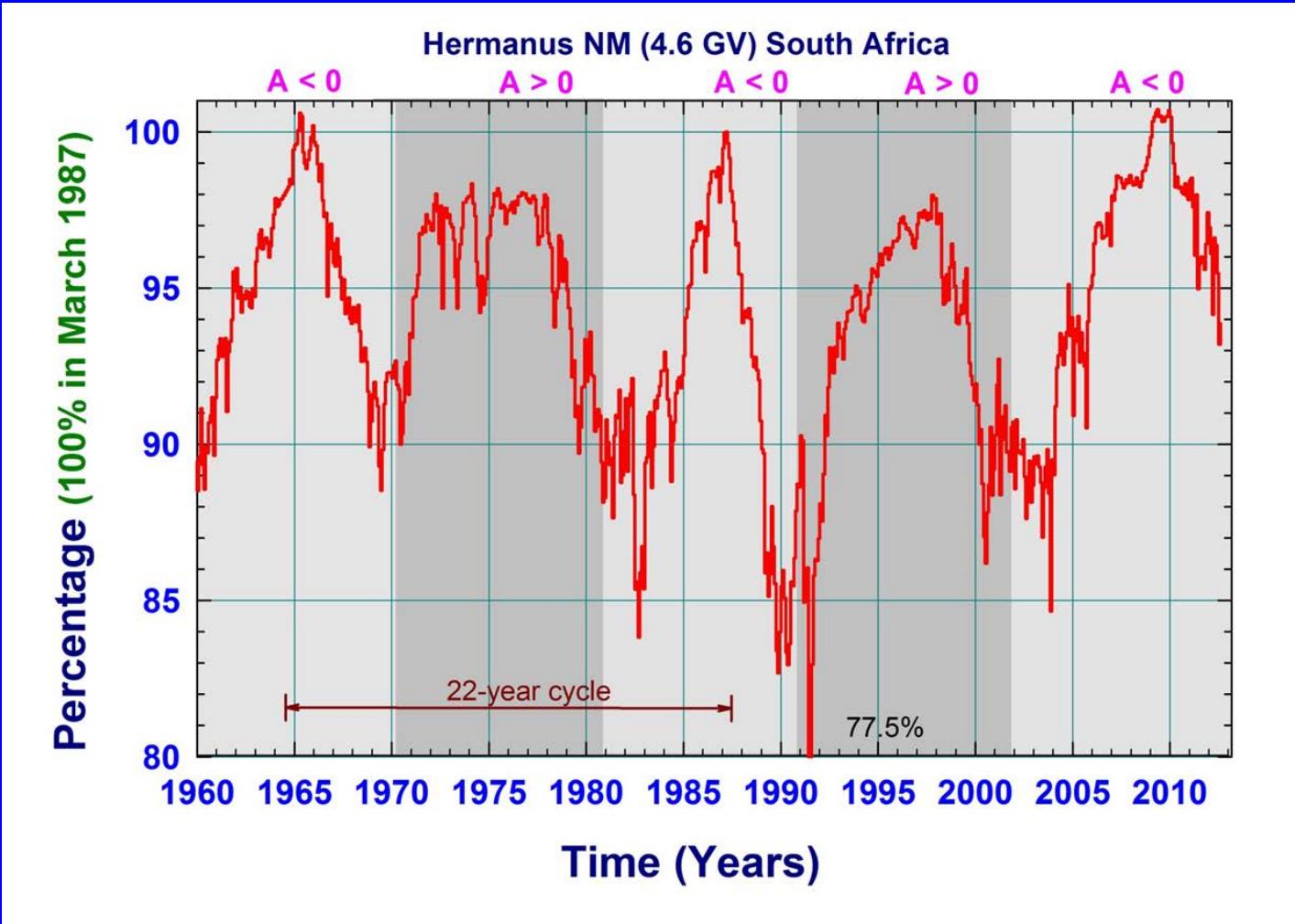


Solar Activity and Cosmic Rays



Cosmic rays as indicators of heliospheric conditions

Modulation of galactic CRs at Earth at NM energies ($E > 10\text{GV}$)



Transport equation for the transport, modulation and acceleration of cosmic rays in the heliosphere

$$\frac{\partial f}{\partial t} = \nabla \cdot [\mathbf{K} \cdot \nabla f] - \mathbf{V} \cdot \nabla f - \langle \mathbf{v}_D \rangle \cdot \nabla f + \frac{1}{3} (\nabla \cdot \mathbf{V}) \frac{\partial f}{\partial \ln p} + Q(r, p, t)$$

Time-dependent, pitch-angle-averaged distribution function

Diffusion

Convection with solar wind

Particle Drifts

Adiabatic energy changes

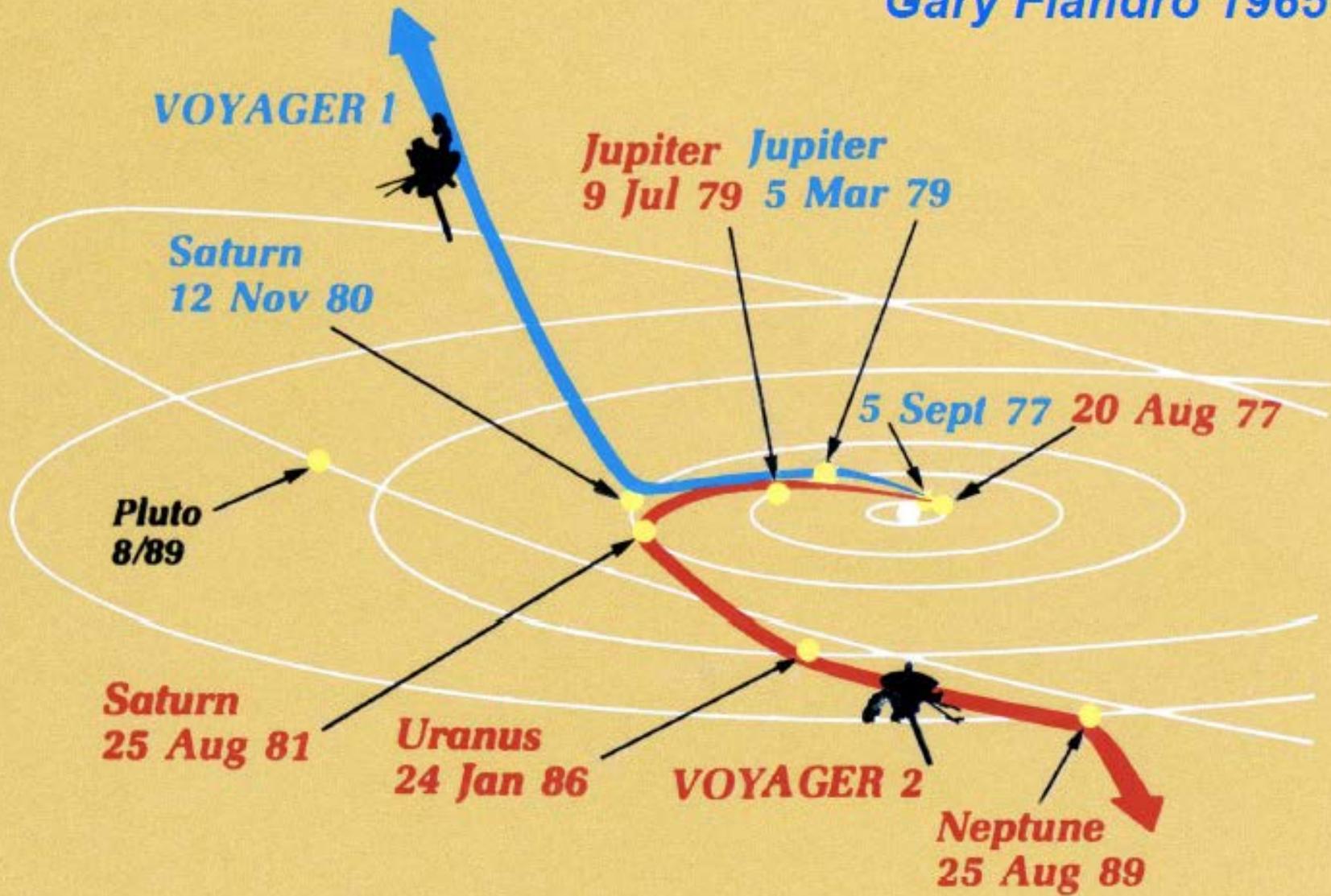
Any local source

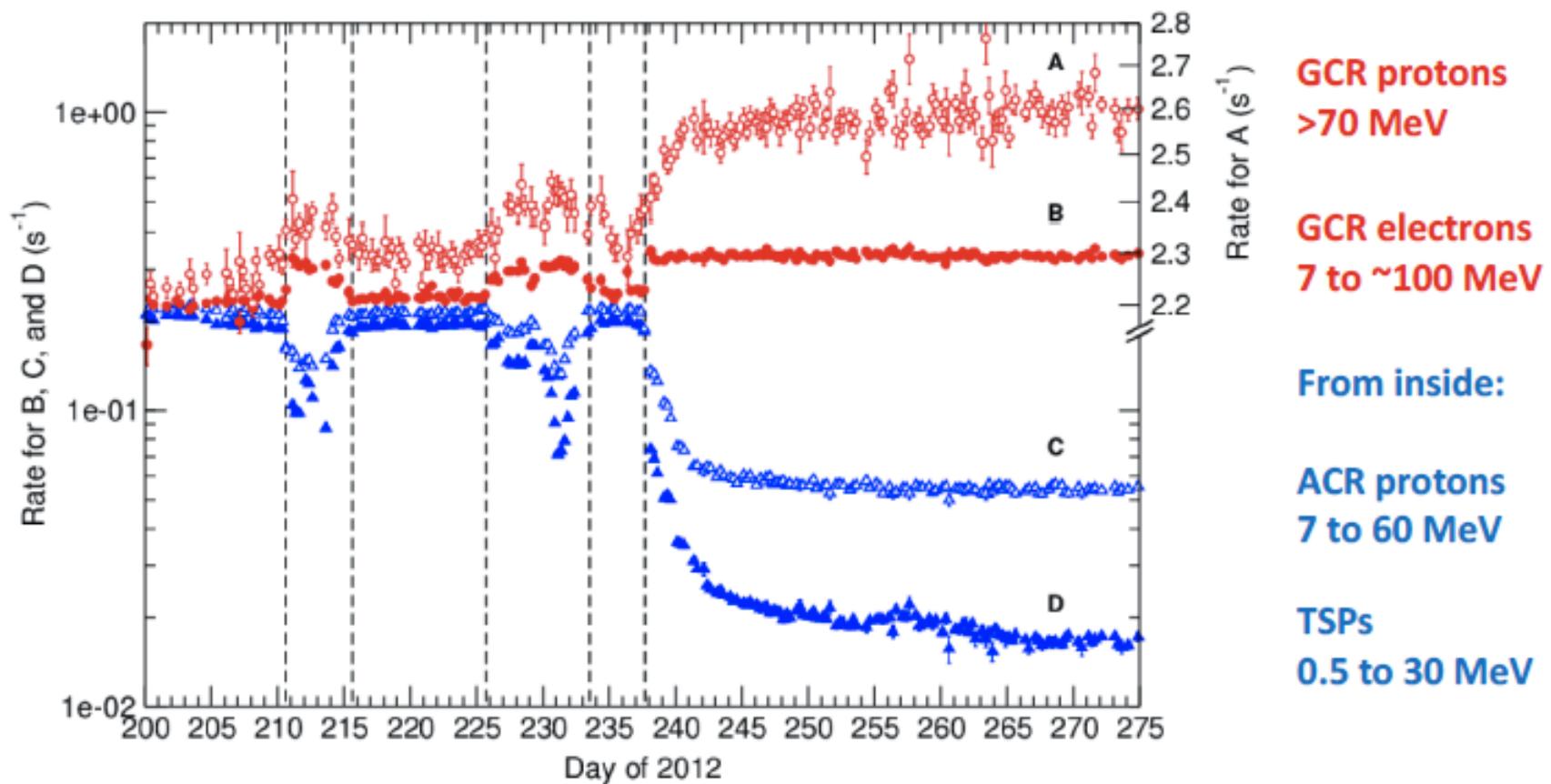
$$\dots = \dots + \frac{I}{p^2} \frac{\partial}{\partial p} \left(p^2 D_{pp} \frac{\partial f}{\partial p} \right)$$

Second order Fermi acceleration

Parker (Planet. Space Science, 13, 9, 1965)

Gary Flandro 1965





From outside:

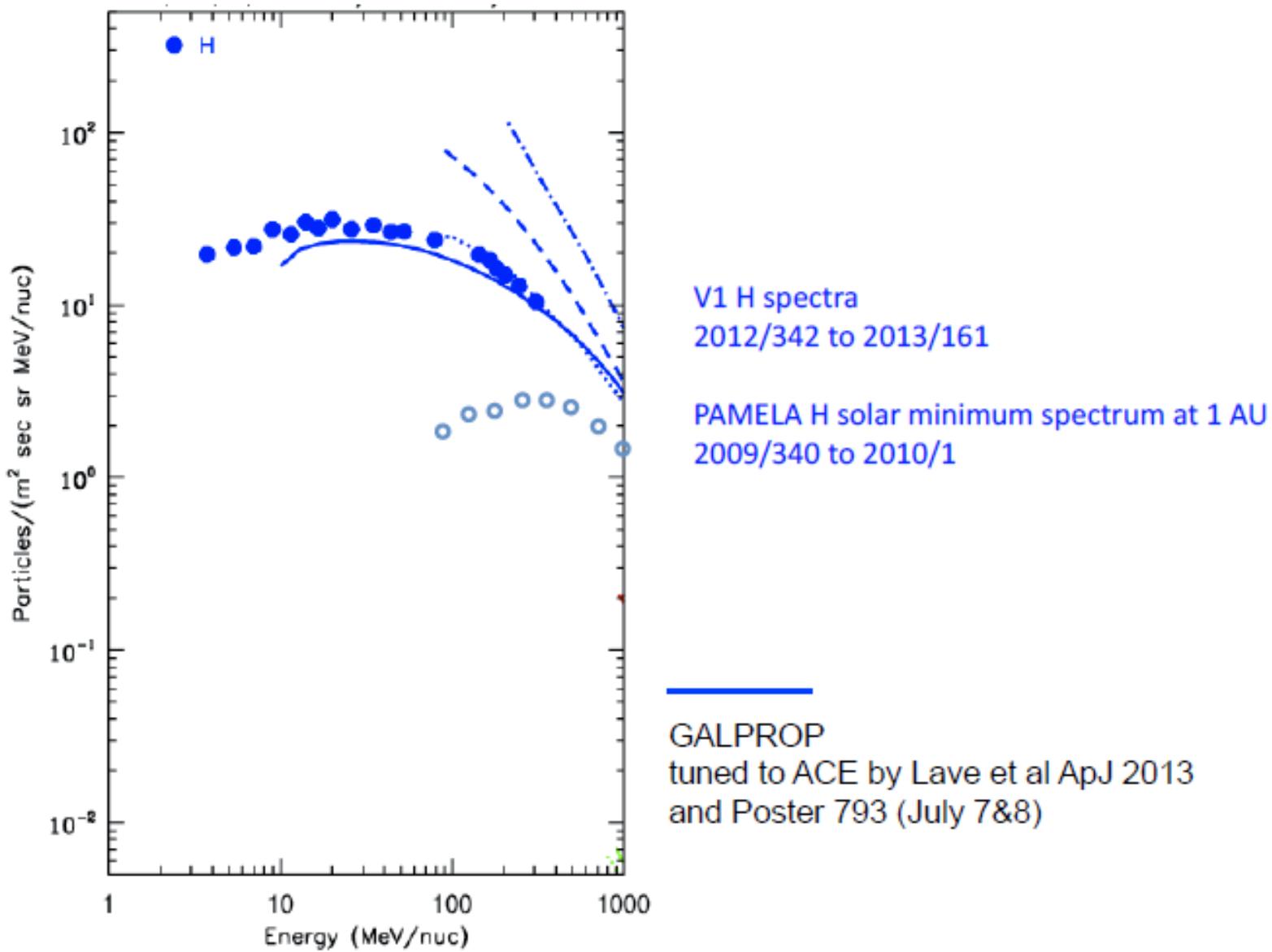
GCR protons
>70 MeV

GCR electrons
7 to ~100 MeV

From inside:

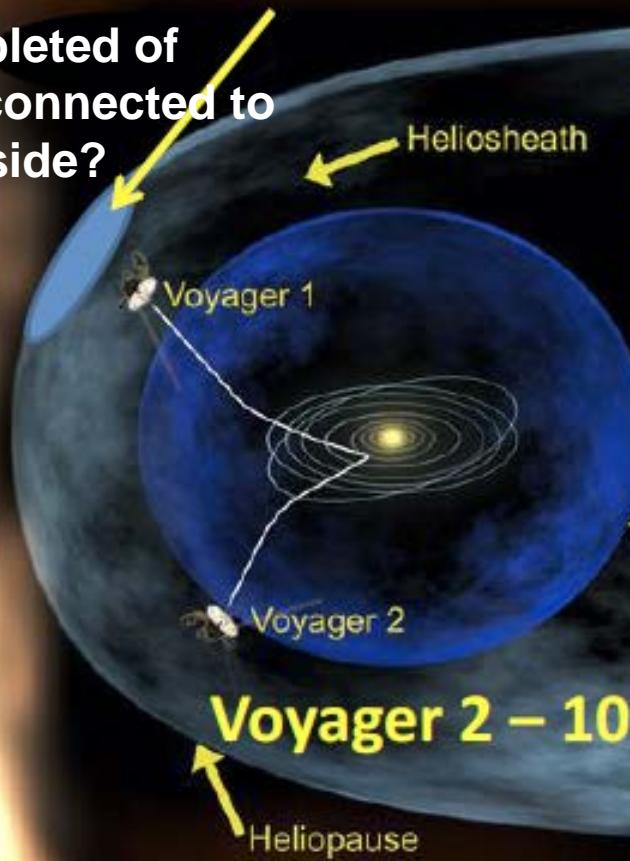
ACR protons
7 to 60 MeV

TSPs
0.5 to 30 MeV



Voyager 1 – 124 AU; 18.6 billion km

Depletion Region: depleted of heliospheric ions and connected to interstellar space outside?

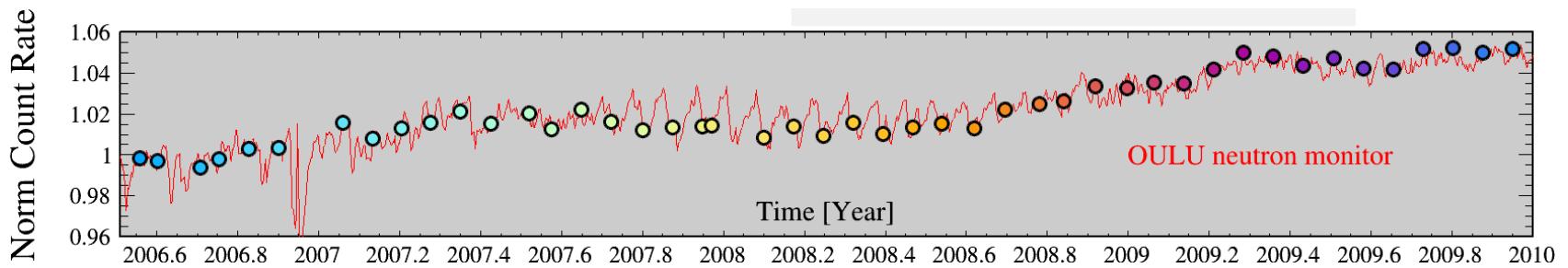
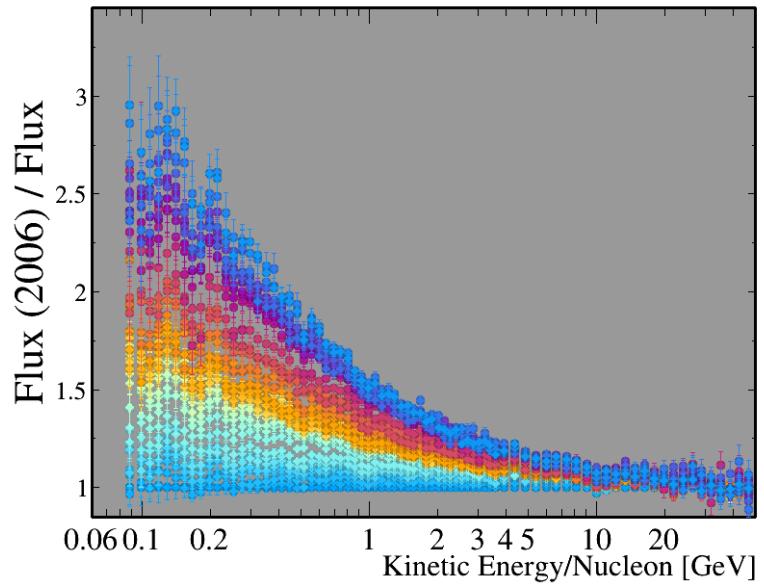
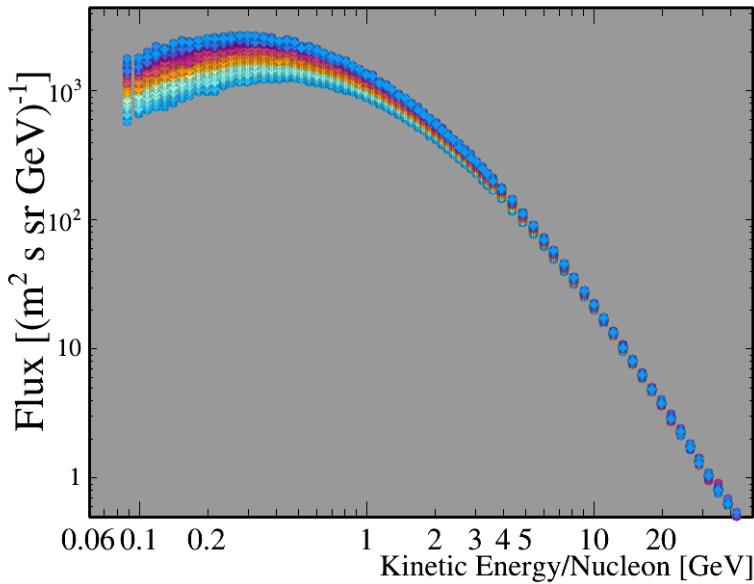


**Power all instruments until 2020 – 150 AU
Turn off final instrument in 2025**

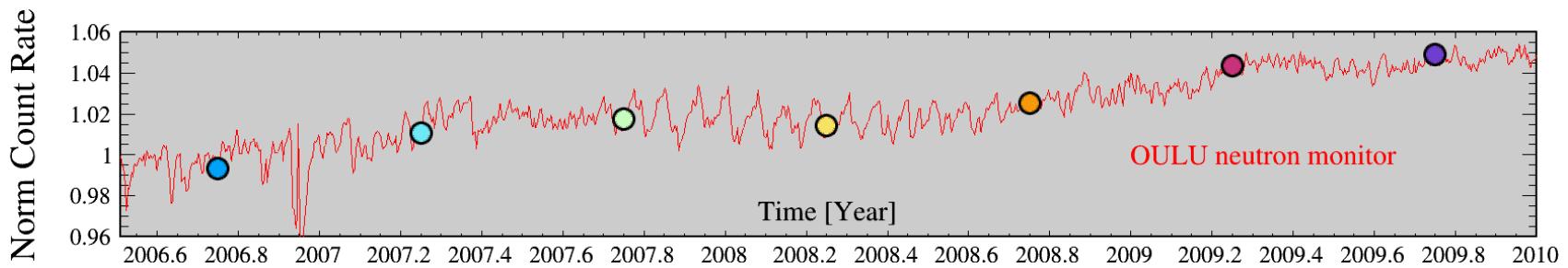
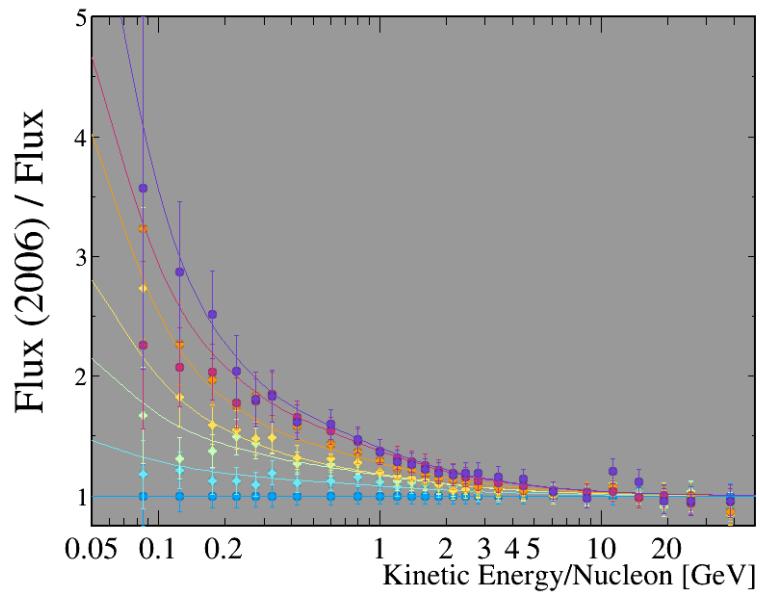
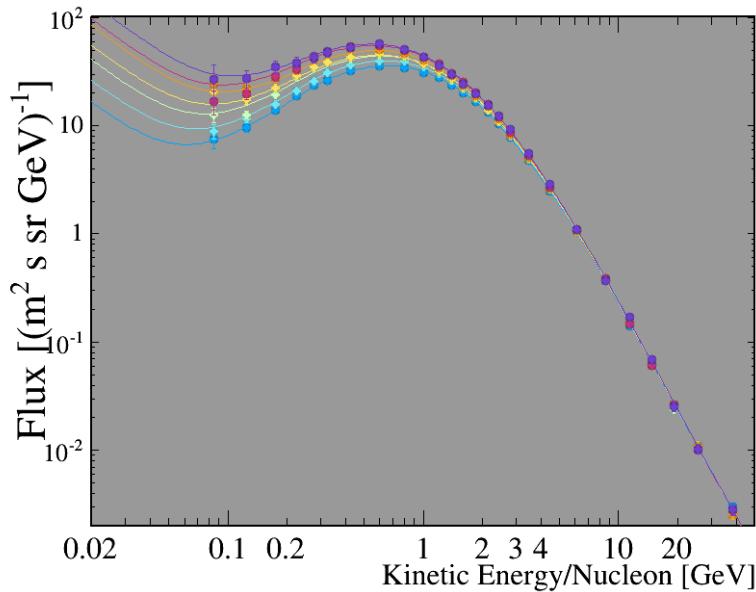
Voyager 2 – 102 AU; 15.2 billion km

Heliosphere

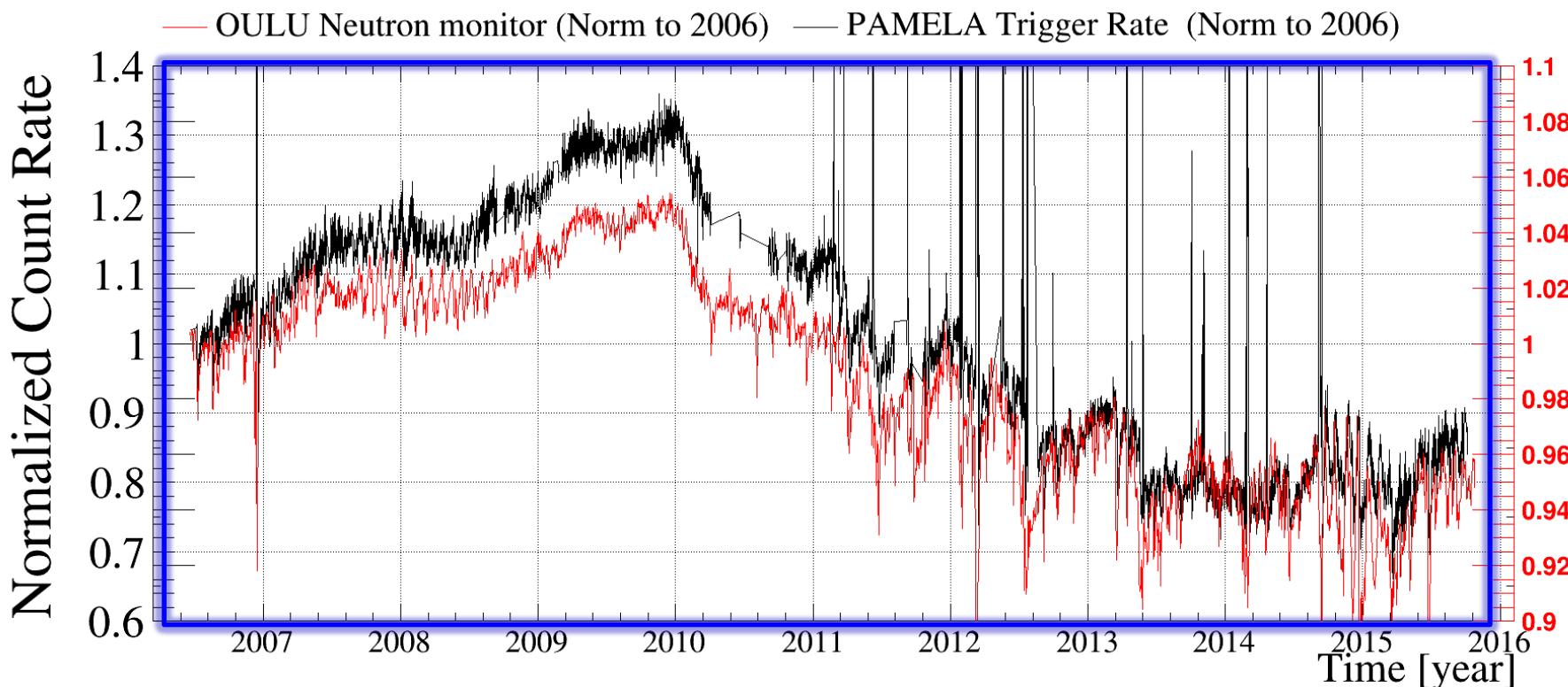
Time dependence - Proton flux



Time dependence - Electron flux



Heliospheric conditions during PAMELA observations

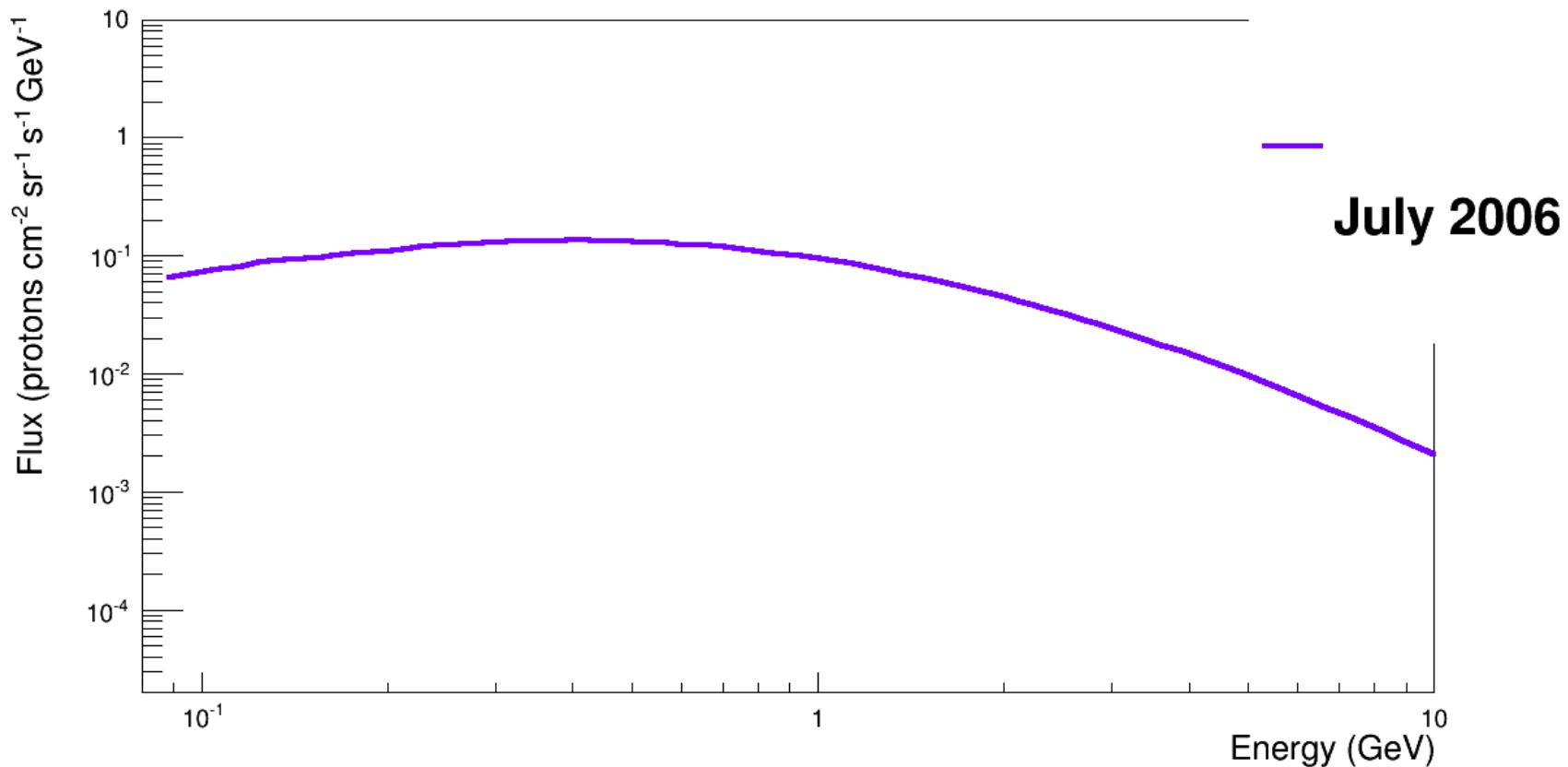


Neutron Monitor counts data from
<http://cosmicrays.oulu.fi/>

PAMELA observations covers ~ one solar cycle

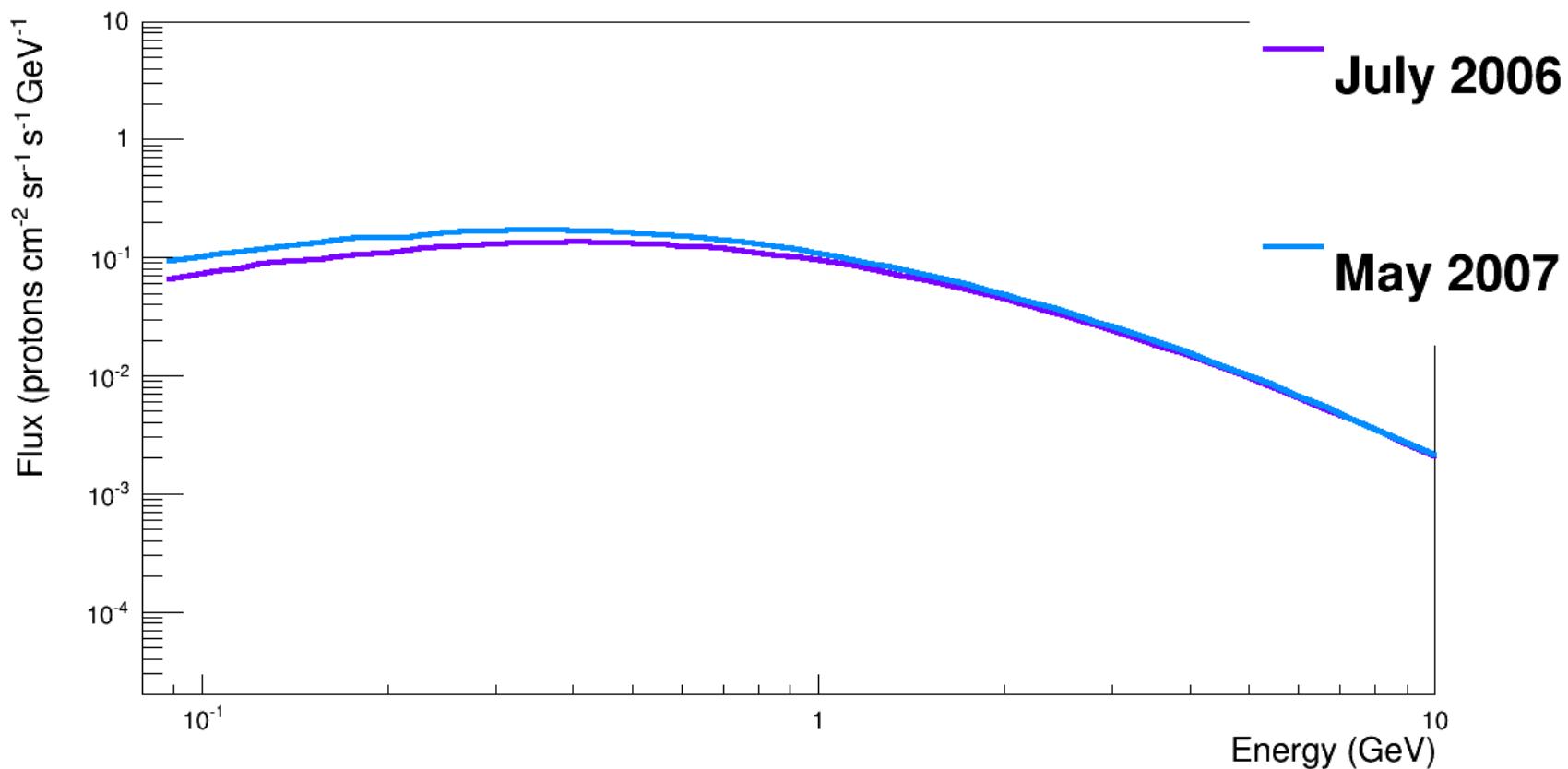
Time dependance of the proton flux

July 2006-January 2014



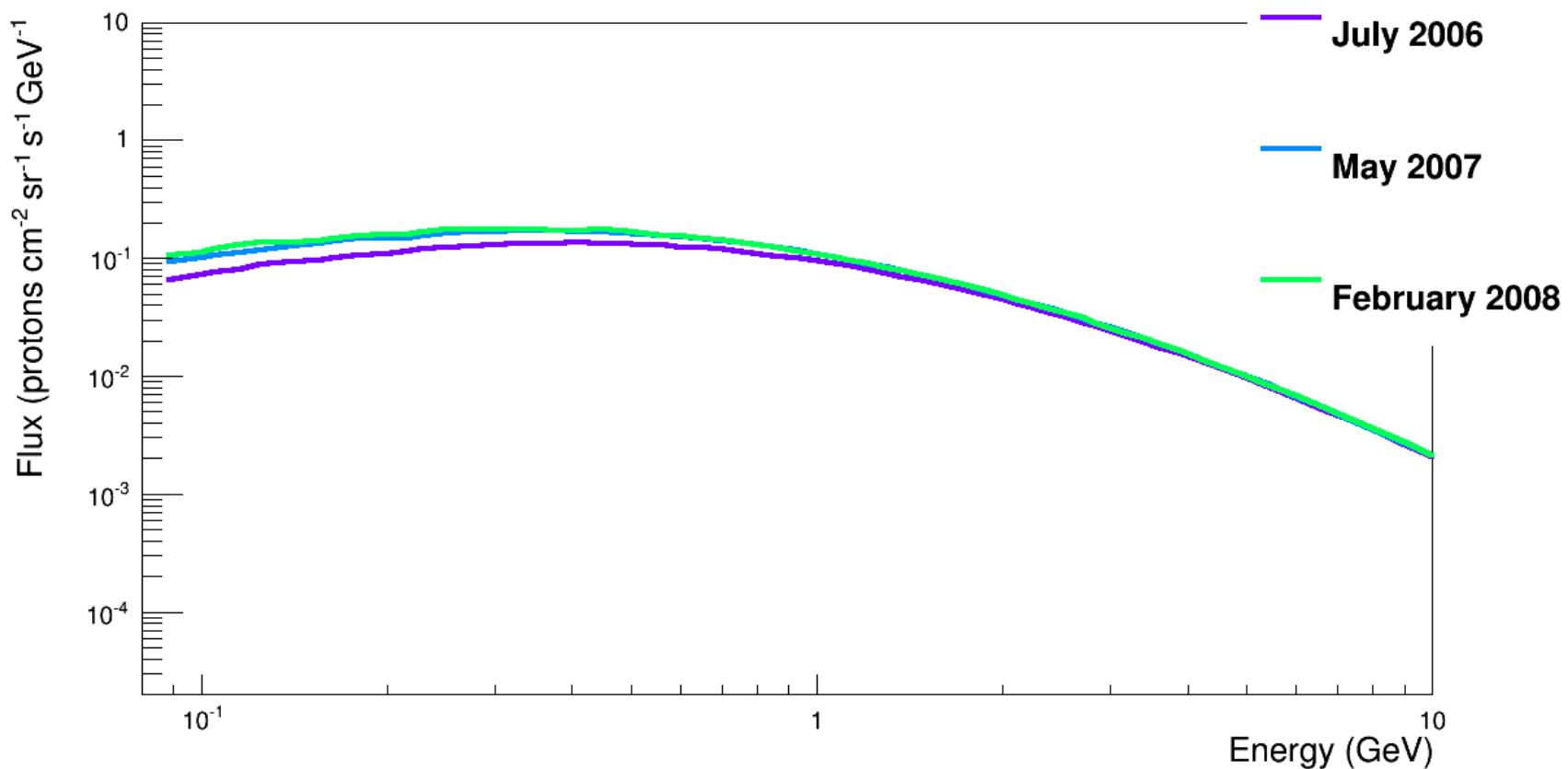
Time dependance of the proton flux

July 2006-January 2014



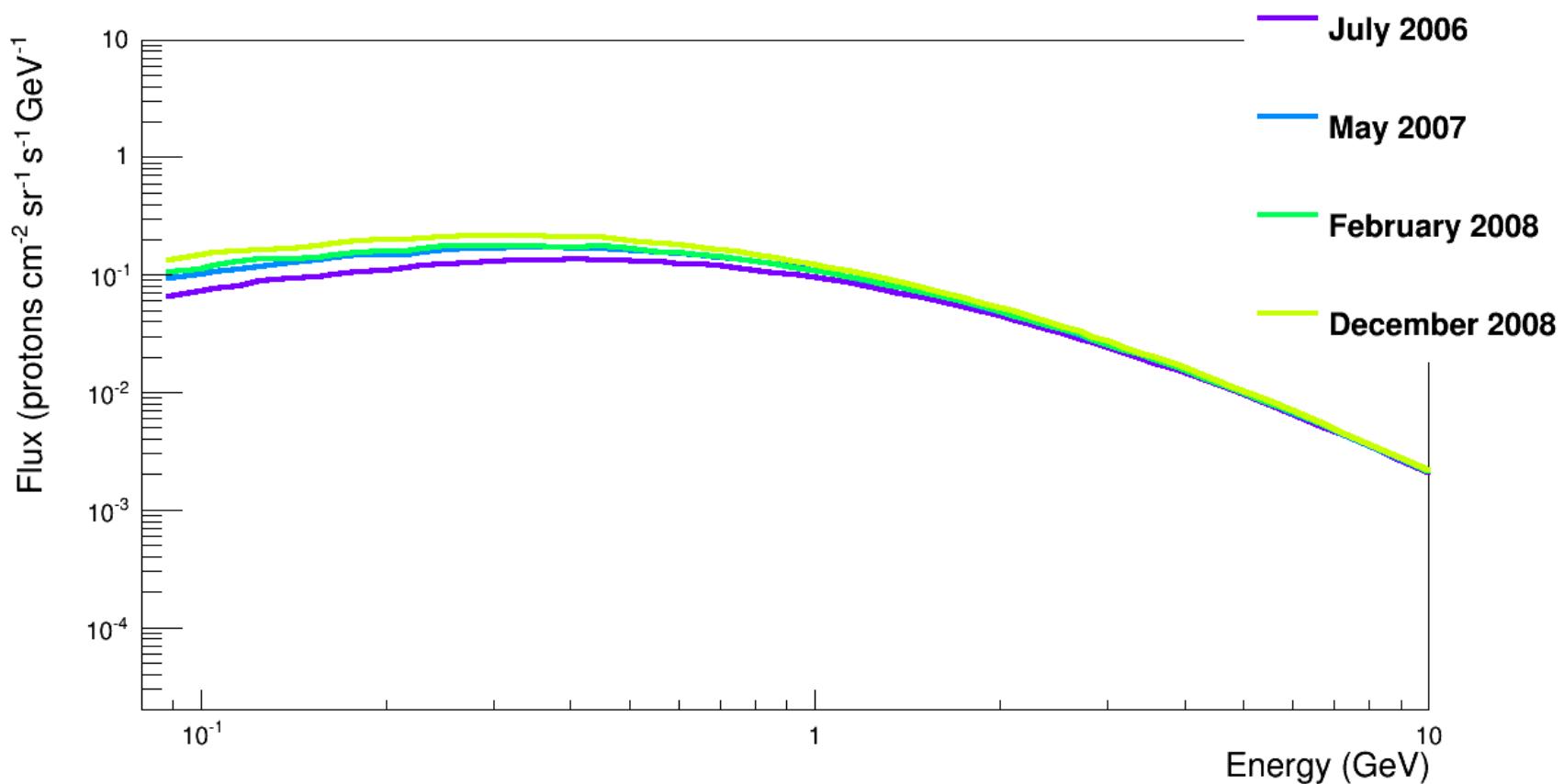
Time dependance of the proton flux

July 2006-January 2014



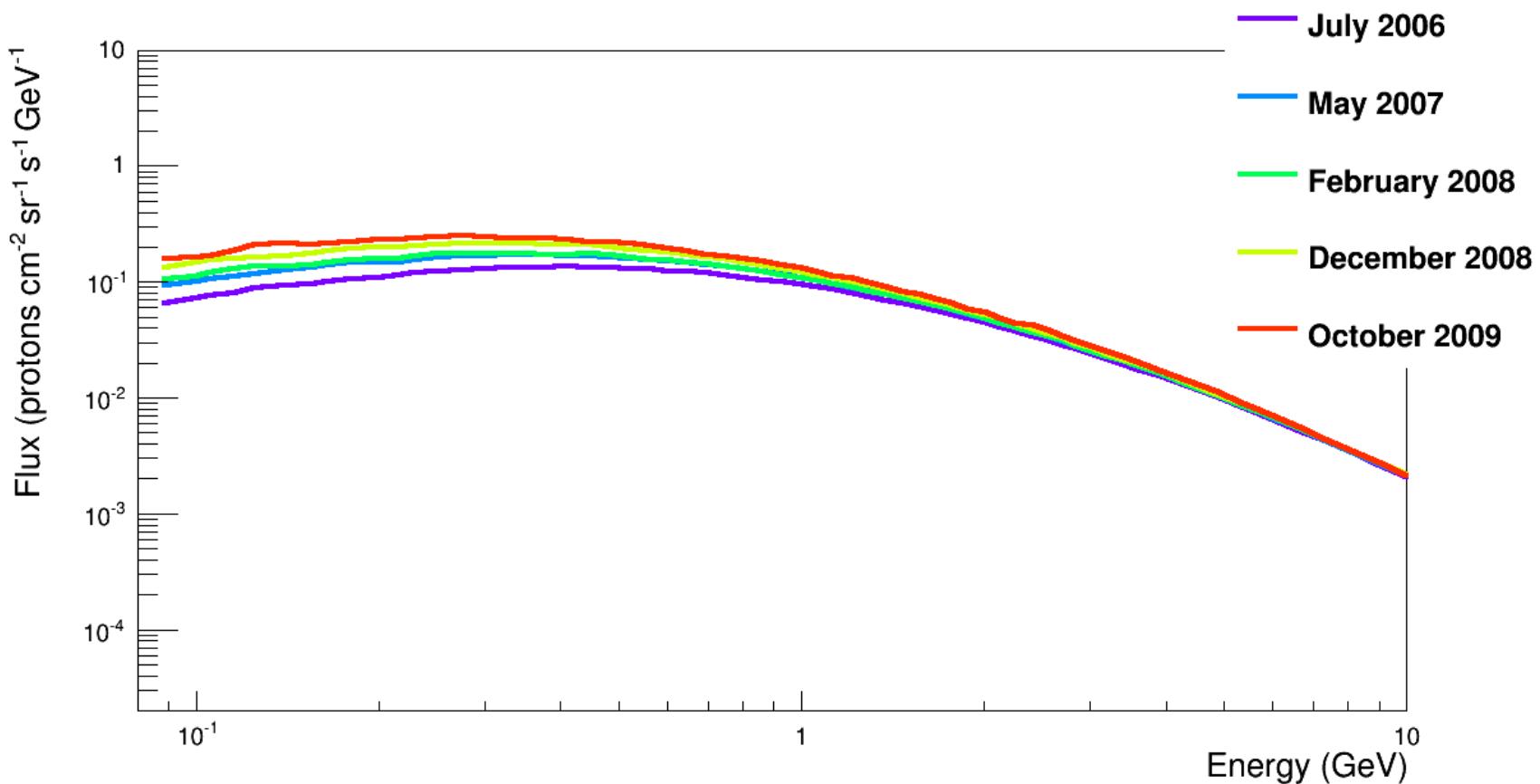
Time dependance of the proton flux

July 2006-January 2014



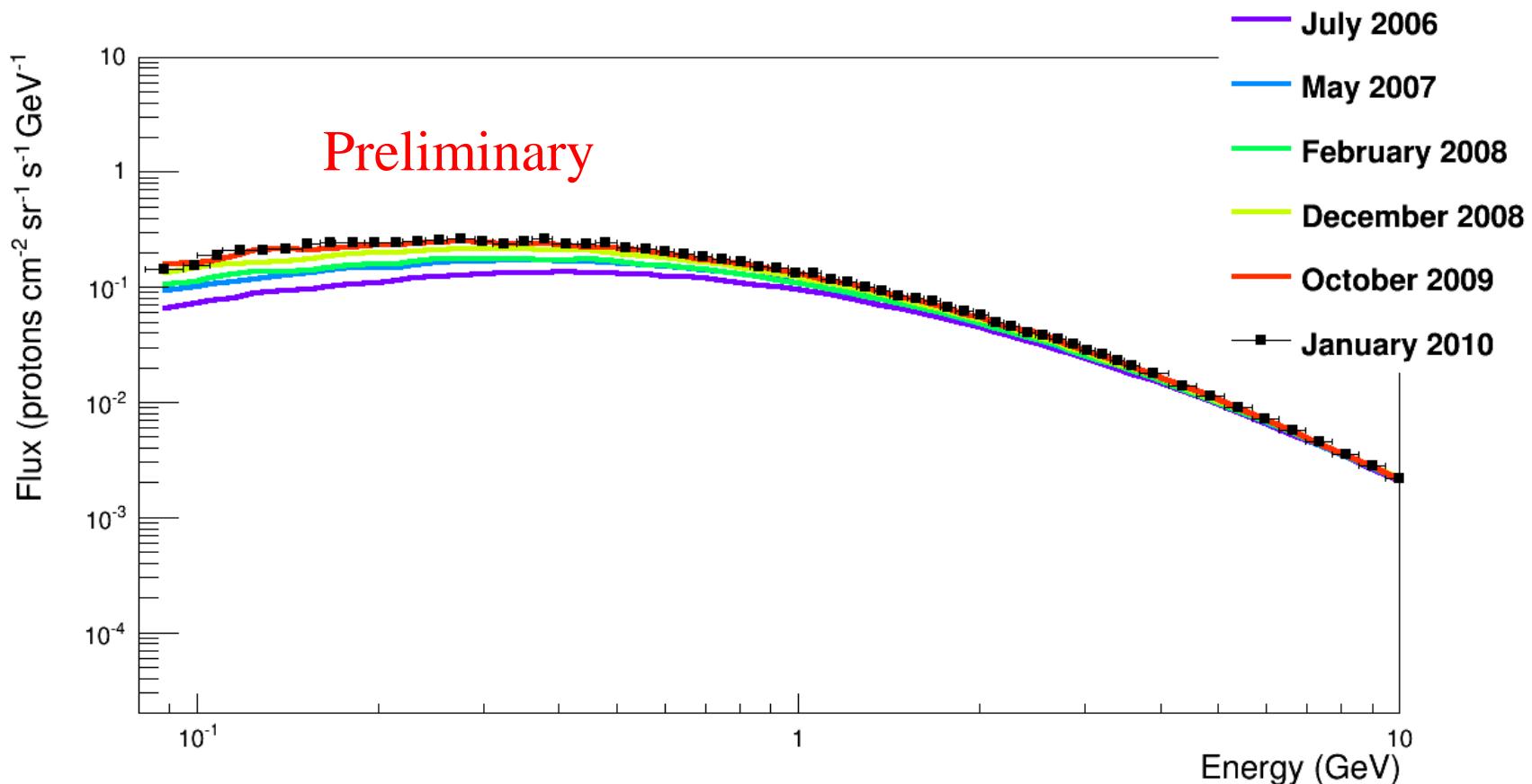
Time dependance of the proton flux

July 2006-January 2014



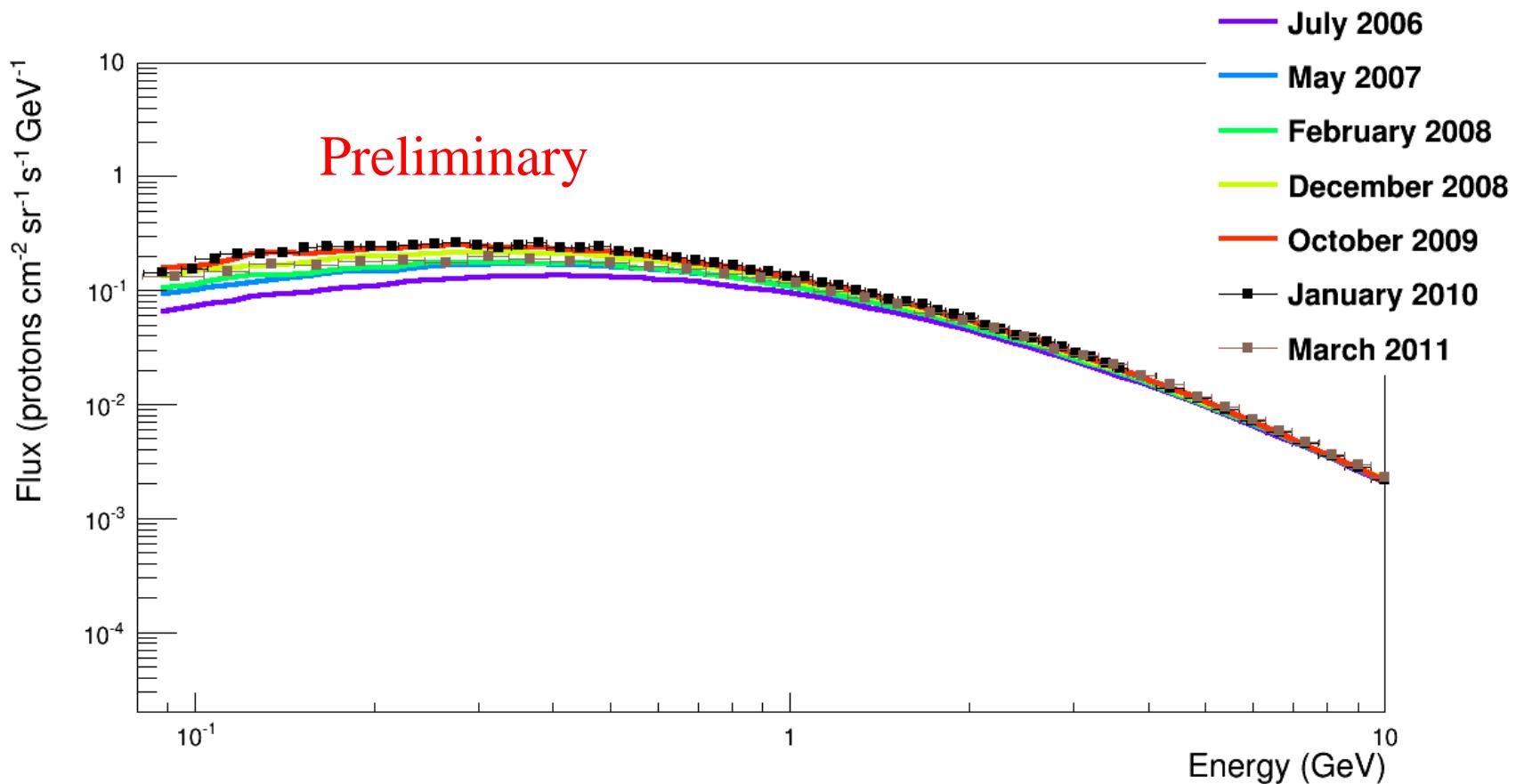
Time dependance of the proton flux

July 2006-January 2014



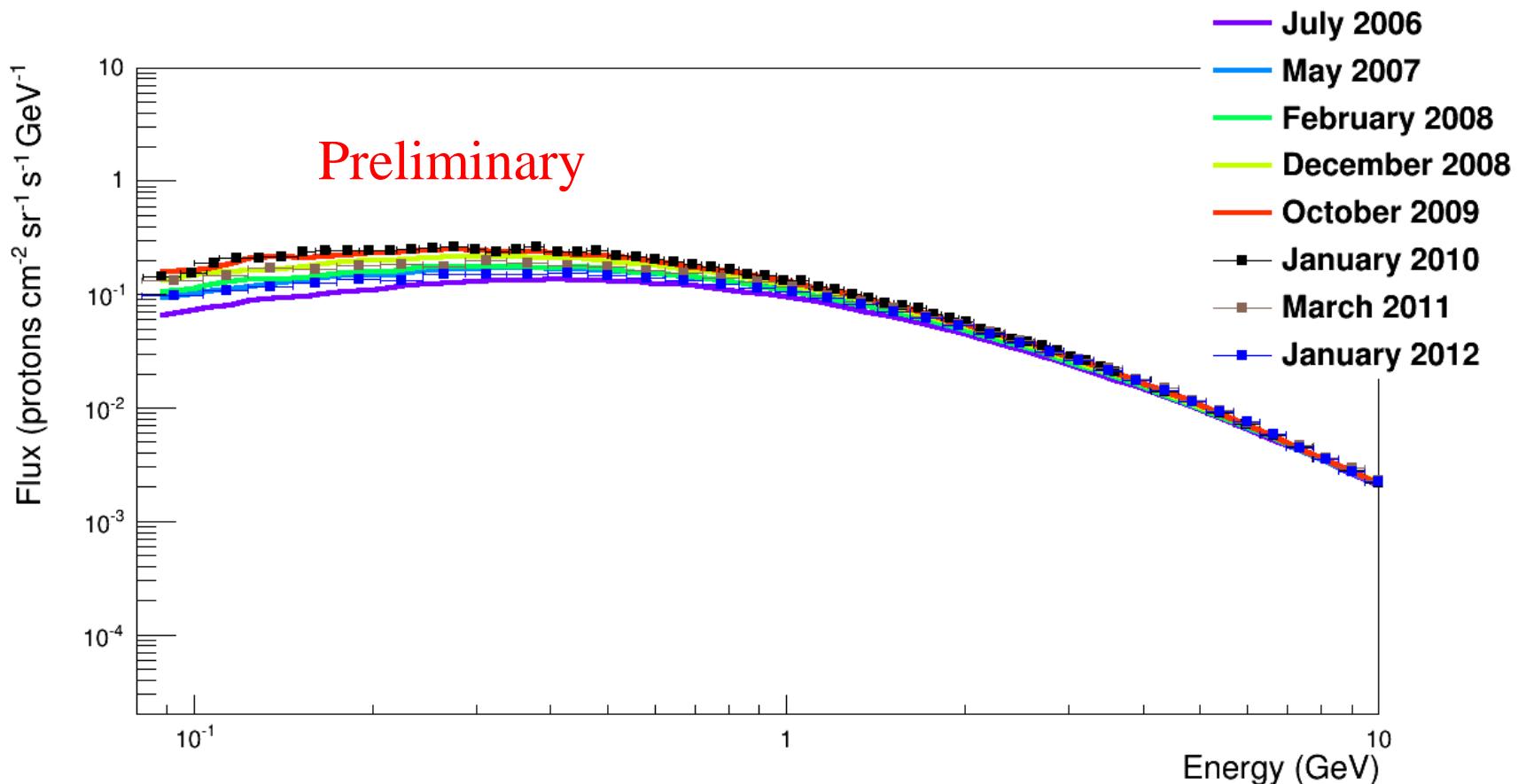
Time dependance of the proton flux

July 2006-January 2014



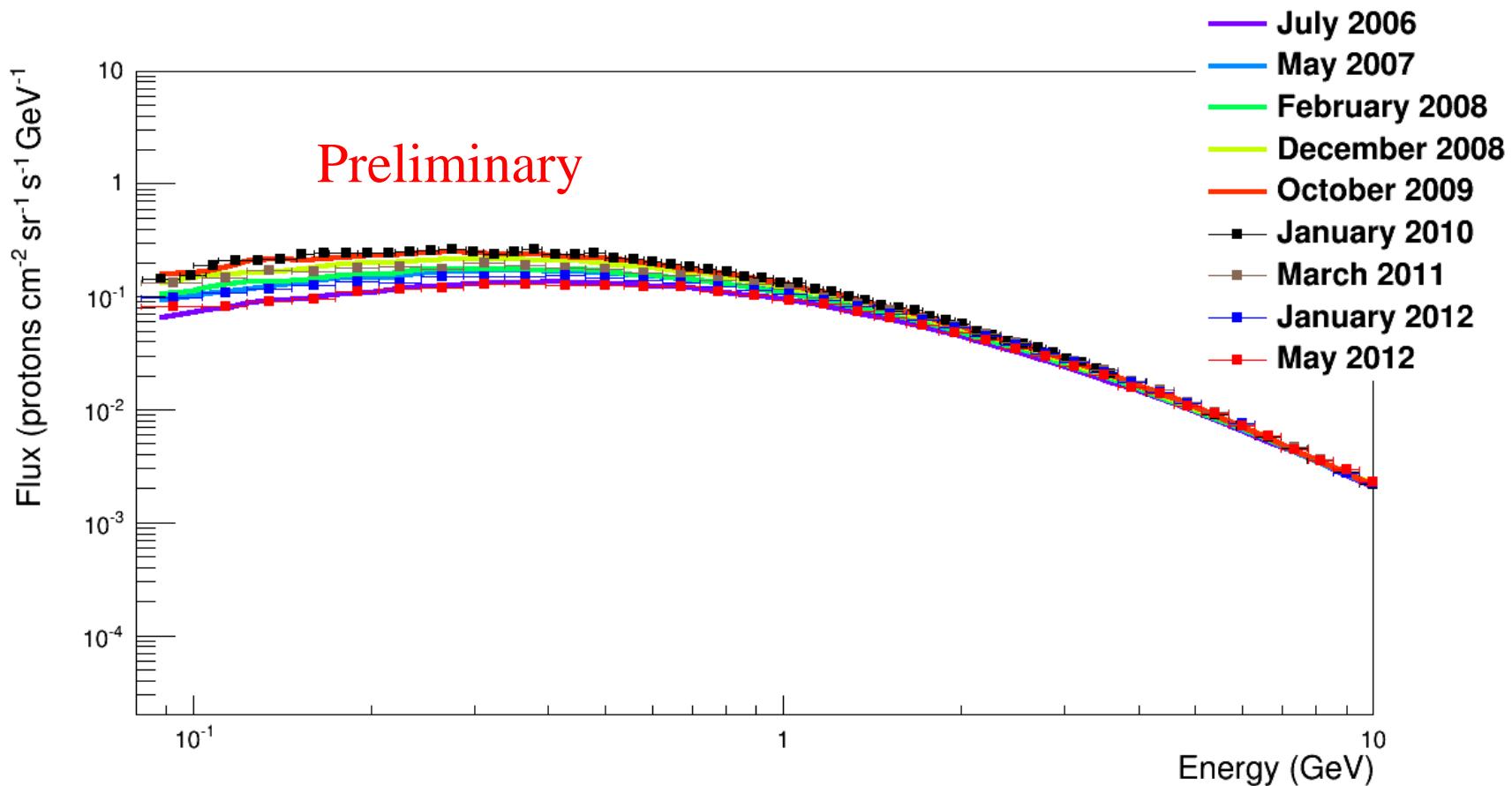
Time dependance of the proton flux

July 2006-January 2014



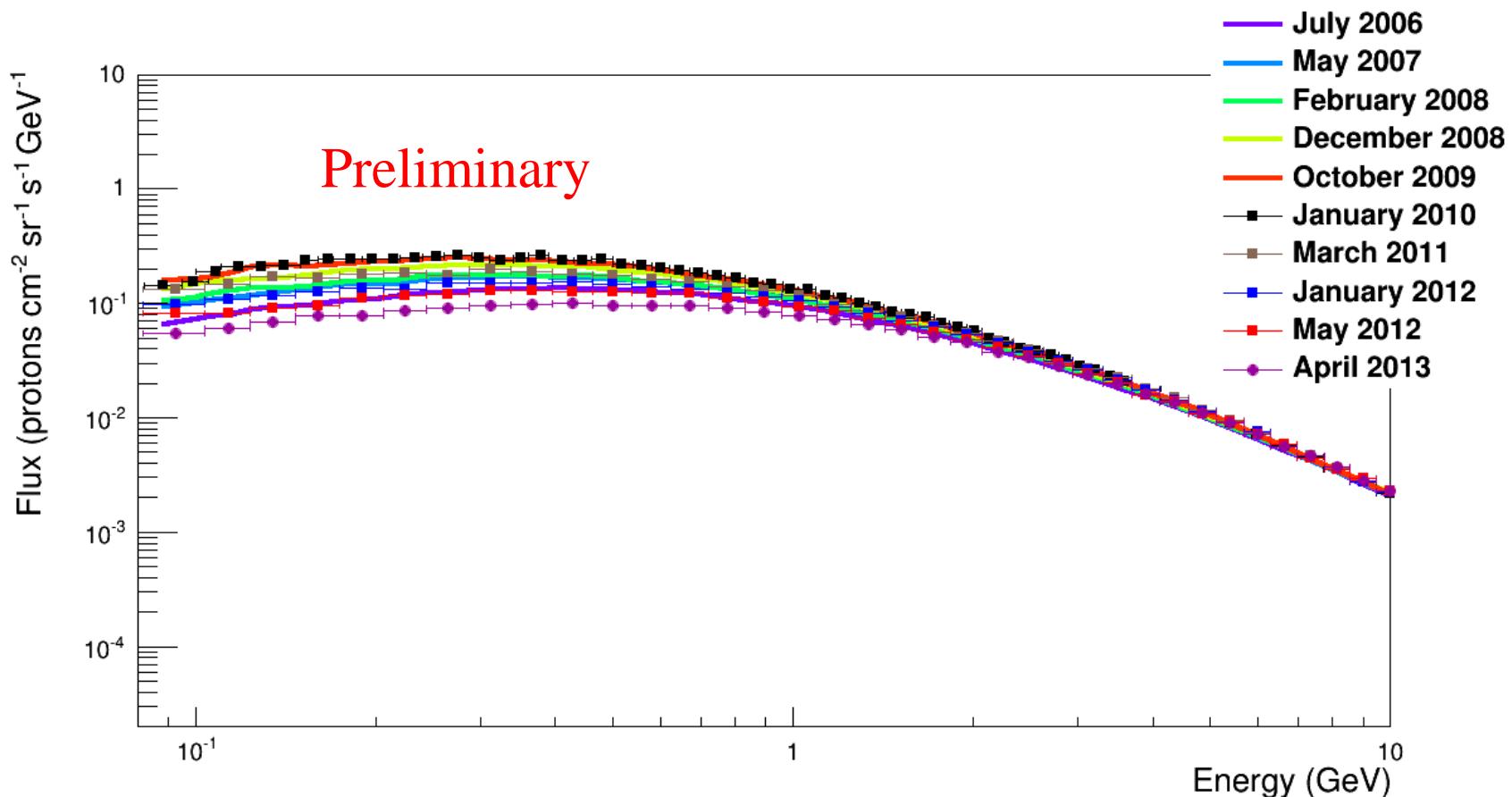
Time dependance of the proton flux

July 2006-January 2014



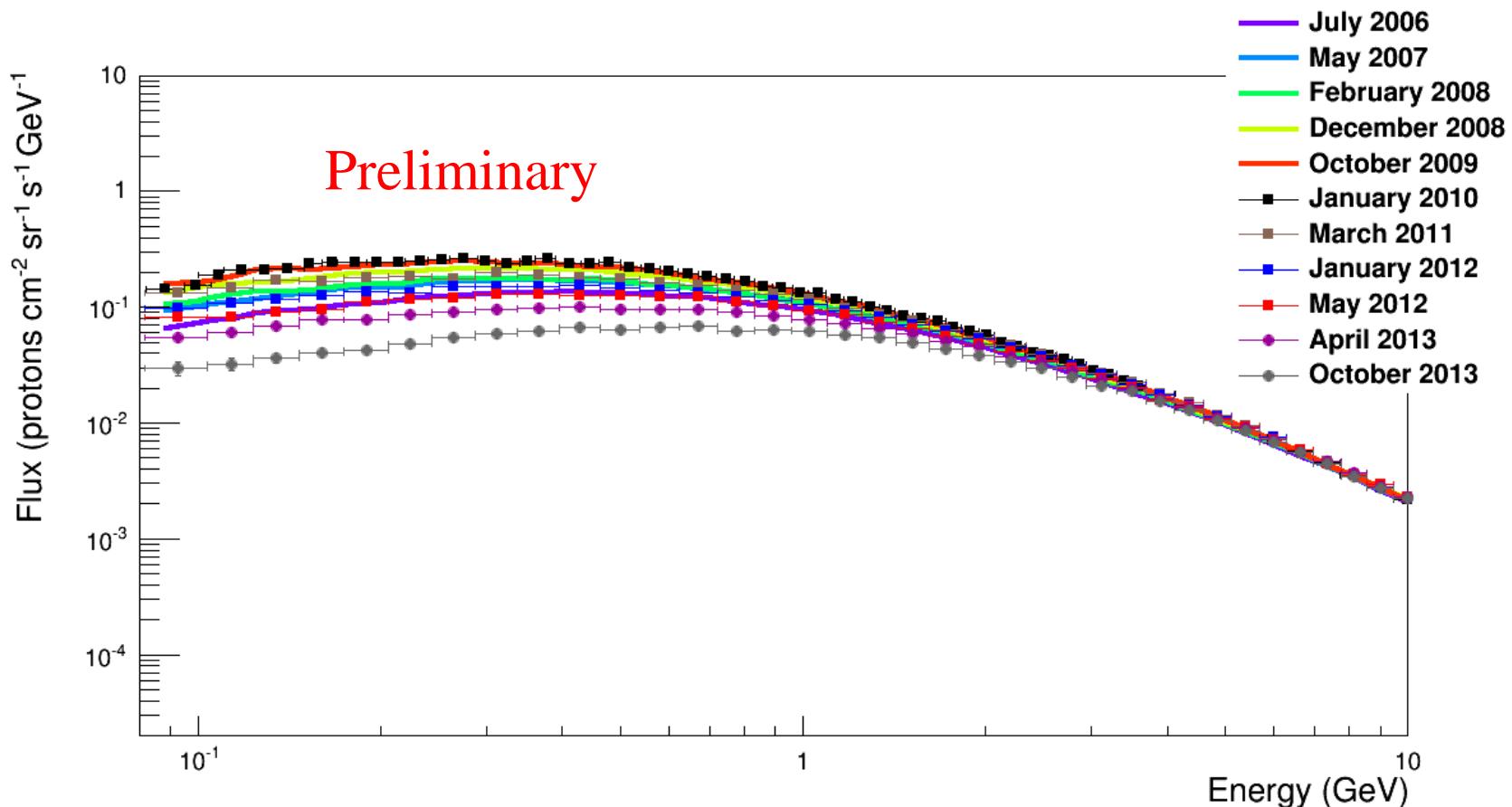
Time dependance of the proton flux

July 2006-January 2014



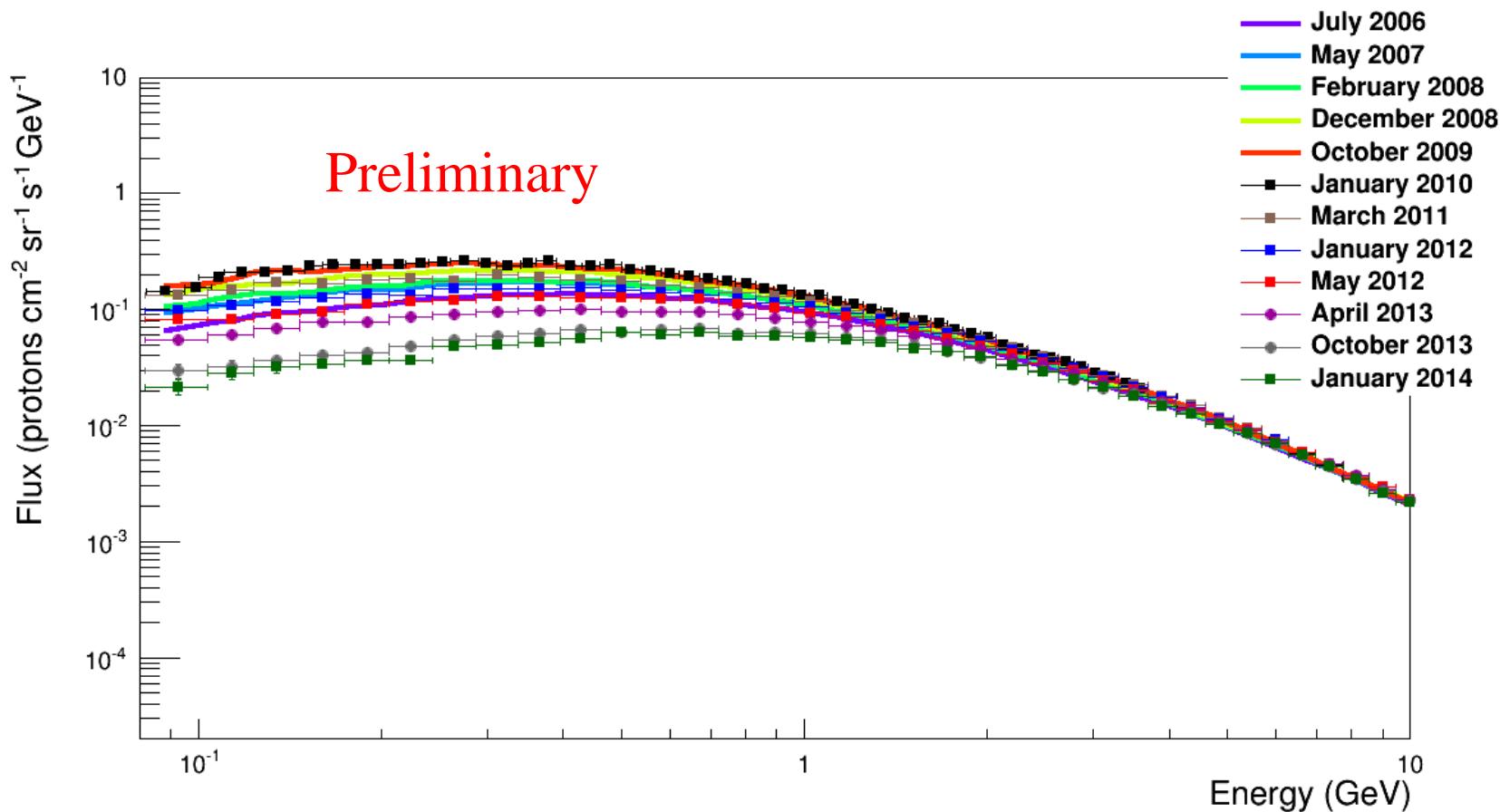
Time dependance of the proton flux

July 2006-January 2014



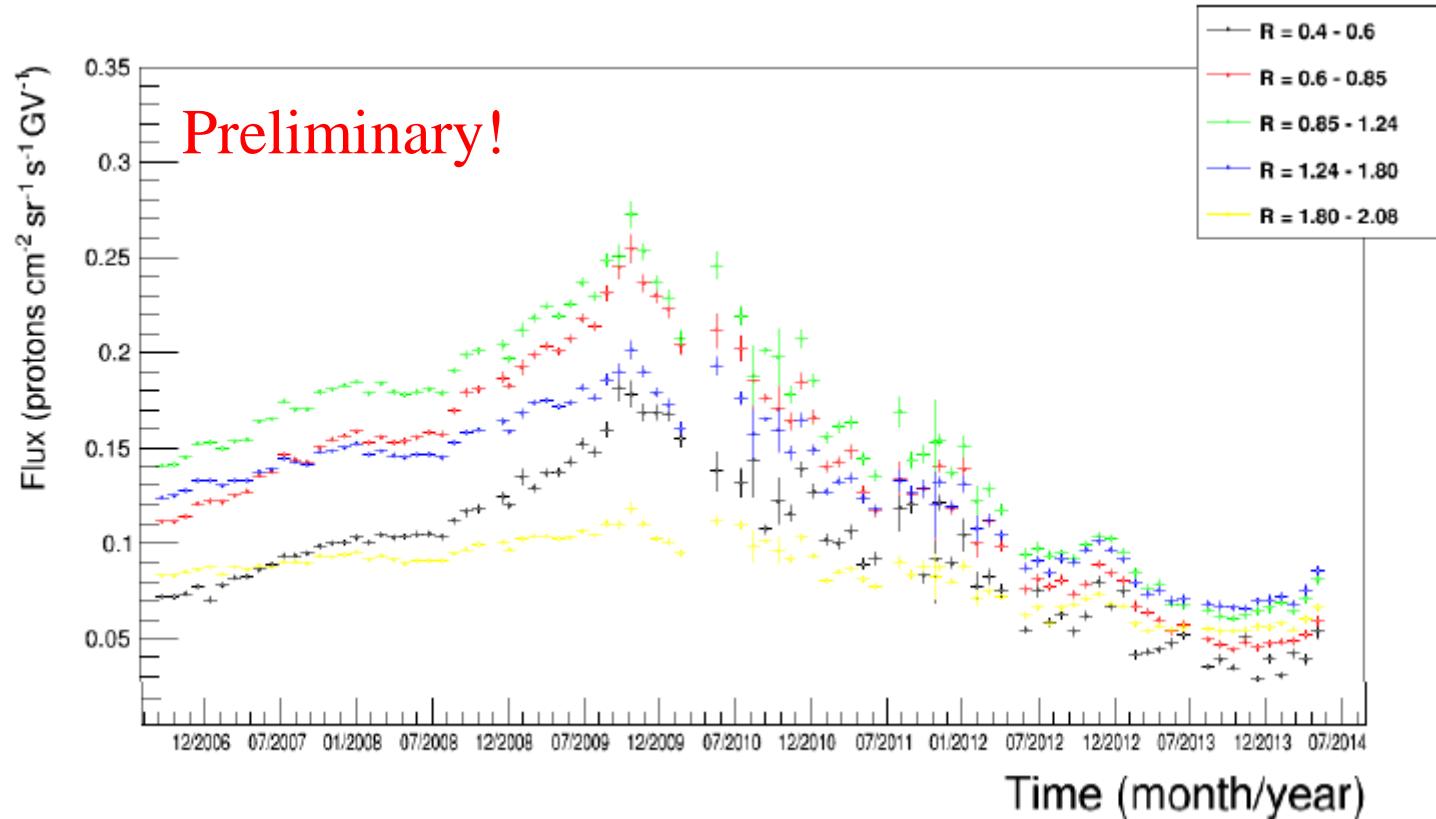
Time dependance of the proton flux

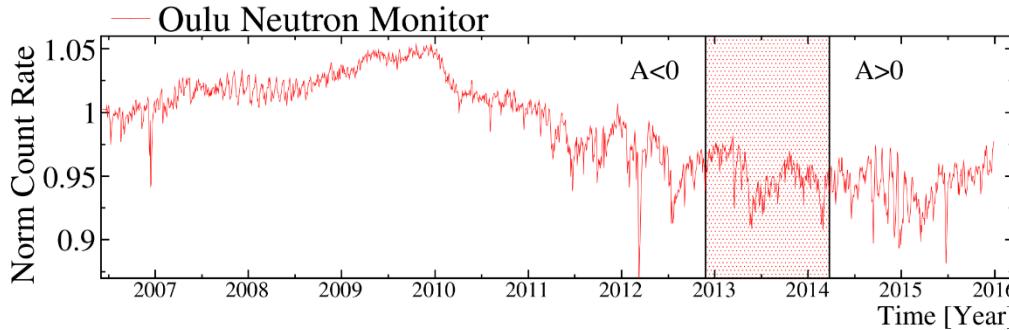
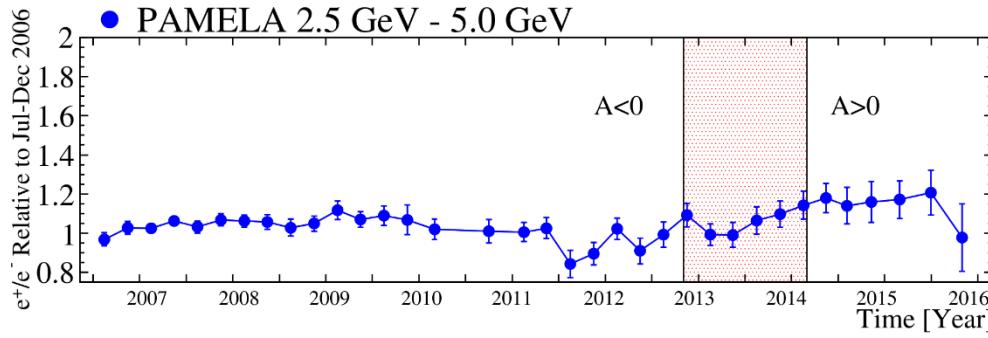
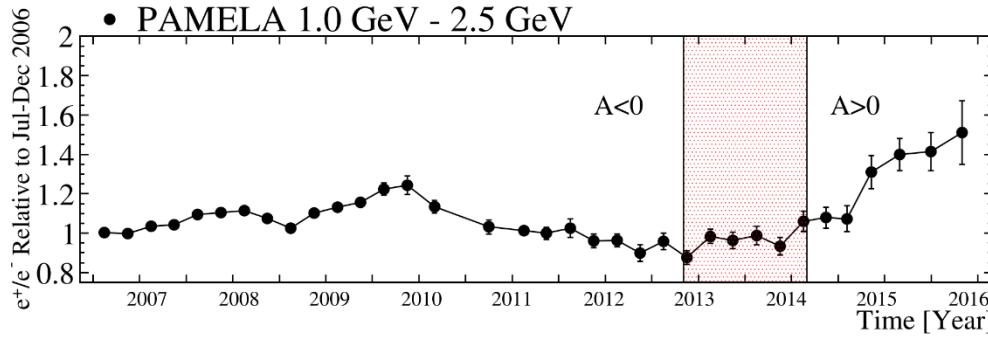
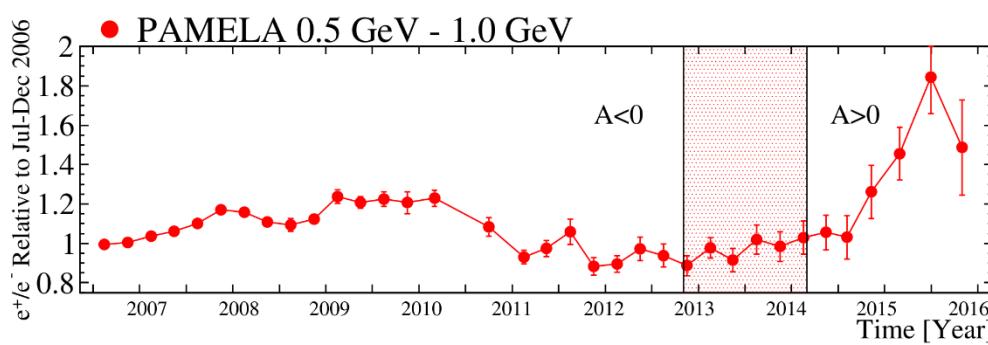
July 2006-January 2014



Time dependance of the proton flux

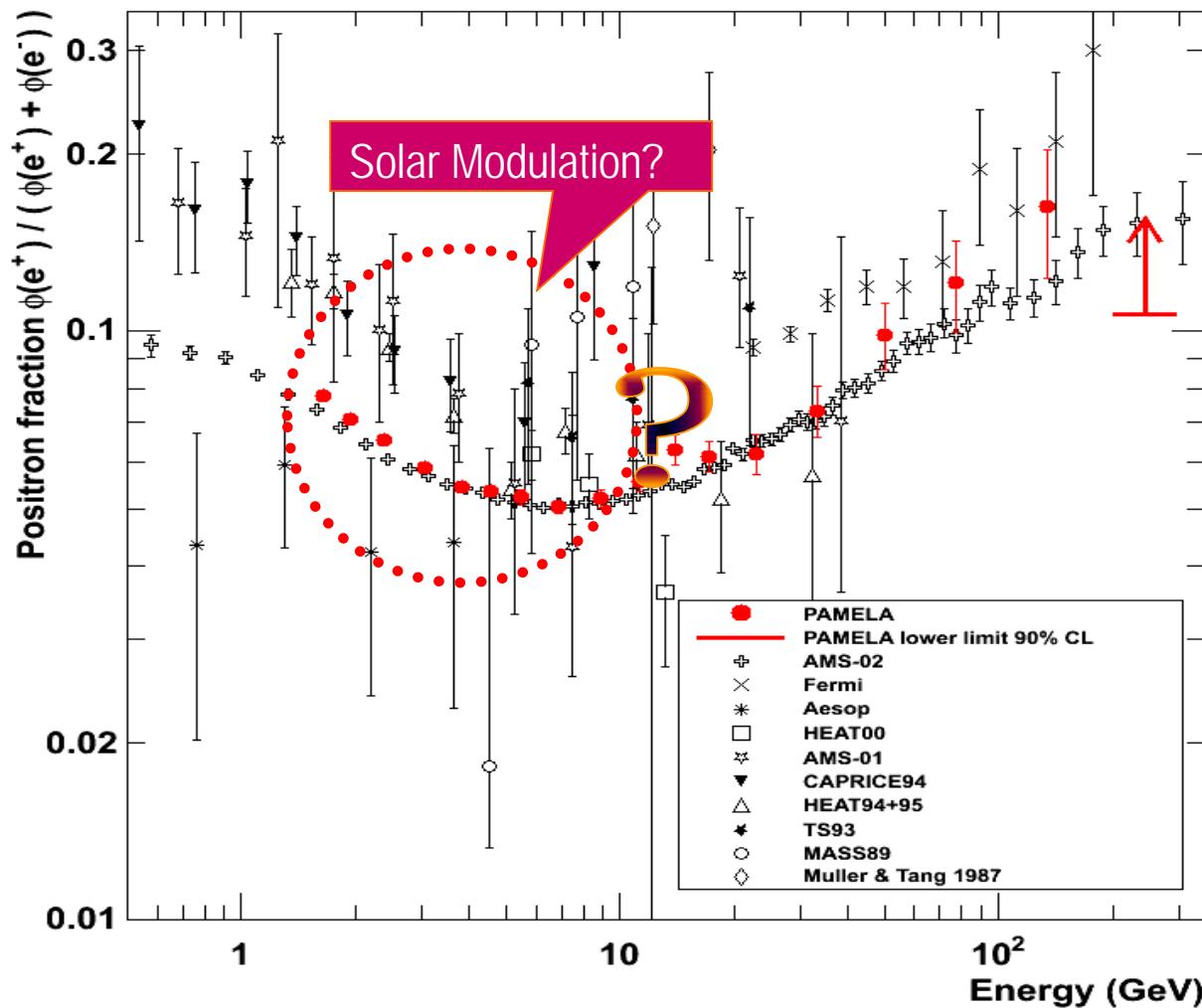
July 2006-January 2014



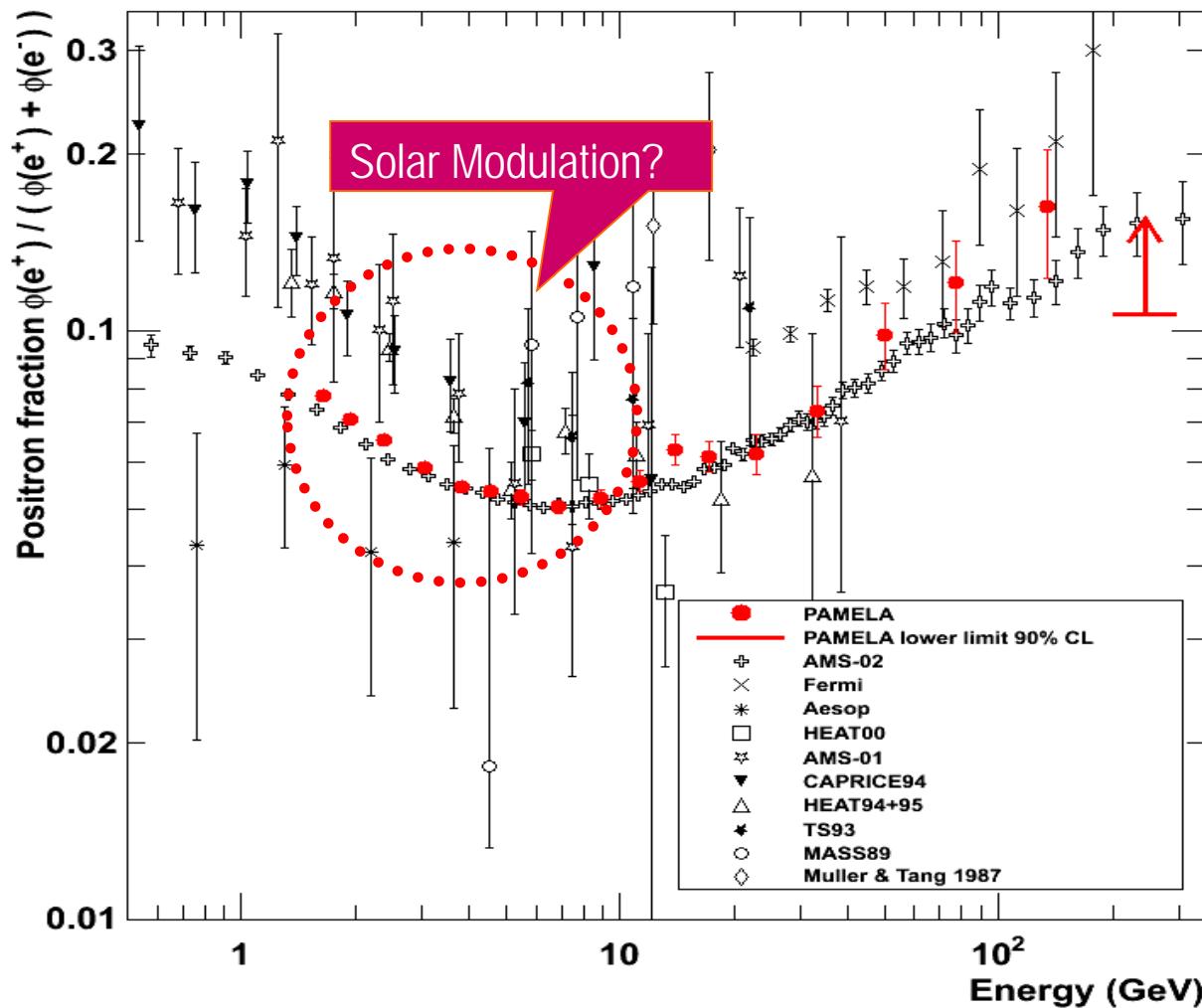


**Time dependance
of the electron and
positron fluxes**

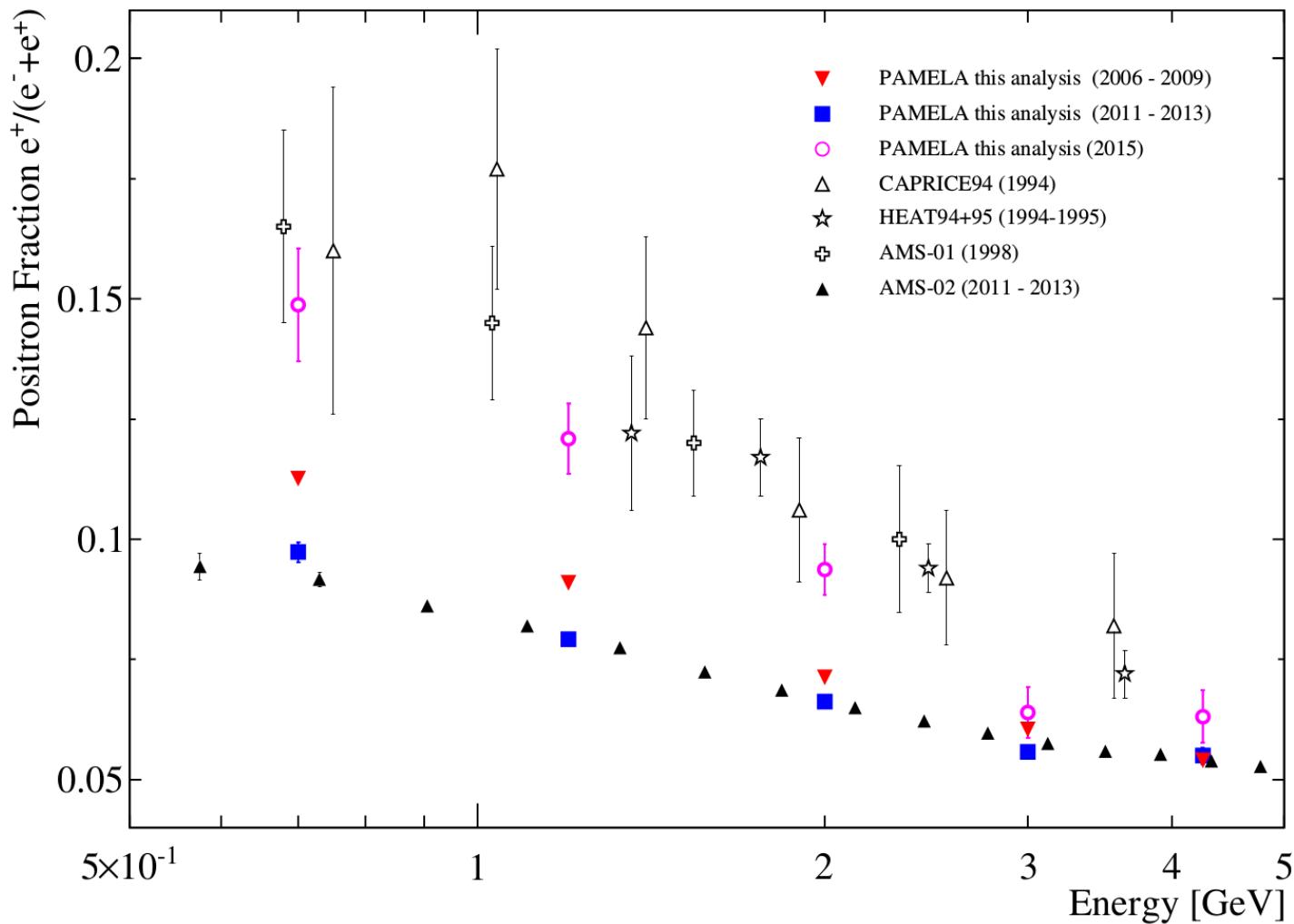
Positron to Electron Fraction



Positron to Electron Fraction



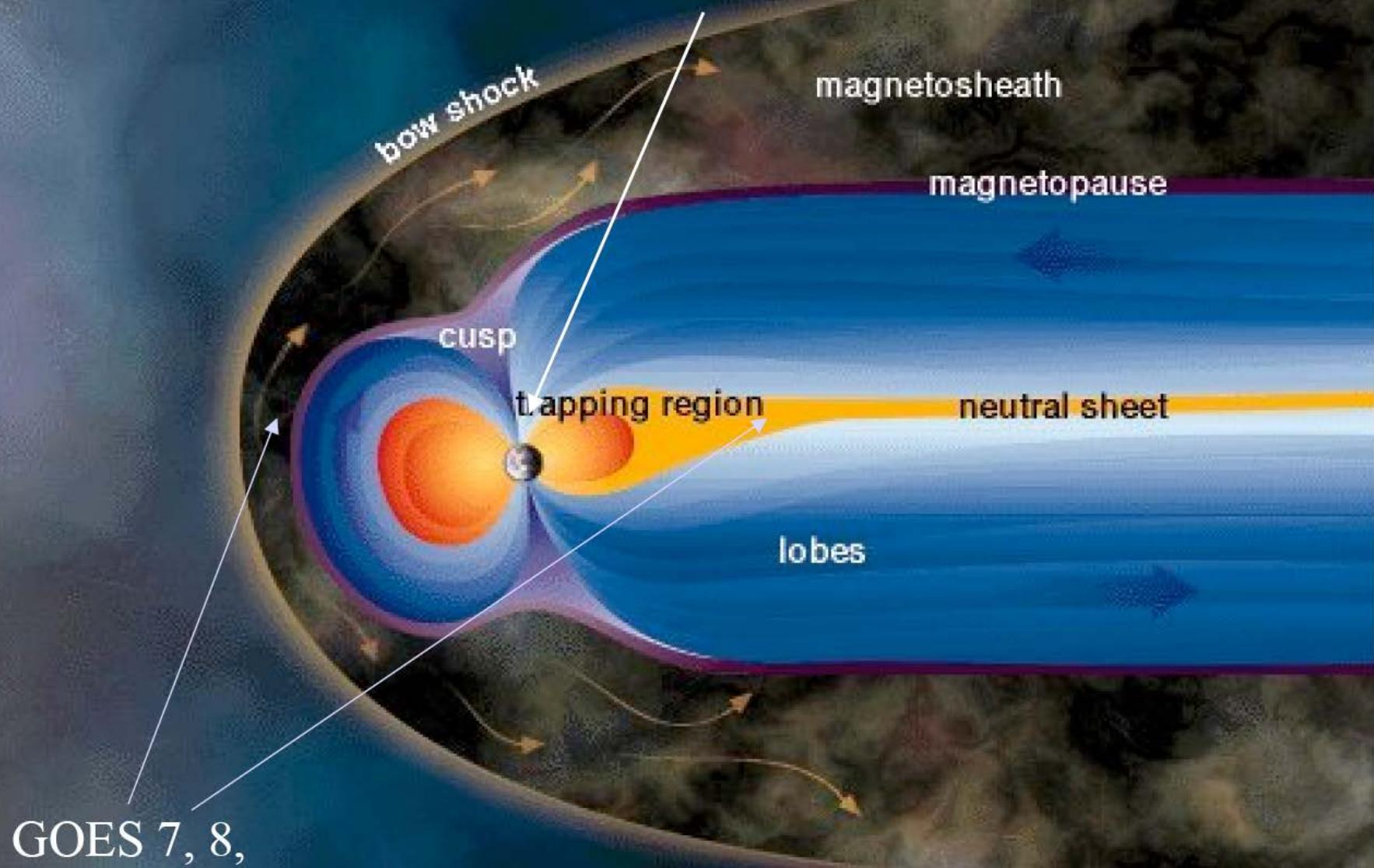
Time dependance of the electron and positron fluxes



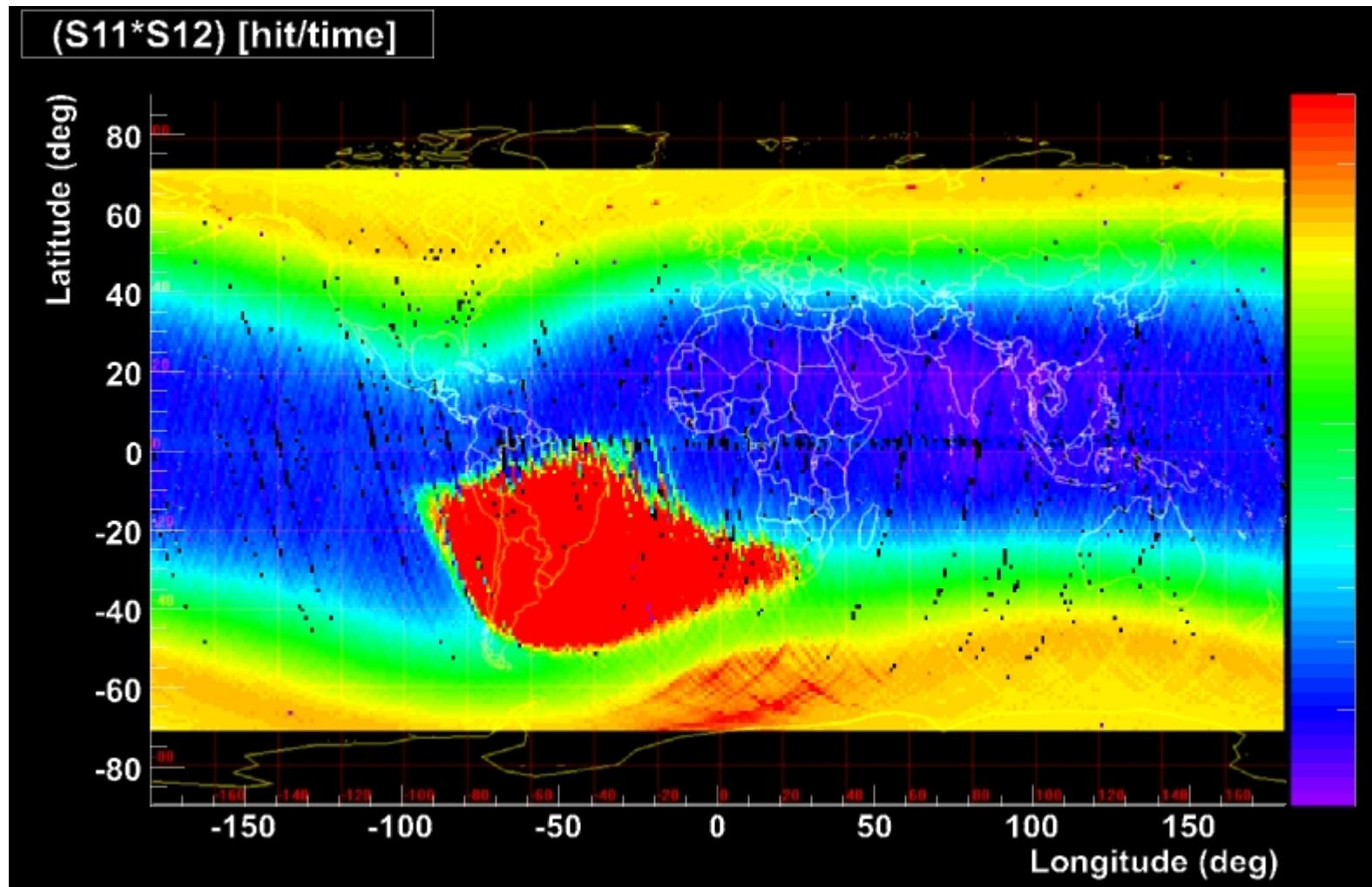
Earth Magnetosphere

Earth's magnetosphere

Sampex, Imp 8, NINA

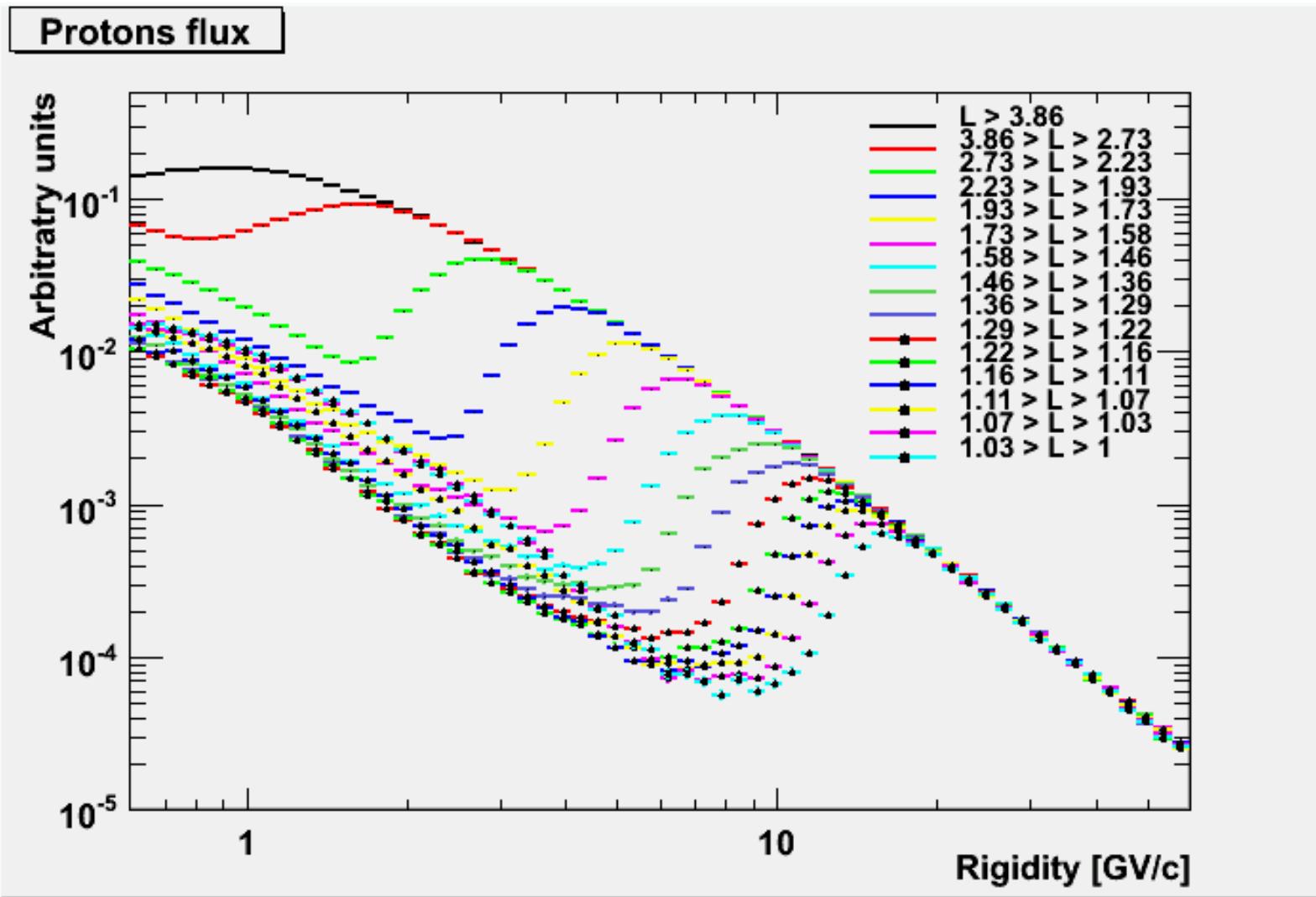


Pamela World Maps: 350 - 650 km alt

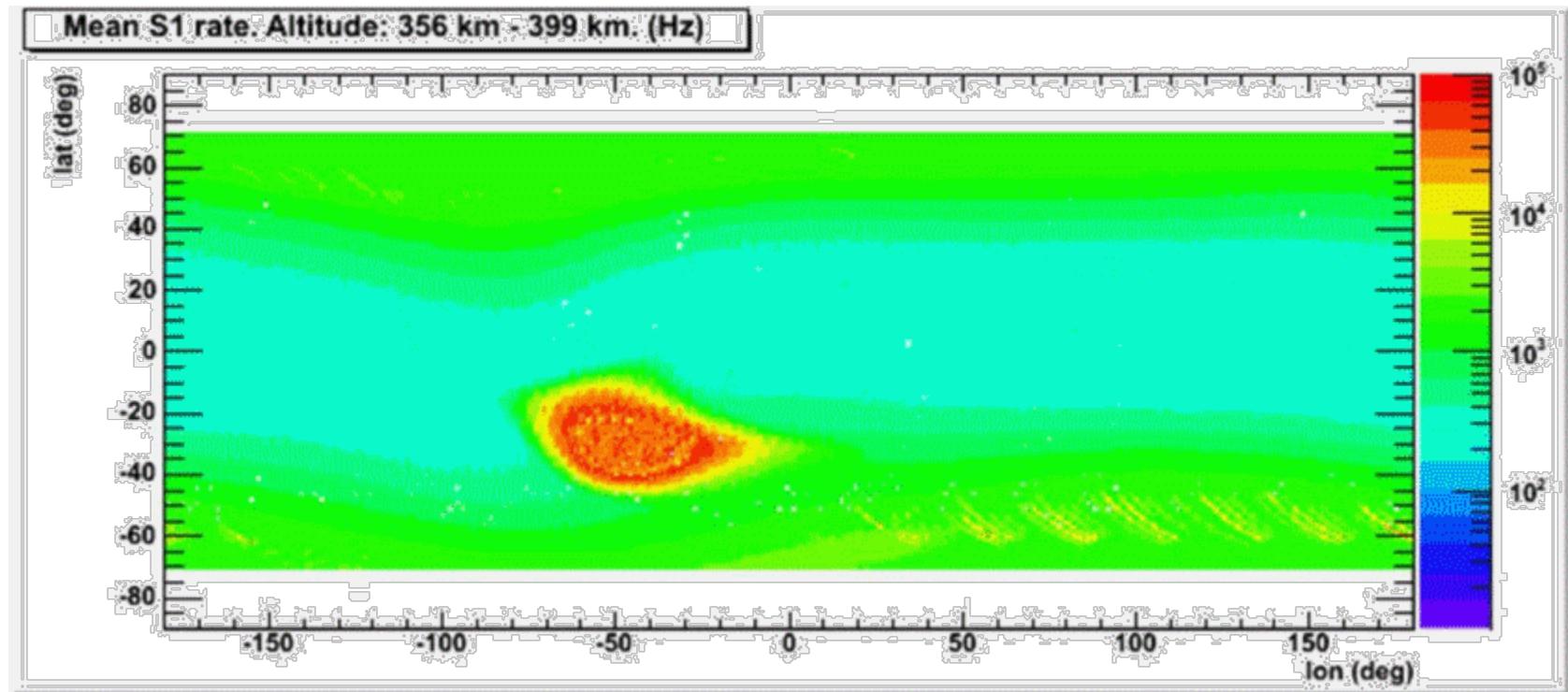


36 MeV p, 3.5 MeV e⁻

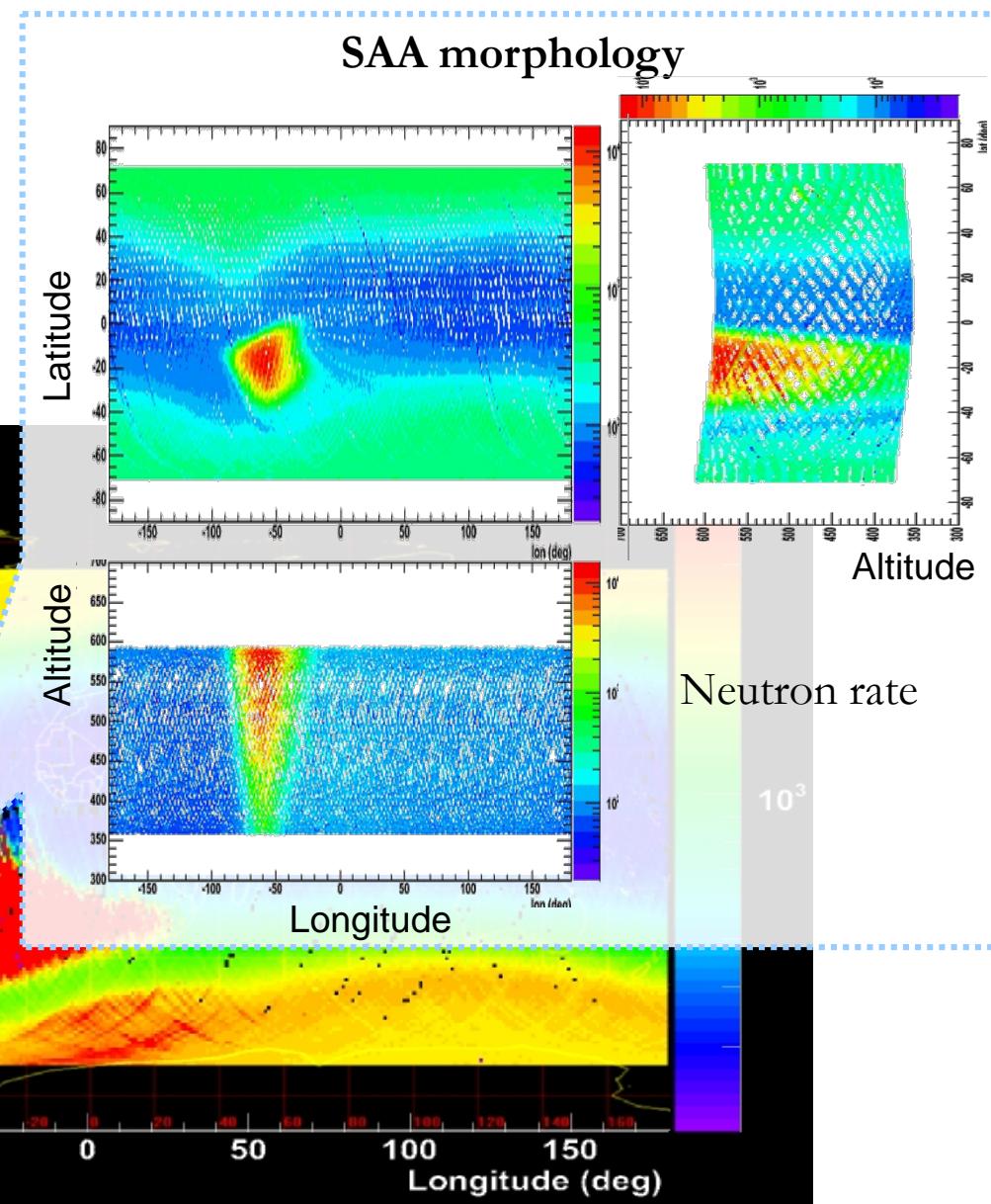
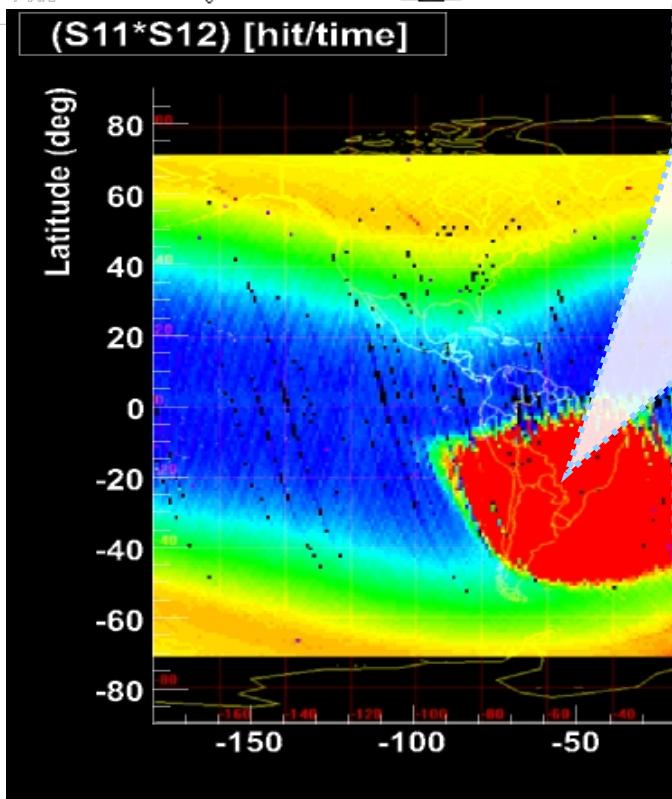
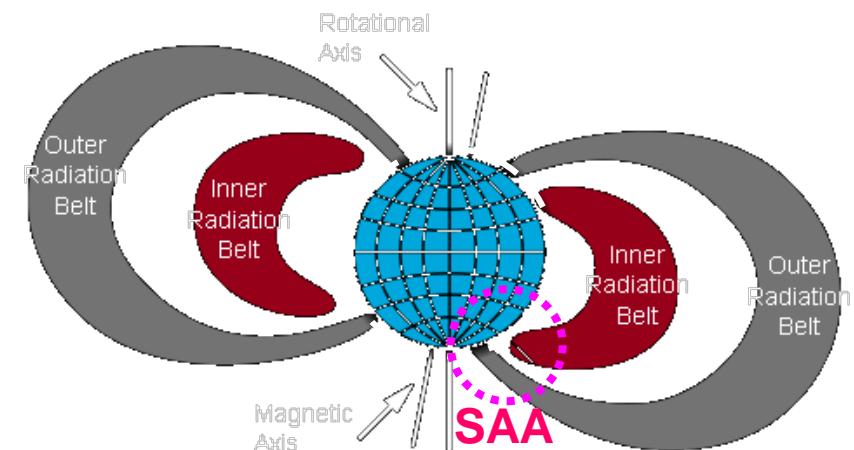
Subcutoff particles

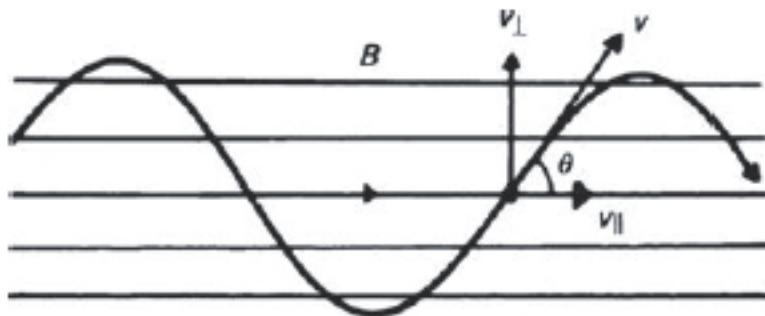


Pamela maps at various altitudes



South-Atlantic Anomaly (SAA)





$$r = \frac{\gamma m_0 v \sin \theta}{ZeB} = \left(\frac{pc}{Ze} \right) \frac{\sin \theta}{Bc}$$

A charged particle gyrating about its guiding centre in a magnetic field is equivalent to a current loop. The equivalent current is the rate at which charge passes a particular point in the loop per second, $i = zev_{\perp}/2\pi r$. The area of the loop is $A = \pi r^2$ and so the magnetic moment μ of the current loop is

$$\mu = iA = \frac{zev \sin \theta}{2\pi r} \pi r^2 = \frac{zev_{\perp}}{2} r .$$

In the non-relativistic limit, $r = m_0 v_{\perp} / zeB$, and therefore

$$\mu = \frac{m_0 v_{\perp}^2}{2B} = \frac{w_{\perp}}{B} , \quad (7.7)$$

where w_{\perp} is the kinetic energy of the particle in the direction perpendicular to the guiding centre.

Magnetic mirroring

$$\Delta B \neq 0 \text{ and } \frac{\Delta B}{B} \ll 1 \rightarrow emf = A \frac{dB}{dt} = \xi$$

Work done on the particle per orbit by emf:

$$Ze\xi = Ze\pi r^2 \frac{dB}{dt} = Ze\pi r^2 \frac{\Delta B}{\Delta T}, \Delta T = \frac{2\pi r}{v_{\perp}}$$

Change in kinetic energy:

$$\Delta w_{\perp} = \frac{ze rv_{\perp}}{2} \Delta B = \frac{m_0 v_{\perp}^2}{2B} \Delta B = \frac{w_{\perp}}{B} \Delta B.$$

The corresponding change in the magnetic moment of the current loop is

$$\Delta \mu = \Delta \left(\frac{w_{\perp}}{B} \right) = \frac{\Delta w_{\perp}}{B} - \frac{w_{\perp} \Delta B}{B^2} = \frac{\Delta w_{\perp}}{B} - \frac{\Delta w_{\perp}}{B} = 0,$$

that is, the magnetic moment of the particle is an *invariant* provided the field is slowly varying. There are other ways of expressing this important result. As illustrated by (7.8), $\Delta \mu = 0$ is equivalent to $\Delta(w_{\perp}/B) = 0$. Since $w_{\perp} = p_{\perp}^2/2m_0$, this is the same as

$$\Delta(p_{\perp}^2/B) = 0.$$

Magnetic mirroring

$$\Delta \left(\frac{p_{\perp}^2}{B} \right) = 0$$

